



Investigating the role of typhoon-induced waves and stratospheric hydration in the formation of tropopause cirrus clouds observed during the 2017 Asian monsoon

Amit Kumar Pandit¹, Jean-Paul Vernier^{1,2}, Thomas Duncan Fairlie^{2,†}, Kristopher M. Bedka²,
Melody A. Avery², Harish Gadhavi³, Madineni Venkat Ratnam⁴, Sanjeev Dwivedi⁵,
Kasimahanthi Amar Jyothi⁶, Frank G. Wienhold⁷, Holger Vömel⁸, Hongyu Liu^{1,2}, Bo Zhang^{1,2},
Buduru Suneel Kumar^{9,†}, Tra Dinh¹⁰, and Achuthan Jayaraman^{a,☆}

¹National Institute of Aerospace, Hampton, VA, USA

²Chemistry and Dynamics Branch, NASA Langley Research Center, Hampton, VA, USA

³Space and Atmospheric Sciences Division, Physical Research Laboratory, Ahmedabad, India

⁴Aerosols, Radiation and Trace Gases Group, National Atmospheric Research Laboratory (NARL),
Gadanki, India

⁵Meteorological Centre, India Meteorological Department, Ministry of Earth Sciences, Bhubaneswar, India

⁶National Centre for Medium Range Weather Forecasting, Noida, India

⁷Institute for Atmospheric and Climate Science (IAC), ETH Zurich, Zurich, Switzerland

⁸Earth Observing Laboratory, National Center for Atmospheric Research, Boulder, CO, USA

⁹Tata Institute of Fundamental Research Balloon Facility, Hyderabad, India

¹⁰Department of Physics, University of Auckland, Auckland, New Zealand

^aformerly at: Aerosols, Radiation and Trace Gases Group, National Atmospheric Research
Laboratory (NARL), Gadanki, India

☆retired

†deceased

Correspondence: Amit Kumar Pandit (amitkandit86@gmail.com)

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Abstract. We investigate the formation mechanism of a tropopause cirrus cloud layer observed during the Balloon measurement campaigns of the Asian Tropopause Aerosol Layer (BATL) over Hyderabad (17.47° N, 78.58° E), India, on 23 August 2017. Simultaneous measurements from a backscatter sonde and an optical particle counter on board a balloon flight revealed the presence of a subvisible cirrus cloud layer (optical thickness ~ 0.025) at the cold-point tropopause (temperature ~ -86.4 °C, altitude ~ 17.9 km). Ice crystals in this layer are smaller than 50 μm with a layer mean ice crystal number concentration of about 46.79 L^{-1} . Simultaneous backscatter and extinction coefficient measurements allowed us to estimate the range-resolved extinction to backscatter coefficient ratio (lidar ratio) inside this layer with a layer mean value of about 32.18 ± 6.73 sr, which is in good agreement with earlier reported values at similar cirrus cloud temperatures. The formation mechanism responsible for this tropopause cirrus is investigated using a combination of three-dimensional back trajectories, satellite observations, and ERA5 reanalysis data. Satellite observations revealed that the overshooting convection associated with a category 3 typhoon, Hato, which hit Macau and Hong Kong on 23 August 2017, injected ice into the lower stratosphere. This caused a hydration patch that followed the Asian summer monsoon anticyclone to subsequently move towards Hyderabad. The presence of tropopause cirrus cloud layers in the cold temperature anomalies and updrafts along the back trajectories suggested the role of typhoon-induced waves in their formation. This case study highlights the role of typhoons in influencing the formation of tropopause cirrus clouds through stratospheric hydration and waves.

1 Introduction

Cirrus clouds are composed of non-spherical ice crystals which exhibit high variability in their size and shapes. Optically thin cirrus coverage is highest over the tropics in a region between the top of maximum convective outflow (between 12 and 14 km) and the cold-point tropopause (CPT, between 16 and 18 km) called the tropical tropopause layer (TTL; Randel and Jensen, 2013). Due to their frequent occurrence at high altitudes and cold temperatures, as well as their thin wispy structures, cirrus clouds trap longwave terrestrial radiation more efficiently than they reflect incoming short-wave solar radiation (Lohmann and Gasparini, 2017). Therefore, high thin cirrus clouds induce a net warming impact on the climate system (Gasparini and Lohmann, 2016; Hong et al., 2016). The radiative effects of cirrus clouds depend on their macrophysical (coverage, altitude, geometrical thickness) and microphysical properties (number concentration, mass density, and size and shape distributions of ice crystals) (Liou, 1986, 2005). Higher cirrus clouds have a larger warming effect on climate than those at lower levels (Lohmann and Gasparini, 2017). Simulations using general circulation models (GCMs) have shown that the radiative impact of cirrus clouds is sensitive to small changes in the number concentration of small ice crystals (Sanderson et al., 2008; Mitchell et al., 2008). In this context, subvisible cirrus clouds (having optical thickness less than 0.03 at visible wavelengths) occurring near the CPT are important because they are the highest type of cirrus clouds with significant horizontal coverage. Subvisible cirrus clouds can have a relatively long lifetime because they usually consist of low concentrations of small ice crystals, which are more slowly removed by sedimentation and yield greater radiative heating than larger ones, given the same ice mass (Fu and Liou, 1993).

Apart from their radiative impact, subvisible cirrus clouds play an important role in regulating the water vapour in the upper troposphere and lower stratosphere (UTLS) region through dehydration (Jensen et al., 1996). Dehydration efficiency depends on microphysical processes such as nucleation, growth, and sedimentation (Rollins et al., 2016). Understanding of microphysical processes under different dynamical forcings and their representation in GCMs remain poor due to numerous factors, one being the lack of accurate measurements of microphysical properties of ice crystals with particle sizes less than 100 μm from current-generation probes (Heymsfield et al., 2017; Baumgardner et al., 2017; Kuhn and Heymsfield, 2016). In situ measurements provide accurate information about the microphysical properties, which are essential for estimating their radiative impact on climate and for the calibration and validation of ice cloud products from passive and active remote sensing sensors. Balloons and aircraft are the most common in situ measurement platforms for ice crystal microphysical properties.

Due to their slow ascent, balloons provide cloud vertical profiles at high resolution, unlike aircraft which provide larger horizontal coverage but miss fine vertical information and detailed imaging of ice crystal habits. Moreover, ice measurements using balloons do not suffer as much from sampling artefacts such as particle shattering as under-wing-mounted aircraft probes (Wolf et al., 2018).

There have been several previous in situ measurements of tropical cirrus clouds using balloon-borne and aircraft-based instruments (Krämer et al., 2016; Jensen et al., 2017; Schoeberl et al., 2019; Krämer et al., 2020, and references therein). However, most of those measurements were over Australia, Brazil, Central America, Europe, and the western Pacific. Only a limited number of field campaigns have occurred over the Asian summer monsoon (ASM) region (Krämer et al., 2020), despite the important role played by the ASM in the transport of aerosols (Vernier et al., 2015, 2018), water vapour (Nützel et al., 2019; Khaykin et al., 2022; Wang et al., 2019), and the formation of cirrus clouds (Ueyama et al., 2018) in the UTLS region. Additionally, the ASM region has the largest coverage of high-altitude cirrus and subvisible cirrus clouds during the boreal summer (Martins et al., 2011). Field campaigns such as the Balloon measurement campaigns of the Asian Tropopause Aerosol Layer (BATAL) (Vernier et al., 2018) and aircraft campaigns such as StratoClim (Krämer et al., 2020) and the Asian Monsoon Chemical and Climate Impact Project (ACCLIP; https://www.eol.ucar.edu/field_projects/acclip, last access: 26 October 2024) have been organized over the ASM region to understand the physical, chemical, and dynamical characteristics of the Asian tropopause aerosol layer (ATAL) along with cirrus cloud microphysics.

Satellite measurements have shown that the UTLS aerosol optical depth between 13 and 18 km altitude over the ASM region has increased 3-fold since the late 1990s (Vernier et al., 2015). A balloon-borne study over Tibetan Plateau (He et al., 2019) also showed evidence for the hygroscopic growth of particles inside the ATAL. These may have a significant impact on the microphysics of cirrus clouds, especially subvisible cirrus clouds that occur in proximity to the CPT during the ASM (Vernier et al., 2018). In this context, long-term (1998–2013) lidar observations over a tropical station in Gadanki (13.5° N, 79.2° E) in southern India have shown an increase in the fraction of subvisible cirrus clouds (Pandit et al., 2015). Also, the StratoClim campaign measurements in July and August 2017 from Nepal have shown the presence of solid ammonium nitrate particles from surface ammonia sources (Höpfner et al., 2019), which act as efficient ice nuclei in the presence of ammonium sulfate particles as shown in cloud chamber experiments (Wagner et al., 2020). In general, anthropogenic activities are expected to influence the occurrence and properties of cirrus clouds by meteorological, aerosol-induced, or cloud-induced changes (Kärcher,

2017). However, these changes over the ASM region, which is one of the most polluted regions of the globe, are not well known and warrant further investigations.

The UTLS region during the ASM is dominated by the complex interplay among frequent deep convection, atmospheric waves, and large-scale updrafts which directly and/or indirectly lead to the formation of cirrus clouds either through liquid-origin clouds or through in situ freezing mechanisms (Krämer et al., 2016). These dynamical processes directly influence the microphysical properties of cirrus clouds. For example, slow updraft leads to thinner cirrus clouds with low ice water content (IWC), while faster updraft leads to thicker cirrus clouds with higher IWC (Krämer et al., 2016). A study (Ueyama et al., 2018) done using a one-dimensional (vertical) time-dependent cloud microphysical model, diabatic back trajectories, and observations of convective clouds has suggested that nearly all the enhancement of water vapour and clouds at the 100 hPa level over the ASM is due to convective saturation with a minimum impact from the convectively detrained ice. This is in contrast to another recent study (Wang et al., 2019) where convective lofting of ice during the ASM is found to be the most important source of water vapour at the 100 hPa level in the 10–40° N region. Khaykin et al. (2022) have recently shown evidence of direct convective hydration (water vapour mixing ratio > 10 ppmv) by overshooting convection over the ASM region using StratoClim aircraft measurements. Such large-scale organized convective systems in the southern ASM anticyclone are also associated with synoptic-scale dehydration near the tropopause. In this context, the role of large-scale organized deep convection over the eastern coast of India in the formation of tropopause cirrus clouds is a subject of interest for the BATAL campaigns organized in Hyderabad (17.47° N, 78.58° E), India (see Sect. 2.1).

In addition to mesoscale deep convection, overshooting convection frequently occurs in tropical cyclones (Romps and Kuang, 2009), especially during their intensification (Horinouchi et al., 2020). Tropical cyclones (typhoons) occurring during the ASM have also been seen to cause dehydration near the CPT (Li et al., 2020) and hydration in the lower stratosphere (Jiang et al., 2020). Climate change is expected to strengthen and increase the occurrence frequency of such typhoons (Stocker et al., 2013). Extreme tropical convection is expected to impact TTL water vapour and cirrus clouds (Aumann et al., 2018). However, a recent climate model study (Smith et al., 2022) suggests that the role of convective ice injection in the stratospheric water vapour budget will remain constrained by large-scale temperatures in a warmer climate. Thus, it is essential to understand how such systems influence the large-scale temperature fields and hence cirrus cloud formation in the TTL. In this context, the ASM, which is influenced by frequent deep convective systems and tropical cyclones, is the preferred region for studying such influence on cirrus clouds as presented here.

In the UTLS region, deep convection generates gravity waves, which are known to influence the cold-point tropopause temperature and cirrus cloud formation and hence cause dehydration in the UTLS region (Schoeberl et al., 2015, 2016). Using radiosonde observations over the tropical western Pacific region, Kim and Alexander (2015) have shown that the CPT temperature is directly modulated by vertically propagating waves irrespective of any change in the mean upwelling. Later, Kim et al. (2016) showed the ubiquitous influence of waves on TTL cirrus clouds using airborne observations of temperature and cirrus clouds obtained over the western Pacific during the Airborne Tropical Tropopause Experiment (ATTREX) campaign. Several studies (Wu et al., 2015; Nolan and Zhang, 2017; Kim et al., 2009) have reported anomalies in temperature and vertical wind speed associated with semi-circular gravity waves generated from the typhoon centre propagating horizontally and vertically through the troposphere, stratosphere, and mesosphere in expanding spirals with horizontal wavelengths typically in the range of 50–500 km and periods from 1 h to 1.6 d.

This paper presents a case study of in situ measurements of microphysical and optical properties of a tropopause cirrus cloud layer observed during the BATAL campaign using balloon-borne instruments (a backscatter sonde and an optical particle counter) capable of detecting very thin cirrus clouds consisting of ice crystals smaller than 100 µm. We also demonstrate the usefulness of combined measurements from these instruments in measuring the range-resolved lidar ratio in situ for a tropical tropopause cirrus cloud layer for the first time (to the best of our knowledge). Using back-trajectory calculations and their intersection with the convective clouds observed by radiometers on the Himawari-8 geostationary satellite, we also investigate the origin of the air masses sampled during the balloon flight. The formation mechanism and properties of tropopause cirrus clouds along the back trajectories are examined using the satellite observations. Section 2 describes the BATAL campaign, balloon-borne instruments, satellite, and model data used in this study. Section 3 discusses the results, and Sect. 4 contains the summary.

2 Data and methods

2.1 BATAL campaign and general meteorological conditions

The balloon measurements were carried out under the framework of the India Space Research Organization (ISRO)–National Aeronautics and Space Administration (NASA) joint project called BATAL (Vernier et al., 2018) during August 2017 from Tata Institute of Fundamental Research Balloon Facility (TIFR-BF, <https://www.tifr.res.in/~bf/>, last access: 26 October 2024) located at 17.47° N, 78.58° E in Hyderabad, India. During the ASM period (June to August),

deep convective clouds developing over the Bay of Bengal (BoB) and over land, especially over the Indian east coast, cause precipitation in the late afternoon and early evening hours. Climatologically, Hyderabad receives the highest rainfall during the months of July and August due to frequent deep convection. Launching a heavy payload on a plastic balloon flight from TIFR-BF during this time is thus quite challenging due to strong south-westerly surface winds and frequent precipitation. During this period, the tropical easterly jet (TEJ) with maximum wind speeds of about 40 m s^{-1} prevails over Hyderabad in the UTLS region (Vernier et al., 2018). The TEJ brings the moisture from the Bay of Bengal, the western Pacific, and Southeast Asia (Das et al., 2011) necessary for cirrus cloud formation. Also, the tropopause over Hyderabad is extremely cold, with a minimum temperature reaching about -86°C as observed during the BATAL campaigns (Vernier et al., 2018), much colder than northern India. The 10-year (2011–2020) mean temperature at the 100 hPa pressure level ($\sim 16.5 \text{ km}$ altitude) from ERA5 (described in Sect. 2.4.3) during the month of August shows a regional pool of cold air, with temperature between -78 and -80°C between 10 and 20°N from the South China Sea to the Indian east coast, with Hyderabad located at the eastern edge (Fig. 1a). This cold region, often influenced by deep convective activities and the TEJ, seems to be a hotspot for laminar cirrus clouds as evident from the observations of two spaceborne lidars (described in Sect. 2.3.1 and 2.3.2), as shown in Fig. 1b for the month of August 2017. The occurrence of these cirrus clouds with their base altitude at or above 16 km is associated with negative temperature anomalies estimated from the ERA5 temperature field by subtracting the monthly mean temperature from each grid ($0.25^\circ \times 0.25^\circ$) in this cold pool (Fig. 1b). All these factors make TIFR-BF an ideal location for cirrus cloud studies during the ASM. A total of nine balloon flights were launched from TIFR-BF during August 2017 (Vernier et al., 2018), and the timings of these flights are shown in Fig. 1b. The balloon-borne instruments used for measuring the cloud properties are described in Sect. 2.2 below.

2.2 Balloon-borne observations

On the night of 23 August 2017 at 20:03 UTC (01:33 LT), balloon-borne instruments on board a heavy flight (HF; see Vernier et al., 2018, for description) were launched from TIFR-BF to measure the vertical distribution of aerosols, clouds, and meteorological parameters as described in the subsections below. Satellite and model data have also been used in this study to support the balloon observations, which are described in the following subsections.

2.2.1 Compact Optical Backscatter Aerosol Detector (COBALD)

COBALD is a unique lightweight (about 500 g) balloon-borne backscatter sonde developed at the Institute of Atmospheric and Climate Science, Swiss Federal Institute of Technology (ETH), Zurich. COBALD enables us to detect optically thin layers of aerosols and clouds at high vertical resolution (Brabec et al., 2012; Cirisan et al., 2014; Reinares Martínez et al., 2021; Vernier et al., 2015, 2016, 2018; Brunamonti et al., 2018). Additional details on COBALD are presented in Appendix A1. COBALD gives the backscatter ratio (BSR) at two wavelengths (455 and 940 nm), which is an optical analogue for the particle mixing ratio. The ratio of $(\text{BSR}_{940} - 1)$ at 940 nm to $(\text{BSR}_{455} - 1)$ at 455 nm , defined as the colour index (CI), gives qualitative information about particle size (Cirisan et al., 2014; Vernier et al., 2015). On substituting the molecular backscatter coefficient at 455 and 940 nm wavelengths, CI is reduced to

$$\text{CI} = 18 \times \left(\frac{\beta_{p940}}{\beta_{p455}} \right) = 18 \times \text{CR}, \quad (1)$$

where β_{p455} and β_{p940} are particulate backscatter coefficients at 455 and 940 nm , respectively, and the ratio $\frac{\beta_{p940}}{\beta_{p455}}$ is called the particulate colour ratio (CR). A CR value near unity indicates the presence of cloud layer, whereas values below 0.7 indicate aerosol presence (Vernier et al., 2015; Brunamonti et al., 2018). The Ångström exponent estimated from the BSR_{940} and BSR_{455} is used to estimate the backscatter ratio at 532 nm (BSR_{532}) following Vernier et al. (2015). After applying the field-of-view (FOV) corrections to the COBALD backscatter coefficient as suggested by Brunamonti et al. (2021), an extinction coefficient profile is obtained at 532 nm wavelength by using the BSR_{532} profile and assuming a constant extinction to backscatter coefficient ratio (lidar ratio) for cirrus clouds following the CALIOP L2 V4 algorithm where the initial lidar ratio is a sigmoid function of centroid temperature $\pm 10 \text{ sr}$ (Young et al., 2018). By integrating these extinction coefficients between the base and top altitudes of the cloud layer, the optical thickness (τ) of the cloud layer is estimated. The extinction coefficient derived from COBALD is used for estimating the IWC of the tropopause cirrus cloud layer using a temperature-dependent parameterized relation derived from in situ aircraft measurements (Heymsfield et al., 2014) as described in the next subsection (Sect. 2.2.2).

2.2.2 SOLAIR Boulder Counter and the estimation of cloud microphysical properties

SOLAIR Boulder Counter

The SOLAIR Boulder Counter is a $\sim 6 \text{ kg}$, portable, forward-scattering-based particle counter built by Lighthouse Worldwide Solutions, USA (<https://www.golighthouse.com/zh/airborne-particle-counters/boulder-counter/>, last access:

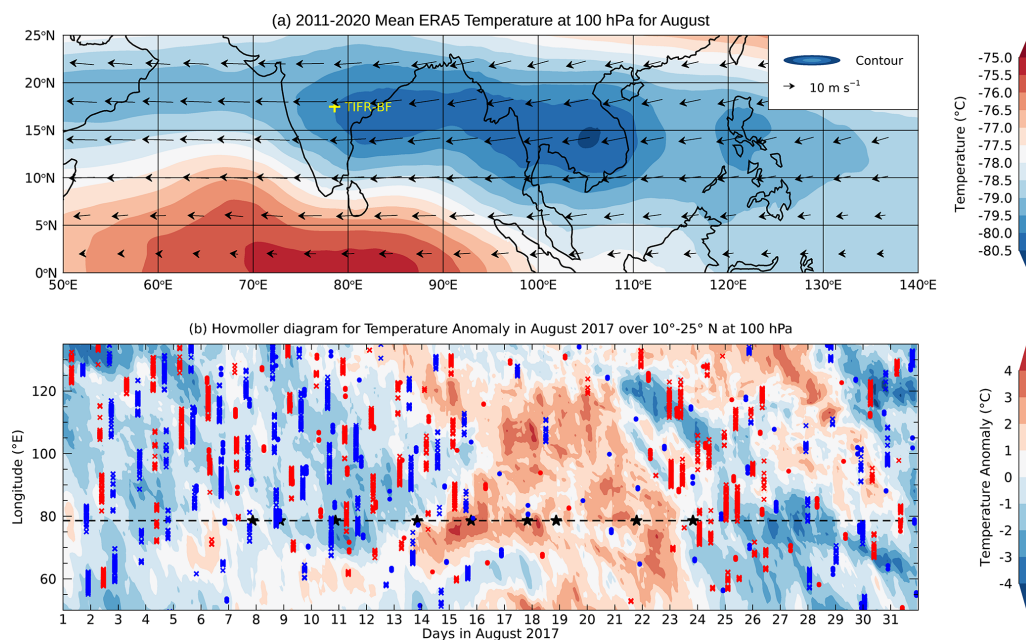


Figure 1. (a) Filled colour contour showing the 10-year (2011–2020) mean ERA5 temperature superimposed, with wind vectors representing the mean ERA5 horizontal wind speed and direction at the 100 hPa pressure level for the month of August. The yellow “plus” sign indicates the location of TIFR-BF in Hyderabad, India. (b) Hovmöller plot for the anomaly in temperature from ERA5 temperature at the 100 hPa pressure level for August 2017. The anomaly is computed by subtracting the monthly mean (for August 2017) temperature for each grid between 10 and 25° N. The horizontal dashed black line marks the longitude (78.58° N) of TIFR-BF, with each black star showing the launch time of the balloon. The filled red (blue) circles and crosses superimposed on the filled colour contours represent the locations of cirrus cloud layers detected by CALIOP and CATS during daytime (nighttime), respectively, with their base altitude located at 16 km or above.

26 October 2024), which uses an extreme-life laser diode at 680 nm as a light source. Particles are sampled at a high flow rate of 28.3 L min^{-1} and are directed to the laser beam. The light scattered at 45° angles from the particles is collected and focused by collection optics onto a photodiode, which converts it into electric pulses. The pulse amplitude is a measure of particle size. The number of pulses as a function of pulse height gives particle counts and size information. Particles are counted in six size channels between 5 and $100 \mu\text{m}$: 5, 10, 25, 40, 50, and $100 \mu\text{m}$, with 50% counting efficiency at $5 \mu\text{m}$ and 100% for the other size bins. The Boulder Counter’s ability to measure large particles allows us to infer the distribution of larger airborne particles such as cloud droplets and ice crystals from a balloon platform. Particle counts in each bin are recorded on board every 5 s, which when divided by the instantaneous sample flow gives the particle number concentration for that size bin at that time. Boulder Counter measurements were synchronized with COBALD measurements. The flow is controlled and monitored through an external mass-flow controller system. During the flight, the sample flow decreased gradually with decreasing pressure, but it remained above 10 L min^{-1} up to 20 km altitude. The entire system was adapted for balloon measurements and enclosed in a foam box to maintain the operating range of conditions during the balloon flight between temperatures of 10 and 40°C . The system was cali-

brated at the company before the flight and the performance of this system was also tested in a thermal vacuum chamber (TVAC) prior to the balloon flight. To the best of our knowledge, this system is used for the first time to measure ice particles in ice clouds from a balloon platform.

Estimation of cloud microphysical properties

Only data during the ascent flight are used for estimating the optical and microphysical properties of cirrus clouds. To estimate the microphysical (particle number concentration, effective diameter, and IWC) and optical properties (extinction coefficient and lidar ratio) from the Boulder Counter data, a lognormal size distribution function was fitted to the cumulative particle number concentration within the cloud layer. A single-mode lognormal size distribution function is expressed as

$$\frac{d}{dr} N(r) = N_0 \frac{\exp \left[-\frac{1}{2} \left\{ \frac{\ln^2 \left(\frac{r}{r_m} \right)}{\ln^2 \sigma} \right\} \right]}{r \ln \sigma \sqrt{2\pi}}, \quad (2)$$

where N is the number concentration, N_0 is the total particle concentration in the mode, r is the radius, r_m is the median radius, and σ is the width of the distribution (Thomason and Peter, 2006). A previous study (Kuhn and Heymsfield,

2016) has shown that a lognormal distribution provides a better match to the data than a gamma distribution. The first and the third moments of the size distribution are calculated from Eq. (2), which are then used to estimate the effective diameter (D_e) and IWC, respectively (described below).

The effective diameter of ice crystals is an important microphysical parameter used for the radiative calculations of ice clouds. It is defined as the ratio of absorption to the extinction cross-section of the ice crystals (Foot, 1988). Approximating ice crystals as spheres, D_e can be estimated using a lognormal particle size distribution from Eq. (3) as shown below:

$$D_e = 2 \times \frac{\int_0^r r^3 N(r) dr}{\int_0^r r^2 N(r) dr}, \quad (3)$$

where the numerator and the denominator are the third and second moments of the particle size distribution function, respectively. We also estimated the effective diameter using its temperature dependence relation given by Heymsfield et al. (2014) as shown below in Eq. (4) to compare with that derived from the size distribution:

$$D_e = \alpha e^{\beta T}, \quad (4)$$

where T is the temperature, $\alpha = 308.4$, $\beta = 0.0152$ for $-56^\circ\text{C} < T < 0^\circ\text{C}$. $\alpha = 9.1744 \times 10^4$, and $\beta = 0.177$ for $-71^\circ\text{C} < T < -56^\circ\text{C}$. $\alpha = 83.3$, and $\beta = 0.0184$ for $-85^\circ\text{C} < T < -71^\circ\text{C}$.

Using the ATTREX campaign data, Thornberry et al. (2017) have found that the effective diameter decreases more sharply at temperatures (T) below -81°C (192 K) than previously found by Heymsfield et al. (2014). This decrease in median effective diameter for $T < 192\text{ K}$ is expressed as

$$D_e = 12 + 28e^{0.625(T-192)}. \quad (5)$$

We also estimated the extinction coefficient from our particle size distribution measurements using a geometric optical approximation which assumes that the extinction coefficient is 2 times the total cross-sectional area of the ice crystals. Assuming that the ice particles are close to being spherical, we use the equation for extinction expressed in Thornberry et al. (2017):

$$\sigma_{\text{ext}} = 2 \times \sum_j N_j \pi r_j^2, \quad (6)$$

where N_j is the ice particle number concentration (L^{-1}) corresponding to the particle radius r_j (μm), with j ranging from 1 to 25 μm . The ratio of σ_{ext} (from Eq. 6) to the particulate backscatter coefficient (β_p) obtained from COBALD gives us the lidar ratio (discussed in Sect. 3.3.4). More details on the lidar ratio are presented in Appendix A2.

IWC is an important microphysical parameter for estimating the radiative impact of cirrus clouds as well as for understanding the microphysical and dynamical processes

occurring within them (Heymsfield et al., 2017). We estimated IWC for the tropopause cirrus independently from the COBALD backscatter and the Boulder Counter measurements. IWC was obtained from COBALD data using the extinction coefficient (σ_{ext}) and the effective diameter (D_e) estimated from Eq. (9e) given in Heymsfield et al. (2014) as shown below.

$$\text{IWC} = \sigma_{\text{ext}} \left(\frac{0.91}{3} \right) D_e \quad (7)$$

Here, D_e is obtained from Eq. (4) using temperature measurements. Using the ATTREX in situ measurements, Thornberry et al. (2017) have given a new parameterization relation for the estimation of IWC from the extinction coefficient in the temperature (T) range between 185 K (-88°C) and 192 K (-81°C) as shown below.

$$\text{IWC} = \frac{0.92}{3} \sigma_{\text{ext}} (12 + 28e^{0.65(T-192)}) \quad (8)$$

IWC is estimated by multiplying the ice density (917 kg m^{-3}) by the total volume obtained from the third moment of the particle size distribution. IWC values for the tropopause cirrus cloud layer derived from these abovementioned methods are compared in Sect. 3.3.5.

2.2.3 India Meteorological Department (IMD) radiosonde profiles from the University of Wyoming

We use daily (00:00 UTC) radiosonde profiles of meteorological parameters over the meteorological centre at the Hyderabad airport (17.45°N , 78.46°E) and its surrounding IMD stations such as Machilipatnam (16.2°N , 81.15°E), Visakhapatnam (17.7°N , 83.3°E), and Jagdalpur (19.08°N , 82.02°E), obtained from the University of Wyoming website of atmospheric soundings (<http://weather.uwyo.edu/upperair/sounding.html>, last access: 26 October 2024). Only profiles during the month of August 2017 are used in this study from these stations. The monthly mean temperature and wind profiles are constructed by using 27 d of available data over Hyderabad during August 2017. The monthly mean temperature and wind profiles are used to obtain anomalies from the profiles obtained on 23 August 2017. In order to obtain wave characteristics, profiles of zonal and meridional wind speed anomalies are used to perform hodograph analysis by following a method described in Leena et al. (2012). The results of hodograph analysis are presented in Sect. 3.4.3 and in the Supplement.

2.3 Satellite observations

2.3.1 Cloud–Aerosol Lidar with Orthogonal Polarization (CALIOP)

CALIOP was a dual-wavelength (532 and 1064 nm), dual-polarization, three-channel spaceborne lidar on board the

Cloud–Aerosol Lidar and Infrared Pathfinder Satellite Observations (CALIPSO) orbiting around the Earth in a sun-synchronous polar orbit at an altitude of about 705 km and an inclination of about 98.2° (Winker et al., 2009). It provided optical properties of aerosols and clouds distributed vertically in the Earth's atmosphere for 17 years starting in June 2006 at unprecedented spatial resolution with a repeat cycle of 16 d at a given location. CALIOP could detect thin and subvisible cirrus clouds (Martins et al., 2011) by measuring the attenuated backscatter coefficients (at two wavelengths and two polarizations at 532 nm), followed by the retrieval of 532 nm extinction coefficients (Young et al., 2018). IWC and effective diameter are estimated from these retrieved extinction coefficients, using a temperature-dependent empirical fit derived from aircraft data (Heymsfield et al., 2014). In this study, we used CALIOP Level 2 version 4.2 cloud profile (CPro) and cloud layer (CLay) data products having a horizontal resolution of 5 km along the orbit track and 60 m vertical resolution between the 8.2 and 20 km altitude range. We obtained scattering ratio profiles from the total backscattering coefficient profiles at 532 nm wavelength and molecular density profiles provided in the CPro data using the method described in the CALIOP Algorithm Theoretical Basis Document (ATBD; Young et al., 2008).

2.3.2 Cloud–Aerosol Transport System (CATS) lidar

CATS was a three-wavelength (355, 532, and 1064 nm), dual-polarization spaceborne lidar intended to provide the vertical distribution of aerosols and clouds from the International Space Station (ISS) platform (Yorks et al., 2014). It was launched in January 2015, operated until October 2017, and provided 33 months of aerosol and cloud data. The orbit of the ISS has an inclination of 51° and an altitude of about 405 km. This geometry provided more coverage over the tropics and mid-latitudes when compared to CALIPSO with a shorter repeat cycle of 3 d and observations at various local times. In this study, we used the 5 km horizontally averaged M7.2 Level 2, Version 3.0 Operational Layer (OL) and Profile (OP) products. These products are obtained from one of the two operational modes (M7.2) of CATS in which the backscatter coefficient and depolarization ratio at 1064 nm are measured. The CATS horizontal and vertical resolutions for the UTLS are the same as those of the CALIOP data. We obtained scattering ratios from the backscatter coefficient measurements at 1064 nm wavelength, combined with molecular density profiles provided in the data and using methods provided in the CATS ATBD (Yorks et al., 2015).

2.3.3 Cloud-top height from Himawari-8 brightness temperature

Himawari-8 is a geostationary satellite launched by the Japan Meteorological Agency in October 2014 to observe weather phenomena in 16 spectral bands (visible, near-infrared, and

infrared) at high resolution (Bessho et al., 2016). Cloud-top temperature observations from the 10.4 μm channel of the Himawari-8 satellite at an interval of 10 min have been used as a proxy for convection. Low cloud-top temperature indicates the presence of deep convective clouds and anvils. The horizontal resolution of cloud-top temperature is ~ 2 km at the satellite nadir for the infrared bands. The influence of deep convection on the tropopause cirrus clouds sampled by our balloon is investigated by using a technique combining back-trajectory analysis (described in Sect. 2.4.1) and cloud-top brightness temperature observations at 10.4 μm from Himawari-8 that was developed by Bedka and Khlopenkov (2016). This technique was used earlier in Vernier et al. (2018) to study the impact of deep convection, including overshooting convection, on air parcels sampled by balloons.

2.3.4 Humidity measurements from the Aura Microwave Limb Sounder (MLS)

Level 2, Version 4.2 relative humidity with respect to ice (RH_i) and water vapour mixing ratio (WVMR) data products (Lambert et al., 2015) from the Microwave Limb Sounder (MLS) instrument on board NASA's Aura satellite are used. These data products are available at 55 pressure levels between 1000 and 10^{-5} hPa. The vertical resolution of MLS RH_i ranges from 3.7 to 4.6 km between the 68 and 100 hPa levels with an accuracy of 20 % to 25 %. The vertical resolution of the MLS H₂O product is ~ 3 km with an accuracy of 8 %–9 % between the 68 and 100 hPa levels. Horizontal resolution along the orbit track is between 190 and 198 km at these pressure levels. The data screening criteria specified for RH_i and H₂O data by Livesey et al. (2020) have been applied to filter out the effect of clouds.

2.3.5 Temperature profiles from Global Navigation Satellite System Radio Occultation (GNSS-RO)

High-resolution Global Navigation Satellite System Radio Occultation (GNSS-RO) dry temperature profile products (Level 2, atmPrf) obtained from COSMIC-1, MetOp-A and MetOp-B, GRACE, KOMPSAT-5, TSX, and TDX missions are used to find the temperature near the tropopause for cirrus clouds detected by CALIOP and CATS. These data are provided by the Constellation Observing System for Meteorology, Ionosphere, and Climate (COSMIC) Data Analysis and Archive Center (CDAAC, <https://cdaac-www.cosmic.ucar.edu/cdaac/products.html>, last access: 26 October 2024) on 100 m altitude grids from 0 to 40 km altitude. Comparison of GNSS-RO temperature profiles with those from radiosondes has shown good agreement (Anthes et al., 2008). The GNSS-RO data set has found many applications in atmospheric science research such as in the study of gravity waves (Liou et al., 2003), the thermal structure of tropical cyclones (Biondi et al., 2013, 2015; Ravindra Babu et al., 2015), volcanic

cloud detection (Biondi et al., 2017), stratospheric thermal perturbation after the 2019/2020 Australian pyrocumulonimbus (pyroCb) event (Khaykin et al., 2020), and tropopause cold anomalies during the 2015 El Niño event over the Pacific Ocean (Ravindrababu et al., 2019).

2.4 Model simulations and reanalysis data

2.4.1 NASA Langley Trajectory Model (LaTM) simulations

The influence of deep convection on the tropopause cirrus cloud is investigated by using a combination of back-trajectory analysis and cloud-top temperature observations from the Himawari-8 geostationary satellite (described in Sect. 2.3.3). Meteorological fields from the NASA Global Modelling and Assimilation Office (GMAO) Goddard Earth Observing System, Version 5.2 Forward Processing (GEOS-5.2 FP; Lucchesi, 2017) reanalysis are used to compute three-dimensional back trajectories using the NASA Langley Trajectory Model (Fairlie et al., 2009, 2014). Back trajectories from the location of the balloon measurement are run at 100 m vertical resolution. Temperature and wind speeds along the back trajectories have also been obtained and used in this study to understand cloud formation in the air masses.

2.4.2 NOAA HYSPLIT trajectories

We used the National Oceanic and Atmospheric Administration (NOAA) Hybrid Single-Particle Lagrangian Integrated Trajectory (HYSPLIT) model (<https://www.ready.noaa.gov/HYSPLIT.php>, last access: 26 October 2024) to compute forward trajectories of air parcels from the cirrus cloud altitudes measured from our balloon payload.

2.4.3 ERA5 reanalysis data

We use hourly gridded ($0.25^\circ \times 0.25^\circ$) temperature, relative humidity, and wind speed from the ERA5 data product to understand the spatiotemporal variation of the meteorological variables and to supplement the satellite observations. ERA5 is the fifth generation of the European Centre for Medium-Range Weather Forecasts (ECMWF) reanalysis data product, which provides improved representations of the troposphere and tropical cyclones at high spatial and temporal resolution compared with its predecessor ERA-Interim (Hersbach et al., 2020). With reference to the high-resolution GNSS-RO temperature profiles, ERA5 temperature data provide the most realistic tropopause temperatures compared to other reanalyses (Tegtmeier et al., 2020). We derive the WVMR using the ERA5 temperature and relative humidity data. The ice saturation vapour pressure is obtained from the relation given by Murphy and Koop (2005). Over the ASM region, the WVMR obtained from ERA5 data provides a good representation of its vertical distribution and variation between

the 60 and 100 hPa pressure levels, with an average overestimation of 0.7–0.9 ppmv (15 %–17 %) compared to that observed by a cryogenic frost-point hygrometer (CFH) during the ASM (Brunamonti et al., 2019).

We list different sources of data and key parameters obtained from different instruments as discussed above in Table 1.

3 Results and discussion

3.1 Radiosonde, COBALD, and Boulder Counter observations

This section describes the measurements of clouds and background meteorological parameters from the balloon-borne instruments on 23 August 2017. The trajectory of this balloon flight and the temporal variation of the atmospheric conditions measured by the radiosonde are respectively shown in Fig. S1a and b in the Supplement. The vertical profiles of BSR₉₄₀, BSR₄₅₅, and CR obtained from COBALD during the ascent flight show several peaks (Fig. 2a). The peaks in the BSR₉₄₀ profile and the corresponding CR values close to unity clearly indicate the presence of five distinct cloud layers labelled as CL1, CL2, CL3, CL4, and CL5 in the free troposphere from bottom to top (Fig. 2a). Simultaneous measurements from the Boulder Counter shown in Fig. 2c reveal the prominence of particles larger than 5 µm (diameter) in these cloud layers. These particles could be water droplets and/or ice crystals depending on the air temperature, which is shown in Fig. 2b along with the wind speed. In what follows we describe the characteristics of each of these cloud layers from the lowest to the highest.

Specifically, CL1 is a thin mid-level cloud layer (Bourgeois et al., 2016) with its base located at about 4.1 km near the 0 °C isotherm and a geometrical thickness of about 0.7 km. It is also associated with a shallow temperature inversion present near its centre, which might be responsible for its formation as mentioned by Bourgeois et al. (2016). Since CL1 is located at a temperature near the 0 °C isotherm, we do not expect ice crystals in it. This layer is characterized by a low concentration (less than 10 L^{-1}) of liquid droplets smaller than 40 µm, with the majority being between 5 and 10 µm in size. CL2 is a cloud layer with optical thickness at 532 nm of ~ 0.37 and cloud-base and cloud-top altitudes at 8.8 and 11.3 km, respectively. It consists of a high concentration of particles spread over a wide size range. Since this cloud layer lies between the temperatures of -27 and -45 °C, it could be a mixed-phase cloud layer having both supercooled water droplets and ice crystals (Korolev et al., 2017). However, its cloud top is colder than -40 °C, indicating the presence of ice crystals which are smaller than 40 µm as shown in Fig. 2c. These ice crystals have likely grown larger at the expense of ambient supersaturation and consequently undergone sedimentation, as evidenced from the increased concentration of larger-sized (> 40 , 50, and

Table 1. List of different data sources and key parameters used in this study.

S. no.	Instrument/model/reanalysis data	Key parameters
Balloon-borne measurements		
1.	COBALD	Backscatter ratios at 455 and 940 nm and colour ratio
2.	iMet-1 radiosonde	Temperature (T), RH, u and v winds, GPS coordinates, and ascent rate
3.	SOLAIR Boulder Counter	Particle number concentration at six size channels, D_e , extinction coefficient, and IWC
4.	India Meteorological Department radiosonde profiles	T , RH, u and v winds
Satellite observations		
5.	CALIOP	Cloud-base and cloud-top height, backscatter ratio, depolarization ratio, colour ratio, D_e , extinction coefficient, and IWC
6.	CATS	Cloud-base and cloud-top height, backscatter ratio, and depolarization ratio
7.	Himawari-8	10.4 μm cloud-top temperature and anvil top height
8.	Aura MLS	Water vapour mixing ratio and RH _i
9.	GNSS-RO	Temperature profiles
Models/reanalysis		
10.	NASA LaTM	Back trajectories
11.	NOAA HYSPLIT	Forward trajectories
12.	ERA5 reanalysis	T , RH, water vapour mixing ratio, u and v winds

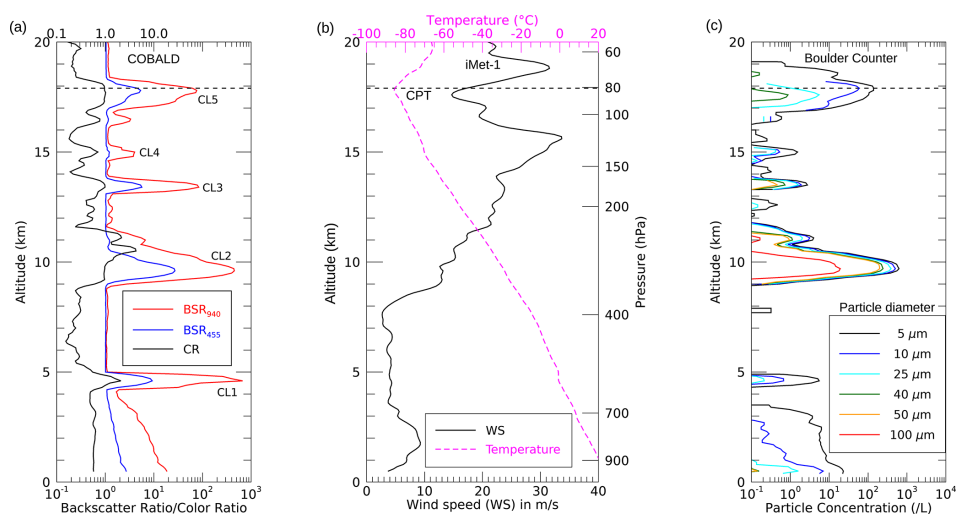


Figure 2. Vertical profile of (a) the backscatter ratio at 455 nm (blue line), 940 nm (red line) as well as the colour ratio (black line) obtained from COBALD. (b) Temperature (dashed magenta line) and wind speed (black line) from radiosonde, with the horizontal dashed black line showing the CPT and (c) cumulative particle number concentration for particles with a diameter greater than 5 (black line), 10 (blue line), 25 (cyan line), 40 (green line), 50 (yellow line), and 100 μm (red line) obtained from the Boulder Counter measurements on 23 August 2017 over Hyderabad, India, during the BATAL campaign. The vertical resolution is 100 m for each profile.

100 μm) ice crystals a few hundred metres below the cloud top and near the cloud base (Fig. 2c). The peak in BSR₉₄₀ almost coincides with the peaks of particle concentration corresponding to each size bin. The peak particle concentration (for diameter > 5 μm) in CL2 is found to be about 732.5 L^{-1} at an altitude of $\sim 9.6 \text{ km}$, with the majority of particles sized between 50 and 100 μm and a small number of particles larger than 100 μm . Above this layer, three layers (CL3, CL4, and CL5) of cirrus clouds (as the temperature is below -40°C) are observed with low ice crystal concentrations of less than 10 L^{-1} in CL3 and CL4 and less than 200 L^{-1} in

CL5. In CL3, ice crystals are smaller than 100 μm , and in CL4 they are smaller than 40 μm . CL5 is a subvisible cirrus cloud layer, with an optical thickness at 532 nm of ~ 0.025 , located at the CPT. The CPT is located at an altitude of about 17.9 km with a temperature of about -86.4°C (Fig. 2b). CL5 has a geometrical thickness of about 2.1 km with its top located at $\sim 18.3 \text{ km}$ in the lower stratosphere, $\sim 400 \text{ m}$ above the CPT. Note that the peak in BSR₉₄₀ and peaks in particle concentration corresponding to the 5 and 10 μm size bins are located exactly at the CPT. CL5 consists of particles smaller than 50 μm , with most of the particles smaller than 25 μm .

This tropopause cirrus is again detected by COBALD during the descent flight between 21:49 and 21:51 UTC at a similar altitude and similar temperature nearly 100 km south-west from the location of CL5 as shown in Fig. 3a. About 5 h later, a similar tropopause cirrus cloud layer was noticed in the CATS lidar observations as discussed in the next section. In the subsequent sections we discuss the properties of CL5.

3.2 CATS lidar observations of the tropopause cirrus cloud layer

On 24 August 2017 at about 01:36 UTC (~ 5 h after CL5 detection), there was an overpass of the ISS nearly 100 km west from the location of the tropopause cirrus cloud detected during the balloon descent flight and ~ 200 km from TIFR-BF. This allowed us to corroborate our balloon observations using the on-board operational CATS lidar data. Linking the data to their larger regional context, Fig. 3a shows a colour contour map of temperature from ERA5 reanalysis data at the 100 hPa pressure level (altitude ~ 16.8 km) averaged between 01:00 and 02:00 UTC on 24 August 2017. Both balloon (ascent and descent) and CATS measurements of clouds with base altitudes above 16 km at different times and different locations along the wind direction confirm the large horizontal extent of the tropopause cirrus cloud. The horizontal extent of this cloud layer along the CATS orbit track is found to be more than 500 km (Fig. S2 in the Supplement).

We further ran forward trajectories using the HYSPLIT model initialized at three different altitudes within the CL5 to see whether the cirrus cloud layer detected by CATS is influenced by the same air mass or not. These altitudes correspond to locations in the vicinity of the CL5 base altitude (16.8 km), the CPT (17.8 km), and the top altitude (18.3 km) of CL5. These trajectories are intersected by the CATS orbit track between 17.40 and 17.65° N after ~ 2.5 h of the passage of air parcels coming from the CL5, indicating their possible influence. The mean vertical profiles of the backscattering ratio and total depolarization ratio at 1064 nm from the CATS lidar averaged between 17.40 and 17.65° N along its orbit track are shown in Fig. 3b including the backscattering ratio profiles from COBALD at 940 nm during the ascent and descent. Although these measurements are at different times, locations, and wavelengths, good qualitative agreement can be seen between the two measurements as far as cirrus clouds near the tropopause are concerned. The value of the total depolarization ratio at 1064 nm from the CATS lidar between 0.4 and 0.5 confirms that this cloud layer consists of non-spherical ice crystals. The CATS cloud-phase algorithm also classifies this layer as an ice cloud layer (Fig. S2). Daily radiosonde observations at 00:00 UTC from the IMD Hyderabad station (17.45° N, 78.46° E) just before the ISS overpass on 24 August 2017 revealed that the CPT altitude was about 17.7 km (Fig. 3b), while the CPT temperature is found to be the coldest (about -81.5 °C) observed during the previous week. This cirrus cloud layer occurring at the tropopause is

the focus of this paper, and we discuss its observed microphysical (ice crystal number concentration, effective diameter, and ice water content) and optical properties (extinction coefficient and lidar ratio) derived from our balloon measurements in the following section (Sect. 3.3). We also investigate and discuss the mechanisms responsible for its formation in Sect. 3.4.

3.3 Microphysical and optical properties of the tropopause cirrus cloud

We only have four altitude bins within the CL5 at 100 m resolution where the estimation of microphysical properties using the size distribution is possible. This is because ice crystals were counted only in a few size bins away from the CPT as shown in Fig. 2c, and particle concentration in at least four size bins is needed for the solution to converge. Assuming spherical ice crystals in CL5, the lognormal size distributions for particle number and volume concentrations at each of these altitude bins are estimated. Figure 4 shows such distributions for the layer mean particle concentration in CL5. From the estimated size distribution, we derive various parameters as discussed below.

3.3.1 Ice crystal number concentration (N_{ice})

As shown in Fig. 4, the cumulative ice crystal number concentration for CL5 derived from the lognormal distribution (blue curve) fits the measured cumulative number concentration well (blue circles). The mean particle number concentration (for diameter > 5 μm) inside the CL5 layer is about 46.79 L^{-1} , with a maximum concentration of $\sim 196.3 \text{ L}^{-1}$ at the CPT (Fig. 2c). Particle concentrations at different altitude bins within CL5 are given in Table 2. These values are in good agreement with those measured inside tropopause cirrus clouds under similar temperatures observed over Bandung in Indonesia (Shibata et al., 2012), the western Pacific (Woods et al., 2018), and Costa Rica (Lawson et al., 2008).

3.3.2 Effective diameter (D_e)

Using Eq. (3), the effective diameters (D_e) of ice crystals within CL5 at different altitude bins from 17.3 to 17.7 km are estimated and are shown in Table 2. We have integrated between the radius from 1 to 25 μm because we do not see ice particles larger than 50 μm (diameter) in CL5. D_e varies between 16.34 and 26.61 μm with a mean plus or minus a standard deviation of $20.22 \pm 4.54 \mu\text{m}$. The mean effective diameter for CL5 between 17.3 and 17.7 km obtained from temperature using Eq. (4) is about $18.01 \pm 0.49 \mu\text{m}$, while the value obtained using Eq. (5) for temperatures between -81.32 and -84.94 °C is found to be about 21.29 μm . D_e estimated from Eq. (3) exhibits an increase in magnitude with decreasing temperature, which is contrary to what is observed in the case of D_e estimated from Eqs. (4) and (5).

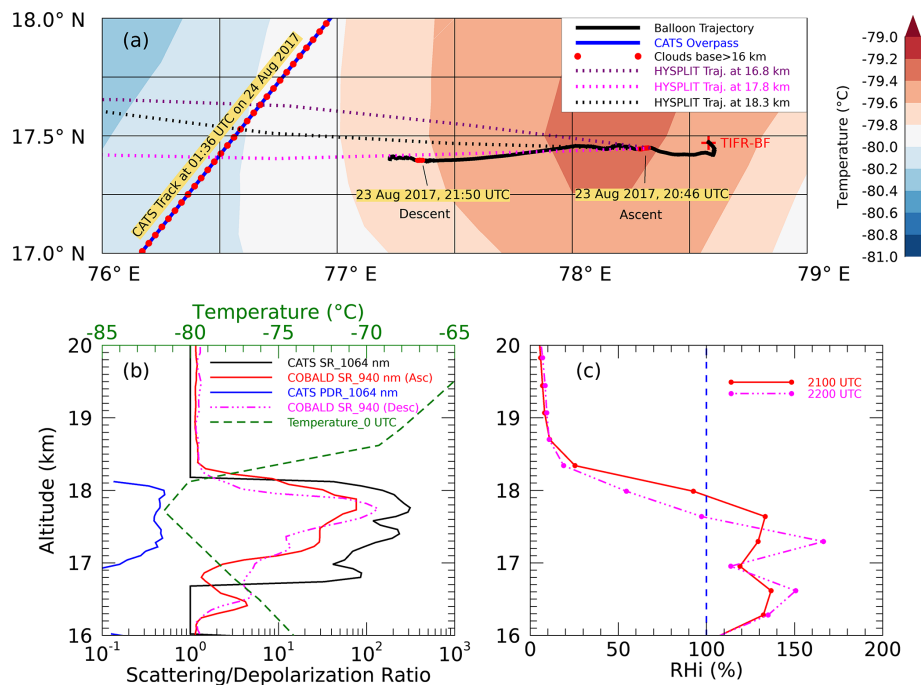


Figure 3. (a) ERA5 temperature at the 100 hPa level on 24 August 2017 averaged between 01:00 and 02:00 UTC. The CATS overpass is shown by a solid blue line with filled red circles superimposed on it representing the locations of cirrus clouds with a base height greater than 16 km. The solid black line shows the balloon flight trajectory from TIFR-BF (shown by a red plus sign) on 23 August 2017, with red circles representing the locations of cirrus clouds with a base altitude above 16 km during the ascent (at around 20:46 UTC) and descent (at around 21:50 UTC). Dashed lines show the forward trajectories from the location of tropopause cirrus at 16.8 km (purple), 17.8 km (magenta), and 18.3 km (black). (b) Vertical profiles of the scattering ratio (solid black line) obtained from CATS (at 1020 nm wavelength) and COBALD (at 940 nm wavelength) measurements during the ascent (solid red line) and descent (dashed–dotted magenta line) flight. The solid blue line represents the profile of the total depolarization ratio from CATS at 1064 nm. CATS profiles are averaged between 17.40 and 17.65° N along its orbit track. (c) Vertical profiles of relative humidity with respect to ice (RH_i) from ERA5 during the ascent (21:00 UTC) and descent (22:00 UTC). The vertical dashed blue line marks 100 % RH_i.

Table 2. Optical and microphysical properties of CL5 derived using COBALD and Boulder Counter measurements on 23 August 2017 from TIFR-BF. Data in italic font show the properties derived using only Boulder Counter data, while data shown in normal font are derived from COBALD and/or iMet-1 radiosonde data. Data shown in bold font are derived using both COBALD and Boulder Counter data.

Altitude (km)	17.3	17.4	17.5	17.6	17.7	Mean ± SD
Temperature (°C)	-81.32	-82.26	-83.23	-84.10	-84.94	-83.17 ± 1.43
No. concentration (#L ⁻¹)	54.23	67.27	99.47	101.77	108.87	85.32 ± 25.12
<i>D_e</i> (μm) using Eq. (3)	16.34	–	17.81	20.11	26.61	20.22 ± 4.54
<i>D_e</i> (μm) using Eq. (4)	18.62	18.36	18.0	17.68	17.41	18.01 ± 0.49
<i>D_e</i> (μm) using Eq. (5)	33.33	25.14	18.56	15.42	14.0	21.29 ± 7.98
σ_{ext} (km ⁻¹) from COBALD	0.013	0.018	0.017	0.018	0.026	0.018 ± 0.005
σ_{ext} (km ⁻¹) using Eq. (6)	0.014	–	0.029	0.032	0.035	0.028 ± 0.009
β_{p532} ($\times 10^{-4}$ km ⁻¹ sr ⁻¹)	5.77	8.13	7.71	8.33	12.06	8.40 ± 2.28
Lidar ratio (sr)	24.17	–	37.44	38.04	29.05	32.18 ± 6.73
IWC (mg m ⁻³) using Eq. (7)	0.072	0.099	0.093	0.098	0.140	0.100 ± 0.025
IWC (mg m ⁻³) using Eq. (8)	0.130	0.137	0.096	0.087	0.114	0.112 ± 0.021
IWC (mg m ⁻³) from the size distribution	0.070	–	0.157	0.195	0.285	0.176 ± 0.089

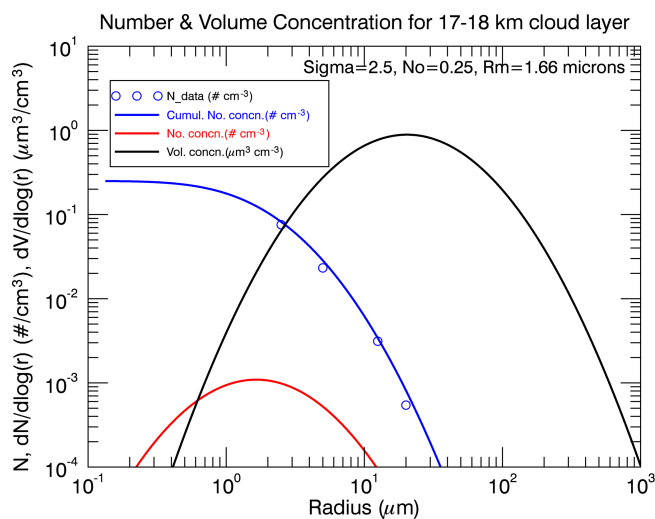


Figure 4. Particle size distribution in the cloud layer (CL5) observed between 17 and 18 km from the Boulder Counter measurements on 23 August 2017. Particle number concentration and volume concentration are shown by solid red and black lines, respectively. The measured cumulative number concentration is shown by open blue circles, while the solid blue line is a lognormal fit. The width of the distribution is given by sigma, which is about 2.5; total particle concentration (N_0) in the mode is 0.25 cm^{-3} , and $R_m = 1.66 \text{ }\mu\text{m}$ is the median radius.

An increase in D_e estimated from Eq. (3) is also consistent with an increase in particle concentration, with gradually increasing contributions coming from larger particles (diameter $> 25 \text{ }\mu\text{m}$) near the tropopause as shown in Table 2 and Fig. 2c. For comparison, the effective diameter of the tropopause cirrus cloud layer observed between 17 and 18 km at temperatures below $-83 \text{ }^\circ\text{C}$ by Shibata et al. (2007) was in the range from 8 to $80 \text{ }\mu\text{m}$. By contrast, ATTREX data showed that the median value of D_e reaches near $12 \text{ }\mu\text{m}$ at temperatures below $-86 \text{ }^\circ\text{C}$ (Thornberry et al., 2017).

3.3.3 Extinction coefficient (σ_{ext})

The values of the extinction coefficients estimated using relation Eq. (6) at different altitude bins within CL5 are given in Table 2. The mean and standard deviation of σ_{ext} ($0.028 \pm 0.009 \text{ km}^{-1}$) are slightly higher than the value of $0.018 \pm 0.005 \text{ km}^{-1}$ obtained from the FOV-corrected COBALD backscatter coefficient derived at 532 nm wavelength using a constant lidar ratio of 22 sr. This value of the lidar ratio is chosen based on the CALIOP L2 V4 algorithm following Young et al. (2018), where the initial lidar ratio at a cloud centroid temperature (for CL5) of about $-81.4 \text{ }^\circ\text{C}$ is about 22 sr. The lower value of the extinction coefficient derived from COBALD could be due to the choice of a constant lidar ratio.

3.3.4 Extinction to backscatter coefficient ratio (lidar ratio)

Table 2 shows the value of the lidar ratio inside CL5 at different altitude bins between 17.3 and 17.7 km with a mean and standard deviation of about $32.18 \pm 6.73 \text{ sr}$. In this study, the mean lidar ratio is estimated for a mean temperature of $\sim -83.2 \text{ }^\circ\text{C}$ (see Table 2). To the best of our knowledge, the in situ lidar ratio for cirrus clouds at this low temperature has not been estimated before. Our estimates of the lidar ratio are in good agreement with values ($40 \pm 10 \text{ sr}$) obtained by He et al. (2013) for cirrus clouds observed near the tropopause (temperature $\sim -80 \text{ }^\circ\text{C}$) over the Tibetan Plateau during July and August 2011. From this study, we demonstrate that the combined measurements from COBALD and the Boulder Counter can be used to estimate the range-resolved lidar ratio for optically thin cirrus cloud layers near the CPT.

3.3.5 Ice water content

The IWC from the size distribution information provided by the Boulder Counter data is estimated by multiplying the total ice volume by ice density (917 kg m^{-3}). By assuming an equivalent spherical volume for the non-spherical ice crystals, the total volume is estimated from their volume distribution calculated for ice crystal radii ranging between 1 and $25 \text{ }\mu\text{m}$. The values of IWC estimated at different altitude bins between 17.3 and 17.7 km using different methods are shown in Table 2. The mean and standard deviations of IWC estimated from the size distribution between 17.3 and 17.7 km are found to be about $0.176 \pm 0.089 \text{ mg m}^{-3}$. We also estimated the IWC from the FOV-corrected extinction coefficient obtained from COBALD data by using Eqs. (7) and (8) as shown in Table 2. Using Eqs. (7) and (8), the mean and standard deviations of IWC obtained from COBALD data are respectively found to be ~ 0.100 and 0.112 mg m^{-3} , which are in good agreement with those obtained from the size distribution data. When considering the entire CL5 layer, mean IWC obtained from COBALD data is found to be about 0.055 and 0.062 mg m^{-3} using the relation in Eqs. (7) and (8), respectively. These values are within the range ($0.001\text{--}10 \text{ mg m}^{-3}$) of IWC for cirrus clouds observed under similar temperature conditions (-90 and $-80 \text{ }^\circ\text{C}$) during several field campaigns over different regions of the globe (Lawson et al., 2008; Krämer et al., 2016; Heymsfield et al., 2014; Jensen et al., 2017; Woods et al., 2018; Thornberry et al., 2017). IWC (estimated by the parameterization Eq. 7) obtained from the CATS lidar observations between 17 and 18 km altitude, a few hours later than our measurements, is in the range $0.1\text{--}0.6 \text{ mg m}^{-3}$. COBALD and CATS measurements are of the same order of magnitude, but it is noted that these measurements are at different wavelengths, different sensitivity levels, different times, and different locations. Overall, our measurements are within the range of IWC measured for cirrus clouds in a similar temperature

range over other regions using aircraft-based instruments and spaceborne lidar. In the next section, we investigate the formation mechanisms of CL5.

3.4 Investigation of the mechanisms involved in the formation of the tropopause cirrus

3.4.1 Back trajectories and their intersection with deep convective clouds observed from Himawari-8

Overshooting convection can directly inject ice crystals into the lowermost stratosphere region (Corti et al., 2008; Khaykin et al., 2009, 2022; Dessler et al., 2016; Smith et al., 2017; Lee et al., 2019). Under the subsaturated conditions of the stratosphere, these ice crystals subsequently undergo sedimentation and sublimation, leading to a hydration patch (Jensen et al., 2020; Lee et al., 2019). This hydration patch could get advected and further lead to secondary ice formation upon cooling near the CPT region. Extreme convective clouds (cloud-top heights exceeding 17 km) can also induce cooling near the tropopause (Kim et al., 2018), which could eventually cause dehydration through ice formation. In addition to this, rapid tropopause cooling can also be caused by waves (Kim and Alexander, 2015; Kim et al., 2016), followed by supersaturation, cloud formation, and dehydration (Dzambo et al., 2019; Schoeberl et al., 2015, 2016). A recent study (Dzambo et al., 2019) over Australia has shown that the magnitude of tropopause cooling is greater during the monsoon period than that during the non-monsoon periods. CL5 could possibly be a result of the direct injection of ice crystals due to overshooting convection or through cooling caused by one of these processes mentioned above.

The influence of deep convection on CL5 is investigated by using a combination of back-trajectory analysis and cloud-top brightness temperature observations at $10.4\ \mu\text{m}$ from the Himawari-8 geostationary satellite data as discussed in Sect. 2.3.3. Figure 5a shows 5 d back trajectories initialized from the balloon measurement sites between 16 and 19 km as a function of latitude, longitude, and time. The locations of the convective anvil tops which intersected these air parcels are also superimposed along the back trajectories. A few isolated convective clouds over the Indian east coast along the back trajectories can be noticed in Fig. 5a. While intersecting the back trajectories, the anvil top of these convective events does not reach 18 km altitude. However, comparisons of GOES-16 (nearly identical to Himawari-8) and the Visible Infrared Imaging Radiometer Suite (VIIRS) temperature (Khlopenkov et al., 2021) show that pixel resolution has a strong impact on observed temperatures within convection. Therefore, the highest overshooting tops likely exceeded 18 km height. The convective clouds near the Indian east coast were also observed by the India Meteorological Department Doppler weather radars located at Machilipatnam (16.12°N , 81.09°E) after 12:00 UTC on 23 August 2017, with maximum echo altitude reaching about 17 km

(Fig. S3). These convective clouds were at their peak altitudes at around 11:40 UTC as seen in the Himawari-8 images, and we do not have radar images prior to 12:00 UTC or collocated CATS–CALIOP observations to confirm this. However, the intersection of air parcels with these convective clouds over the Indian east coast took place nearly 6 h later at around 16:00 UTC, when their anvil top altitude was reduced. Based on the Himawari-8 images and back trajectories, the direct injection of ice crystals by these local convective clouds up to 18 km altitude along the Indian east coast seems likely. Thus, we cannot fully rule out the influence of local convection because we notice enhanced WVMR at the 70 hPa pressure level on 23 August 2017 at 12:00 UTC over the surrounding regions of Hyderabad and the east coast (see Fig. S3). This enhanced water vapour might be associated with these convective clouds (Jensen et al., 2020; Ueyama et al., 2018). Enhanced water vapour might have led to secondary ice formation after enduring cold tropopause conditions, with subsequent advection towards Hyderabad by the TEJ. Ueyama et al. (2018) showed that the deep convective clouds occurring above the 380 K potential temperature level ($\sim 17\ \text{km}$) increase in situ formed cirrus clouds by 38 % compared to 21 % formation by convective clouds with cloud tops in the range 375–380 K. This could be the case for CL5.

In addition to hydration, these convective clouds could have also caused cooling near the CPT as discussed by Kim et al. (2018), which could have facilitated the in situ formation of ice crystals. Temperature measurements on 23 August 2017 at different times in and around the 400 km radius of TIFR–BF suggest that such cooling most likely took place before the arrival of the air parcels as shown in Table 3. Lack of accurate RH_i and ice measurements near the CPT over these convective clouds prevents us from confirming in situ formation. Thus, the role of local deep convective clouds in the formation of CL5 is not quite clear due to the lack of observations. High-resolution model simulations of these convective clouds could help ascertain their role in the formation of tropopause cirrus clouds, which is beyond the scope of this study.

In Fig. 5a, we note that the air parcels have already been influenced by deep convective clouds over the South China Sea before their intersection with local convective clouds over the Indian east coast. The altitude distribution of the back trajectories shown in Fig. 5b suggests that the air parcels initialized from the balloon measurement altitudes between 16 and 18 km originated from lower levels (below 10 km; not shown here) in the region between 115 and 125°E on 22 August 2017. This large-scale upward movement of the air parcels indicates a synoptic-scale convective system, such as a tropical cyclone or a typhoon. The Himawari-8 brightness temperature image on 22 August 2017 at 19:00 UTC in Fig. 6 does reveal such a synoptic-scale system named typhoon Hato, a category 3 tropical cyclone in the South China Sea, which severely affected the coastal cities of Macau, Zhuhai, and Hong Kong in southern China on 23 August

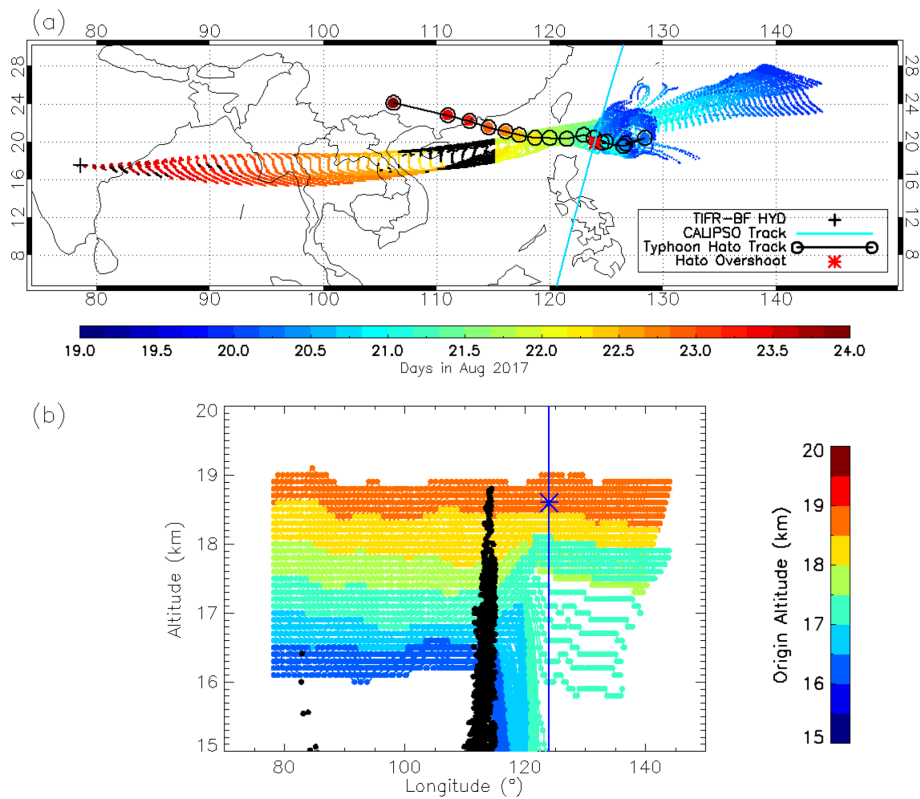


Figure 5. (a) 5 d back trajectories initialized on 23 August 2017 at 20:00 UTC from the balloon measurement locations at different altitudes between 16 and 19 km represented by coloured dots as a function of latitude and longitude, with colour showing the days in August 2017. The locations of deep convective anvil tops observed after 21 August 2017 from the Himawari-8 brightness temperature images at the $10.4\ \mu\text{m}$ wavelength channel that intersected the back trajectories are shown by black dots. The typhoon Hato track is shown by black-outlined coloured circles connected by a black line, with colour showing the days in August 2017. The nighttime CALIPSO orbit track on 20 August 2017 between 17:27 and 17:40 UTC is shown by a cyan line, with the location of the overshooting cloud top shown by a red asterisk. (b) The 5 d back trajectories are represented as a function of longitude and altitude, with colour showing their origin altitude. The locations of the anvil top altitude that intersected with the back trajectories are represented by black dots. The locations of the CALIPSO overpass and the overshooting cloud-top altitude are shown by vertical blue line and blue asterisk, respectively.

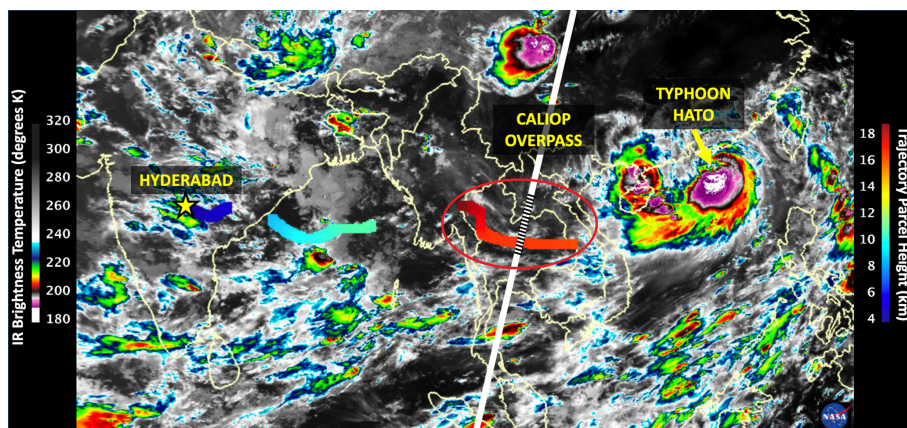
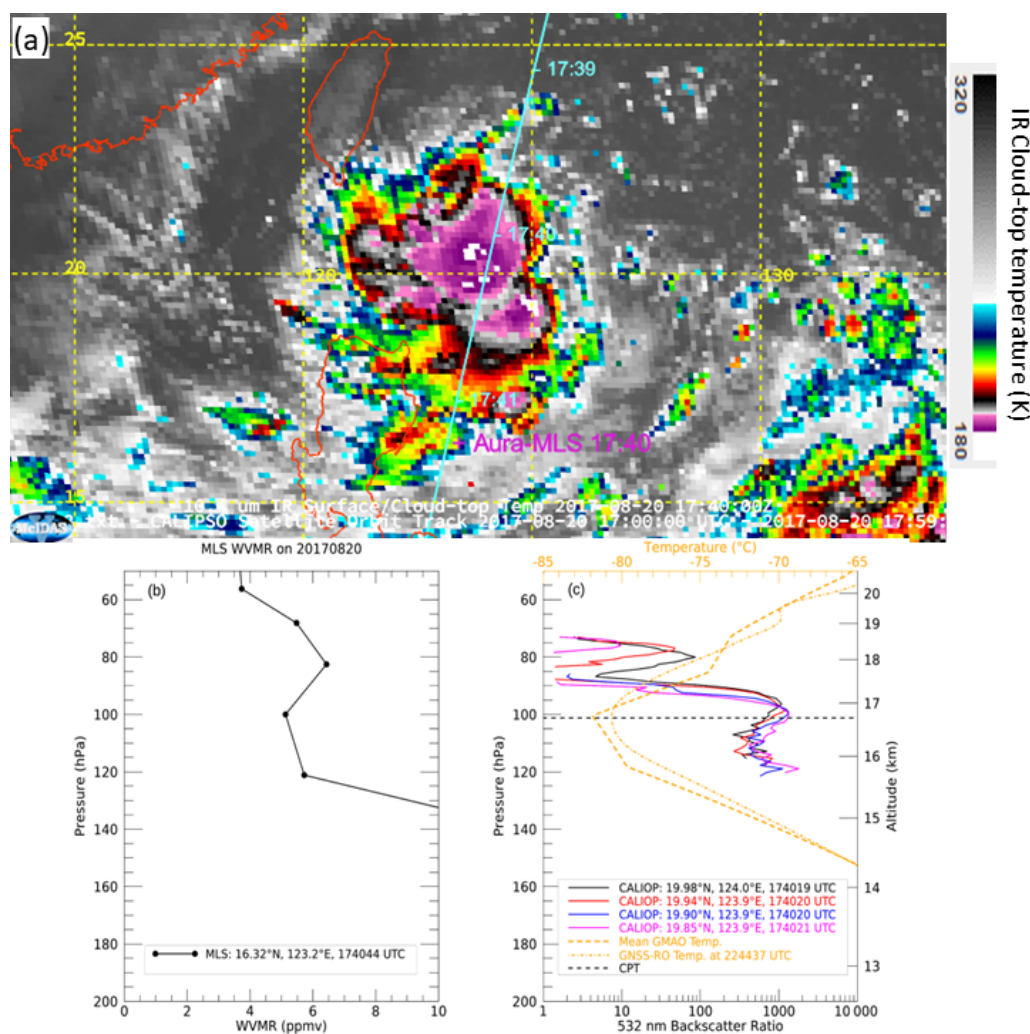


Figure 6. Image of brightness temperature (in Kelvin) at the $10.4\ \mu\text{m}$ infrared channel of the Himawari-8 satellite valid for 22 August 2017 at 19:00 UTC over East Asia. Typhoon Hato is associated with very low brightness temperature. The altitudes (km) of back-trajectory air masses initialized from the balloon measurement site (Hyderabad, India; denoted as a yellow star) are superimposed on the image and are represented by thick coloured curves. The air mass back trajectories initialized between 16 and 18 km altitude above Hyderabad at the time of the balloon flight on 23 August 2017 are encircled in red. The thick white line shows the nighttime CALIPSO overpass from 19:06 to 19:20 UTC on 22 August 2017, with black stripes showing the region where tropopause cirrus was sampled.

Table 3. Variation of the CPT temperature and altitude as measured by radiosondes and GNSS-RO in and around Hyderabad during 23 and 24 August 2017.

Date/time	23 August 2017, 00:00 UTC (radiosonde)	23 August 2017, ~ 14:00 UTC (GNSS-RO)	23 August 2017, 20:55 UTC (radiosonde)	23 August 2017, 21:50 UTC (radiosonde)	24 August 2017, 00:00 UTC (radiosonde)
Location	17.45° N, 78.46° E	18.45° N, 81.40° E	17.44° N, 78.30° E	17.40° N, 77.36° E	17.45° N, 78.46° E
Distance from TIFR-BF (km)	13.5	332	38.9	135.5	13.5
CPT altitude (km)	17.20	17.30	17.9	17.75	17.7
CPT temperature (°C)	−77.2	−82.5	−86.4	−86.1	−81.5

**Figure 7.** The (a) 10.4 μm infrared cloud-top temperature from Himawari-8 at 17:40 UTC on 20 August 2017 showing the overshoots during the development phase of typhoon Hato. This image is created with McIDAS-V software. The cyan line represents the CALIPSO orbit track that observed the overshoot on 20 August 2017 between 17:40 and 17:41 UTC, while the magenta plus symbol denotes the location of the Aura MLS profile at 17:40 UTC. (b) Vertical profile of WVMR from MLS on 20 August 2017 at 17:40 UTC near the overshoot region observed during the nighttime overpass of the Aura satellite. (c) Vertical profiles of the 532 nm backscatter ratio obtained from CALIOP on 20 August 2017 at around 17:40 UTC. Dashed orange lines show the mean temperature profile obtained from GMAO along the CALIPSO track and from GNSS-RO at 22:44 UTC. The horizontal dashed black line represents the CPT.

2017 after its landfall (Li et al., 2018). Typhoon Hato intensified to a category 3 typhoon just before its landfall (Pun et al., 2019). Hato was one of the strongest typhoons in the last several decades, causing widespread flooding in the Pearl River Delta region in southern China, with Macau being the hardest hit city (Li et al., 2018). The anvil cloud-top altitude derived from the Himawari-8 cloud-top temperature images along the back trajectories is found above 18.5 km (Fig. 5b), with the highest cloud-top altitude of about 18.8 km observed on 22 August 2017 at 09:10 UTC over 18.85° N, 114.48° E. It appears that typhoon Hato most likely injected ice crystals in the lower stratosphere, as the CPT altitude observed from the nearest radiosonde temperature profile over Haikou station (20.03° N, 110.35° E) at 12:00 UTC was about 17.4 km (Fig. S4 in the Supplement). Moreover, the tropopause temperature and tropopause altitude might be substantially reduced by the overshooting convection as pointed out by Sun et al. (2021).

Ice crystals injected by typhoon Hato in the subsaturated lower stratosphere might have caused a hydration patch through sedimentation and sublimation while being advected westward by the prevailing easterly winds. We did not have collocated CATS–CALIOP observations at that time to verify this overshooting event. However, there was a nighttime CALIPSO overpass on 20 August 2017 through the developing typhoon between 17:27 and 17:40 UTC (a few hours before the passage of the air parcels) that observed a cloud-top altitude of about 18.6 km (see Figs. 5a–b and S5). Daytime and collocated nighttime profiles of the WVMR from the MLS instrument confirm the presence of such hydration over typhoon Hato, with peak WVMR reaching ~ 6.5 ppmv at the 82.5 hPa pressure level near the tropopause (Fig. 7b). The location of the clouds observed by CALIOP is well above the CPT observed from the nearest GNSS-RO temperature profile at 22:44 UTC as shown in Fig. 7c. This WVMR value could have been even higher after Hato intensified to a category 3 typhoon because intensification can be followed by frequent long-lasting intense convective bursts (Horinouchi et al., 2020) that increase the possibility of overshooting and direct hydration in the lower stratosphere (Romps and Kuang, 2009; Jiang et al., 2020). Aircraft observations over the western Pacific region during NASA's POSIDON campaign have provided evidence for such direct hydration in the outskirts of typhoon Haima on 15 October 2016 (Jensen et al., 2020). They observed a layer (~ 1 km thick) of enhanced water vapour in the lower stratosphere with a peak WVMR of about 7 ppmv at an altitude of ~ 17.5 km, which could not be resolved by MLS owing to its coarse vertical resolution (see Fig. 2 of Jensen et al., 2020).

Jiang et al. (2020) studied the impact of 30 tropical cyclones (TCs, many of them observed during the ASM) on the UTLS water vapour over the tropical north-western Pacific Ocean using MLS and CloudSat observations during 2012–2016. They found that the lower-stratospheric water vapour over the TC area increased by an average value of

0.75 ppmv compared to the non-tropical cyclone area, indicating direct hydration. It seems that the hydration caused by typhoon Hato followed the anticyclonic flow and may have influenced the formation of CL5. We investigate this advection with further water vapour observations from MLS and ERA5-derived WVMR in the next subsection.

3.4.2 Advection of injected ice and hydration caused by typhoon Hato

As discussed in the previous section, the trajectories and the Himawari-8 cloud-top images suggest that the ice crystals detected by our balloon measurements are likely influenced by the convective outflow of typhoon Hato, which could possibly be advected to the measurement site by the TEJ or formed in situ from the moist air mass (from sublimated overshoot) brought by the TEJ near the cold tropopause. We discuss these aspects in this section. The TEJ is known to play a significant role in the redistribution of upper-tropospheric moisture and the formation of cirrus clouds during the ASM (Das et al., 2011). However, it can be questioned whether the ice crystals sampled by our balloon measurements could have survived that long distance (about 4000 km) if they were advected from the outflow of a typhoon because larger particles would eventually sediment and sublimate in a subsaturated region. Assuming the ice crystals to be spheroids (aspect ratio 1), the terminal velocity calculated by using the relationship given by Heymsfield and Westbrook (2010) suggests that both 5 and 10 μm particles will fall less than 1 km in 2 d. Thus, they are expected to survive within the layer if they were advected from the outflow of typhoon Hato by the TEJ provided they do not grow or sublimate and the background conditions remain the same. These theoretical calculations are consistent with the Boulder Counter measurements, which show the dominance of ice crystals smaller than 25 μm in CL5 (Fig. 2c). The presence of ice crystals larger than 10 μm in the middle of the CL5 could be explained by the constant growth, aggregation, and consequent sedimentation within the layer in the outflow.

A thin cirrus cloud layer with properties like CL5 can be seen near the tropopause between 17 and 18 km over Laos and Thailand in the nighttime observations from CALIOP on 22 August 2017 at around 19:00 UTC along the air parcels coming from the typhoon (Fig. 8a–d). This cirrus cloud layer extended horizontally to more than 200 km along the CALIPSO orbit track in the same latitude band (17–18° N) where the balloon measurements were taken. The values of the depolarization ratio between 0.3 and 0.6 and the colour ratio close to 1 (Fig. S6) in this layer confirm the presence of non-spherical ice crystals. The CALIOP IWC to extinction coefficient ratio provides an estimated effective diameter ranging between 18 and 20 μm , which is consistent with our balloon measurements ($20.22 \pm 4.54 \mu\text{m}$). The peak extinction coefficient value lies between 0.08 and 0.1 km^{-1} and is in the altitude range of 17.6–17.9 km, consistent with the

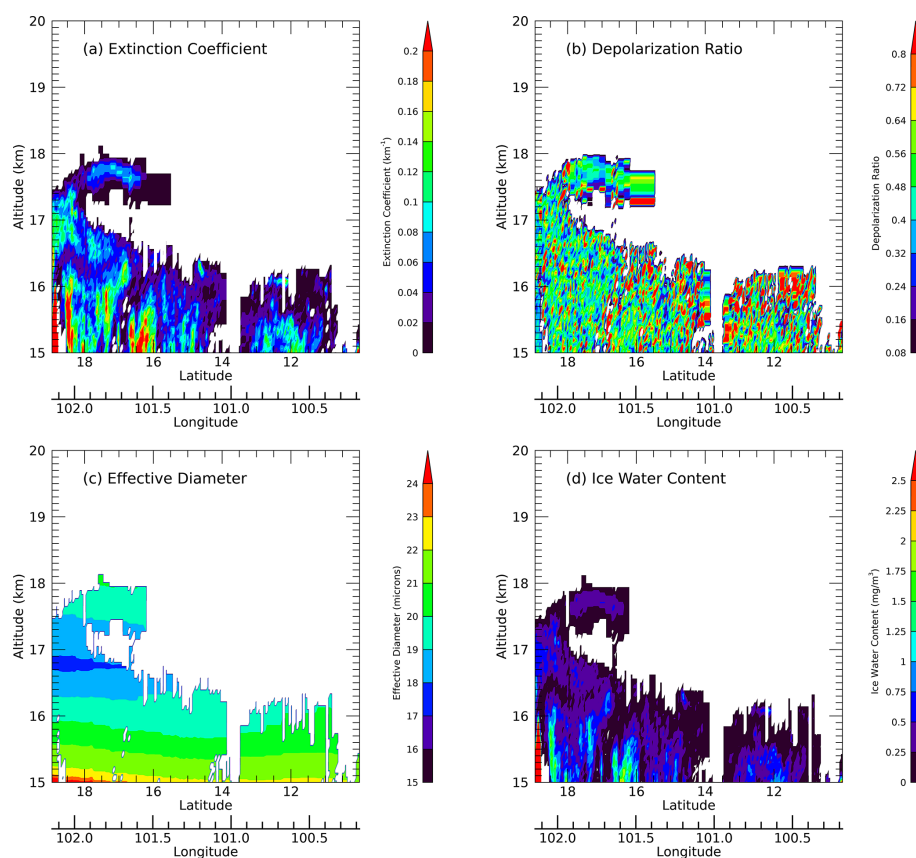


Figure 8. Optical and microphysical properties of cirrus clouds observed by CALIOP on 22 August 2017 at 19:00 UTC over Southeast Asia represented by black stripes in Fig. 6.

extinction coefficient maxima (0.06 km^{-1} at 17.9 km from COBALD) for CL5. IWC estimated by CALIOP for this layer using an empirical parameterization developed from aircraft data (Heysfield et al., 2014) ranges between 0 and 0.5 mg m^{-3} . These similarities in optical and microphysical properties suggest that the two layers detected by balloon and CALIOP could possibly be formed by the same mechanism under similar background conditions. However, the lifetime of ice crystals in cirrus clouds depends on several local environmental factors such as temperature, relative humidity with respect to ice, horizontal wind and updraft speed, and the type and quantity of ice nuclei. These factors may vary rapidly, especially in the TTL under the influence of deep convective clouds, TEJ, waves, and relatively high aerosol concentration during the ASM, which could influence the ice microphysics. This means that the ice crystals formed at one location through one mechanism may sublimate and/or crystallize at other locations while moist air is being transported to another location depending upon the local background conditions which force its distribution between ice and gas phases. It is therefore important to keep a track record of these parameters along the back trajectories. We discuss the role of environmental conditions such as temperature, water vapour,

and vertical wind speed in the in situ formation of tropopause cirrus clouds along the back trajectories in the next subsection.

As mentioned earlier, the convectively injected ice by typhoon Hato in the subsaturated lower stratosphere may undergo sedimentation and sublimation, leading to a hydration patch in the lower stratosphere (Lee et al., 2019). To investigate this, we derive WVMR from ERA5 relative humidity and temperature hourly data at the 70 hPa level over typhoon Hato on 22 August 2017 at 09:00 UTC when the peak overshoot was observed from the combined analysis of Himawari-8 and the back trajectories (Fig. 9a). We see an enhancement in the WVMR located between $15\text{--}20^\circ \text{ N}$ and $107\text{--}114^\circ \text{ E}$. As we move forward in time, this enhanced moist plume is seen to be advected towards the west. However, a weaker enhancement in WVMR can also be seen over the Indian east coast surrounding Hyderabad, probably due to afternoon deep convective clouds, which likely have also assisted the formation of CL5. The advection of a moist plume can also be seen in the averaged WVMR at the 70 hPa level for 22 and 23 August 2017 with high WVMR between 15 and 20° N latitude as shown in Fig. 9b.

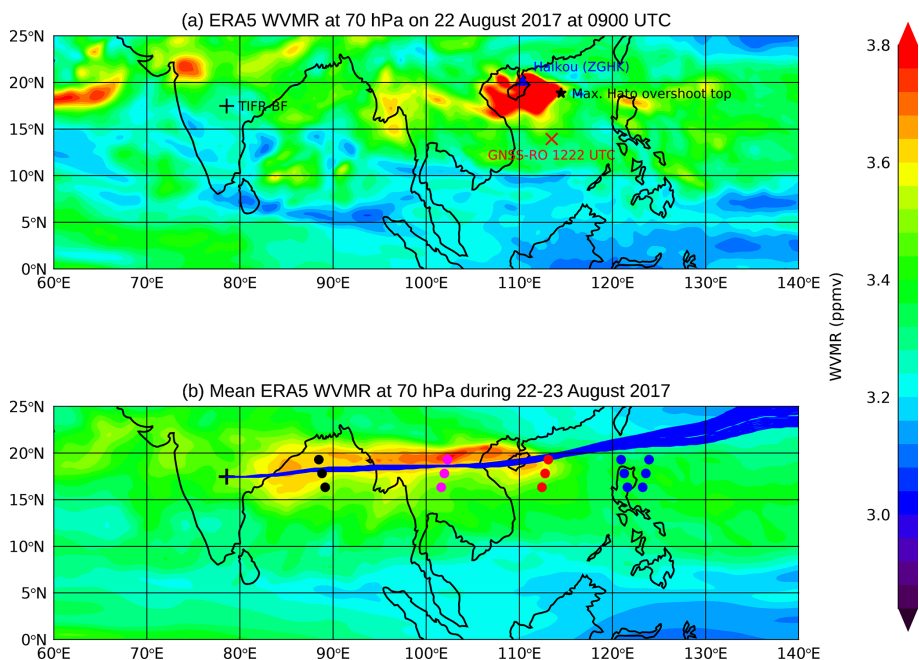


Figure 9. (a) Spatial map of WVMR derived from ERA5 data at the 70 hPa level on 22 August 2017 at 09:00 UTC showing large enhancement over the South China Sea. The locations of the TIFR Balloon Facility, Haikou radio sounding station, maximum overshoot top altitude, and GNSS-RO profile obtained at 12:22 UTC are represented by black plus, blue triangle, black star, and red cross symbols, respectively. (b) Average water vapour mixing ratio derived from ERA5 for 22 and 23 August 2017 at the 70 hPa level with back trajectories initialized between 18 and 19 km (between the 79 and 68 hPa levels) from the balloon site represented by blue lines superimposed on it. Nearest MLS footprints with respect to the back trajectories between 16 and 20° N on 20, 21, 22, and 23 August 2017 are represented by filled blue, red, magenta, and black circles, respectively.

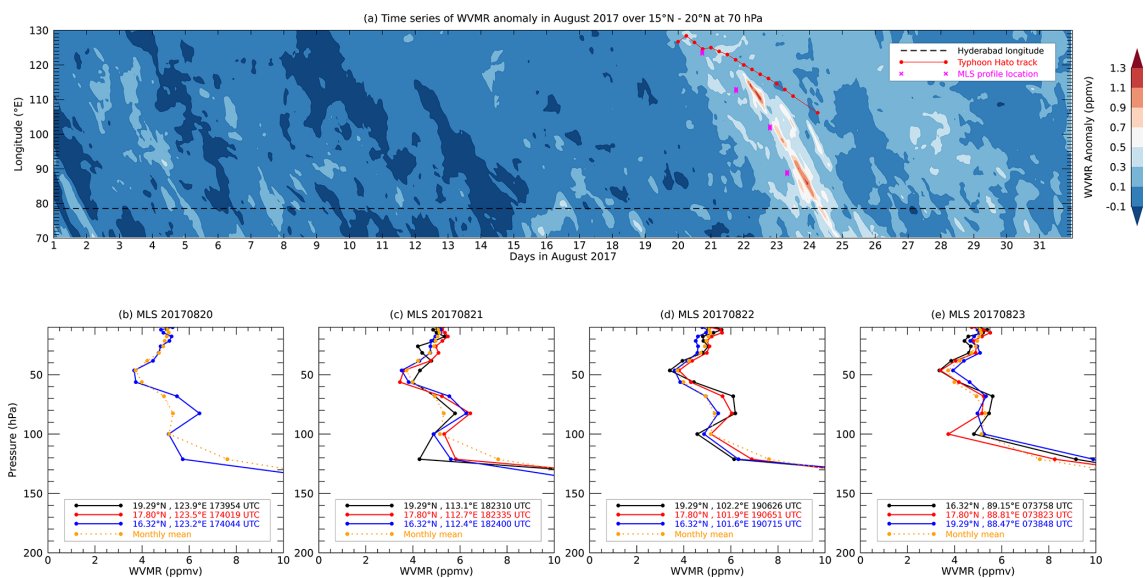


Figure 10. (a) Hovmöller plot for the anomaly in WVMR derived from ERA5 temperature and relative humidity at the 70 hPa pressure level for August 2017. The anomaly is computed by subtracting the monthly mean WVMR for each grid box between 15 and 20° N. The horizontal dashed black line marks the longitude of TIFR-BF in Hyderabad. Red dots represent the track of typhoon Hato, while magenta crosses show the location of the nearest MLS profiles on different dates. Vertical profiles of WVMR from MLS between 15 and 20° N observed on (b) 20 August 2017, (c) 21 August 2017, (d) 22 August 2017, and (e) 23 August 2017 at different times in UTC shown by different colours in the legend. The profiles affected by clouds are filtered out following MLS data screening criteria. The orange line with orange circles in each panel represents the mean MLS profile for August 2017 obtained by averaging all profiles located in the region bounded by 15–20° N and 70–130° E.

To quantify the transport of this enhanced moisture relative to the monthly mean WVMR at the 70 hPa pressure level and non-TC area, we estimated the WVMR anomaly for the month of August 2017 by subtracting the monthly mean WVMR for each grid box in the region between 15 and 20° N. The Hovmöller plot in Fig. 10a clearly shows the transport of enhanced moisture (0.6–0.9 ppmv) relative to the monthly mean reaching Hyderabad on 23–24 August during the night. We also notice that the magnitude of this enhanced moisture is in good agreement with the results found by Jiang et al. (2020) using MLS observations. We also verified this transport with the MLS WVMR observations during 20 and 23 August 2017 over this region between 15 and 20° N that transected the back trajectories nearest in space and time (Fig. 10b). We clearly see enhanced moisture (greater than the monthly mean value) at the 82.5 and 68 hPa pressure levels whose magnitude decreases as we move forward in time (from Fig. 10b–e). While being advected, this moist plume may get transported downward towards the CPT and may lead to in situ secondary ice formation upon cooling. We present the evidence for such downward transport and cooling near the tropopause along the back trajectories in the section below.

3.4.3 In situ ice formation due to cooling induced by waves associated with typhoon Hato

Figure 11 shows the temperature history of the air parcels along the back trajectories initialized from the balloon measurement site between 16 and 19 km altitude every hour. After being influenced by typhoon Hato, the air parcels have experienced several cooling and warming phases before reaching the balloon site. We observe a quasi-periodic occurrence of cold regions (represented by encircled regions in Fig. 11) with temperature below -81°C along the back trajectories after being influenced by typhoon Hato. Colder air parcels at these temperatures are susceptible to generating supersaturation that can trigger the formation of ice crystals through ice nucleation. Nearly coincident and collocated CATS and CALIOP observations on 22 and 23 August 2017, which intersected the back trajectories, are used to investigate the presence of cirrus clouds in these cold regions. CATS and CALIOP observations on 23 August 2017 are almost coincident with the air parcels along the back trajectories, whereas the CALIOP overpass on 22 August 2017 was a few hours later than the passage of the air parcels, and thus there is a difference in time (Fig. S7). It is interesting to see the presence of cirrus cloud layers with their base altitude greater than 16 km in and around these cold regions along the CALIPSO and CATS tracks.

The periodicity and the vertical extent of these cold regions can be clearly seen in Fig. 12a, where hourly temperature along the back trajectories initialized from the balloon site between 16 and 19 km altitude every 100 m is shown as a function of altitude (and potential temperature;

see Fig. S8). The periodic cold regions are confined in the altitude range between 16.5 and 18 km (potential temperature between 370 and 400 K; see Fig. S8) containing the CPT. The anvil top height derived from Himawari-8 cloud-top temperature, which intersected the back trajectories initialized between 16 and 19 km, is also superimposed (Fig. 12a). The coldest temperatures are found over the strongest updraft regions of typhoon Hato, with the highest cloud tops reaching ~ 18.5 km. The occurrence of cirrus clouds between 16 and 18 km altitude in these cold regions is also clear from the nearly collocated and coincident measurements from CATS and CALIOP (Fig. 12a). We observe westward-propagating wave-like patterns in the temperature of the air parcels after they are influenced by typhoon Hato. There is also an upward propagation of these wave patterns, which are more prominent near 18 km and altitudes above originating right from the region of typhoon Hato. Such periodic temperature fluctuations near the tropopause have also been observed along diabatic back trajectories (derived from ERA5 data) after being influenced by a typhoon during the ASM in a recent study (Li et al., 2020). Li et al. (2020) investigated the influence of two typhoons on the dehydration and transport of low-ozone air masses near the tropopause during the ASM over Kunming, China, using balloon measurements of temperature, ozone, and water vapour. However, such wave-like patterns in temperature and their role in cirrus cloud formation along the back trajectories are not discussed in their study. These wave-like temperature fluctuations or temperature anomalies near the tropopause could be due to the gravity waves generated by the deep convection associated with typhoon Hato.

We derived vertical wind speed along the back trajectories (Fig. 12b), which shows correspondence between the periodic updrafts and cold anomalies. This suggests that the areas of cold anomalies are most likely formed due to the cooling induced by the updraft. It has been found that the updraft motion slows down the process of sedimentation, leading to longer lifetime of ice crystals in cirrus clouds (Podglajen et al., 2018). The downdrafts along the back trajectories might have brought the moisture from the lower-stratospheric hydration patch towards the CPT, resulting in secondary ice formation. This is evident from Fig. 5b where air parcels that originated between 17 and 17.5 km can be seen descending from higher altitudes at around 120° N. One more interesting point to note is that the ascent speed of the balloon was increased from ~ 6 to $\sim 8\text{ m s}^{-1}$ near the tropopause region (17.5–18.5 km) with its peak value at the top of CL5, above which it decreases sharply (see Fig. S9), and beyond that it exhibits oscillatory behaviour with increasing amplitude. It is interesting to note the presence of downdrafts (Fig. 12b) and a sudden drop in balloon ascent speed above the CL5 top. Such vertical wind shear near the CPT could influence ice nucleation, with sudden rising motion triggering it while sinking motion inhibits it. The association of cold anomalies with updrafts indicates the possible role of wave influence, which we investigate below.

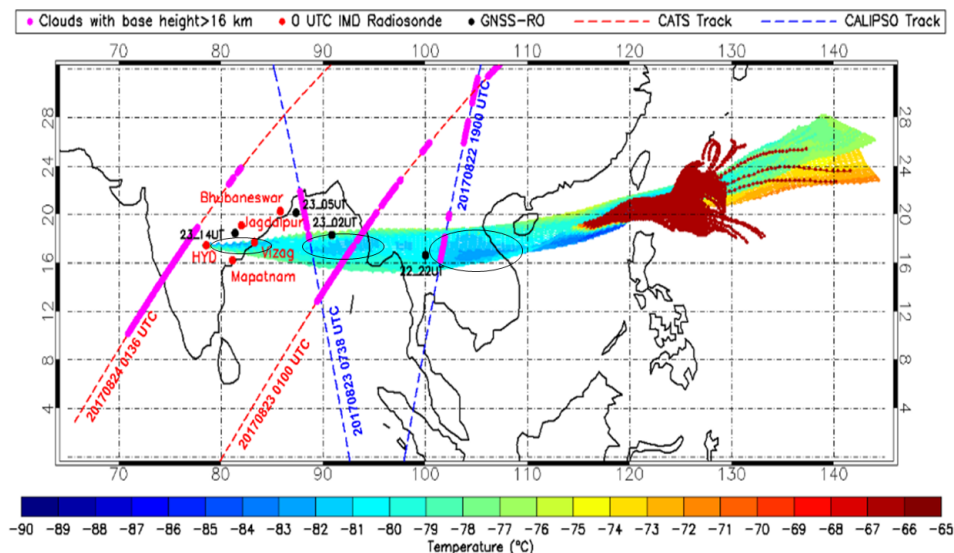


Figure 11. Filled coloured circles on the map show the temperature and location of the air parcel every hour along the back trajectories initialized from the balloon measurement sites between 16 and 19 km on 23 August 2017 at 20:00 UTC. The encircled regions along the back trajectories represent quasi-periodic features with temperature below -81°C . Dashed red and blue lines represent the orbit tracks of CATS (on 24 August 2017 at 01:36 UTC and 23 August 2017 at 01:00 UTC) and CALIPSO (23 August 2017 at 07:38 UTC and 22 August 2017 at 19:00 UTC), respectively, with respect to the air parcels. Filled magenta circles superimposed on the orbit tracks represent the locations of cirrus cloud layers with base altitude greater than 16 km. Filled red circles on the map show the locations of the IMD stations from which daily radiosondes are launched at 00:00 UTC. Filled black circles show the locations of available GNSS-RO temperature profiles closest in time and space to CATS and CALIPSO overpasses.

The temperature along the back trajectories in our analysis is obtained from the GEOS-FP reanalysis data, which may not be accurate enough to resolve these temperature fluctuations near the CPT when compared with the observations (Tegtmeier et al., 2020). To confirm the robustness of these cold anomalies, we use high-resolution and relatively more accurate temperature observations from GNSS-RO near these colder regions. We found temperature profiles within a 6 h interval and within a 400 km radius from the locations of the mean scattering ratio profiles of CATS and CALIOP intersecting the back trajectories between the 16 and 18° N latitude band as shown in Fig. 13b–d arranged according to their observation time (from latest to farthest). The temperature near the tropopause is better resolved in the GNSS-RO temperature profiles compared to the GMAO temperature profile averaged between the 16 and 18° N latitude band along the CATS and CALIOP orbit track. It is clear from the scattering ratio profiles obtained from COBALD (at 940 nm; Fig. 13b), CALIOP (at 532 nm; Fig. 13c and e), and CATS (at 1064 nm; Fig. 13d) that layered cirrus clouds or laminar cirrus are observed near the CPT in these cold anomalies. Profiles of RHi observed from the Aura MLS on 22 and 23 August 2017 collocated with CALIOP show RHi values greater than 100 % at the 100 hPa (~ 16.8 km) level, indicating supersaturation that resulted in ice formation (see Fig. S10).

The CPT temperature and CPT altitude over Hyderabad (within a 400 km radius around TIFR-BF) on 23 August 2017 exhibited a drastic change during the preceding 24 h as observed from balloon and GNSS-RO observations as discussed earlier (in Table 3). Deep convective clouds that occurred over land along the Indian east coast in the late evening hours (as mentioned in Sect. 3.4.1) might have helped strengthen this tropopause cooling. CPT temperature was at the minimum and CPT altitude was at the maximum for this BATAL flight. All the IMD stations near Hyderabad (Machilipatnam, Vizag, and Jagdalpur) recorded lower tropopause temperature on 24 August 2017 at 00:00 UTC compared to those observed on 23 August 2017 at 00:00 UTC. The magnitude of tropopause cooling can be expressed in terms of an anomaly in the observed temperature with respect to the mean observed temperature over Hyderabad. We constructed a mean temperature profile over the Hyderabad region during August 2017 using the temperature profiles obtained from all (nine) flights of the BATAL-2017 campaign over TIFR-BF and daily 00:00 UTC radiosonde data from the Hyderabad airport. The vertical profile of the observed temperature anomaly for our balloon flight on 23 August 2017, estimated after subtracting the mean temperature profile, clearly shows a cold anomaly of $\sim -8^{\circ}\text{C}$ near the CPT region (Fig. 14a) where CL5 was found. We also notice a wave-like pattern in the temperature anomaly profile above 15 km altitude. Using the airborne data from

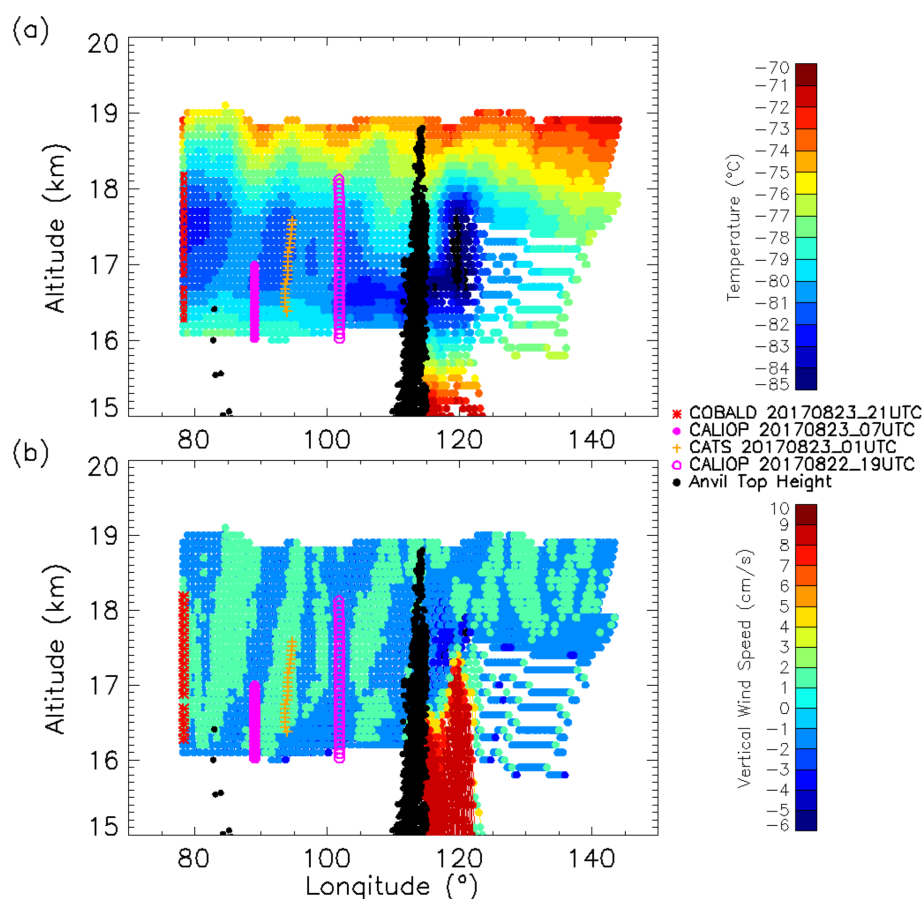


Figure 12. (a) Temperature history of the air parcels along the back trajectories initialized from the balloon measurement sites between 16 and 19 km on 23 August 2017. Each filled coloured circle represents a temperature value every 100 m of vertical resolution at each hour. (b) Same as panel (a) but for the vertical wind derived from the location and time of the back trajectories. Different coloured symbols represent the location of the layered cirrus clouds observed near the tropopause from CALIOP, CATS, and balloon measurements nearly collocated in space and time with respect to the back trajectories. The locations of anvil cloud-top altitude are represented by black dots.

the ATTREX campaign, Kim et al. (2016) showed that about 86 % of the cirrus clouds between 16.5 and 18 km over the Pacific are formed in cold anomalies induced by waves. They also found that these clouds more often form in the negative slope of the temperature anomaly, which is also true for CL5 (see the cloud shading in Fig. 14). However, a part of CL5 is also present in the positive anomaly above the tropopause up to the CL5 top altitude where ascent speed peaks and drops (see Fig. 14d). Similar wave-like patterns are also noticed in the anomaly profiles of zonal and meridional wind speed (Fig. 14b and c). They are estimated by subtracting the mean zonal and meridional speeds for August 2017 obtained from BATAL flights and Hyderabad airport radiosonde data for August 2017. Such a wave pattern is also reflected in the spatial movement of the balloon as shown in Fig. S1. In addition to this, the ascent speed of the balloon also shows oscillation above the CPT, with its amplitude increasing with altitude (Fig. 14d). All these observations indicate the possible influence of waves in cooling the tropopause. This may

have led to ice formation after reaching supersaturation with respect to ice. Recently, Reinares Martínez et al. (2021) reported a case of an extremely thin cirrus cloud layer formed near the CPT over the south-western Indian Ocean using COBALD and CFH observations. Using the spectral analysis (S-transform; Stockwell et al., 1996) of temperature anomaly and ascent speed, Reinares Martínez et al. (2021) suggested a role for homogeneous freezing under the influence of a high-frequency gravity wave with a vertical wavelength of 1.5 km. Using hodograph analysis of radiosonde data following Leena et al. (2012), we estimated the wave characteristics (see Fig. S11 and description). We found that the wave was propagating from the south-east direction to a north-west direction in the stratosphere with a horizontal phase speed of about 18.8 m s^{-1} and a horizontal wavelength of about 1770 km, indicating that it was generated from typhoon Hato.

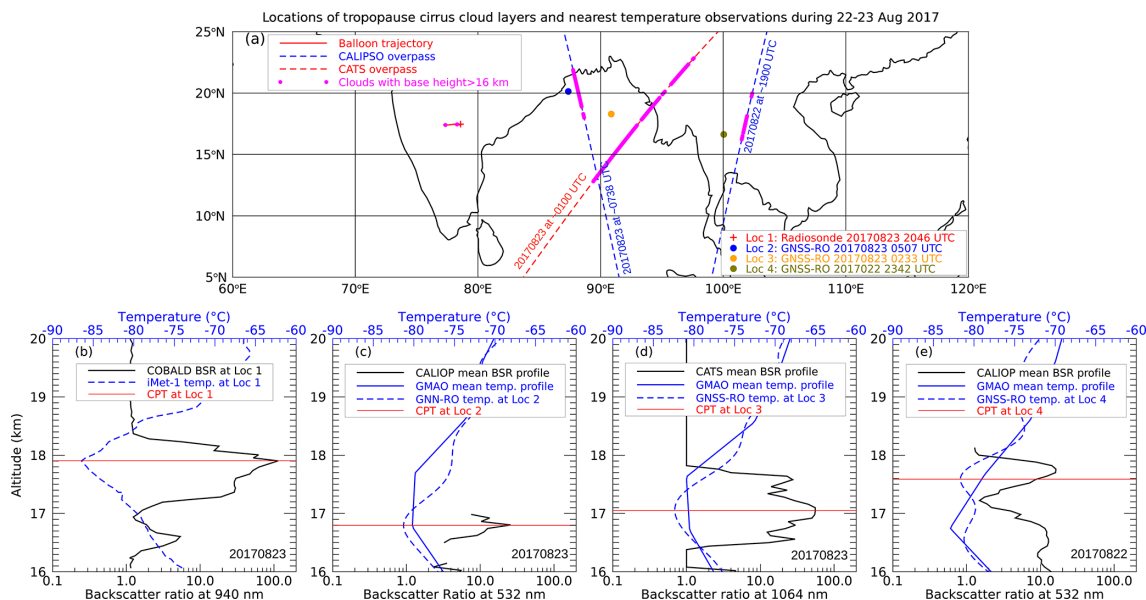


Figure 13. (a) Map showing the locations of tropopause cirrus clouds (filled magenta circles) observed along the balloon trajectory (solid red line) as well as during CATS (dashed red line) and CALIOP (dashed blue lines) overpasses on 22 and 23 August 2017. Vertical profile of the backscatter ratio (shown by black line) obtained from (b) COBALD data on 23 August 2017 at around 20:40 UTC, (c) CALIOP on 23 August 2017 at around 07:38 UTC, (d) CATS on 23 August 2017 at around 00:59 UTC, and (e) CALIOP on 22 August 2017 at around 19:07 UTC. CALIOP and CATS profiles are averaged between 17 and 18° N. The blue line in panels (c), (d), and (e) represents the mean temperature profile obtained by averaging the corresponding GMAO GEOS-FP temperature profiles provided with CALIOP and CATS data. The dashed blue line shows the nearest available coincident temperature profile obtained from (b) radiosonde measurement over location 1 (red plus symbol), (c) COSMIC GPS RO on 23 August 2017 at 05:07 UTC over location 2 (filled blue circle), (d) COSMIC GPS RO on 23 August 2017 at 02:33 UTC over location 3 (filled orange circle), and (e) COSMIC GPS RO on 22 August 2017 at 23:42 UTC over location 4 (filled olive circle). The horizontal red line in each panel marks the CPT.

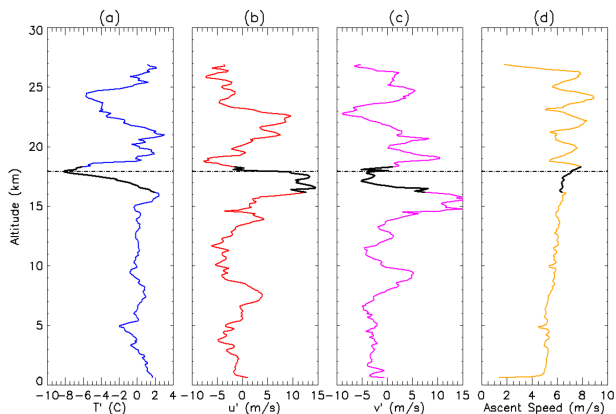


Figure 14. Vertical profile of (a) temperature anomaly (T'), (b) zonal wind speed anomaly (u'), (c) meridional wind speed anomaly (v'), and (d) ascent speed of the balloon observed on 23 August 2017. Anomalies are estimated by subtracting the monthly mean profile from the profile obtained on 23 August 2017. Mean profiles are obtained from all the radiosonde profiles at 00:00 UTC and BATAL flights over Hyderabad for August 2017. The horizontal dashed-dotted black line in each panel denotes the location of the CPT (at 17.9 km) observed on 23 August 2017 at around 21:00 UTC from the balloon. Black shading in each profile shows the location of the tropopause cirrus cloud.

4 Summary

In situ measurements in cirrus clouds with ice crystals smaller than $100\ \mu\text{m}$ near the tropopause over the ASM region are sparse, and the available observations mainly come from probes on aircraft that are travelling at rapid speeds through the clouds, leading to shattering of ice crystals and failure to resolve small ice crystal habits. In this paper, we have presented balloon-borne measurements of the optical and microphysical properties of a tropical tropopause cirrus cloud layer obtained using a backscatter sonde and an optical particle counter on 23 August 2017 during the BATAL campaign over Hyderabad, India. This layer was a sub-visible cirrus cloud located at an extremely cold tropopause temperature of $-86.4\ ^\circ\text{C}$. The top of this layer was found in the lower stratosphere at about 18.3 km. Ice crystals in this cloud layer were smaller than $50\ \mu\text{m}$ in diameter. Nearby spaceborne lidar backscatter measurements from the CATS on board ISS confirmed the presence of this tropopause cirrus that extended more than 500 km along the ISS orbit track. IWC for this layer was estimated independently from the backscatter measurements and optical particle counter and found to be less than $0.2\ \text{mg}\ \text{m}^{-3}$. Simultaneous measurements of the backscatter coefficient from COBALD and the

extinction coefficient from a Boulder Counter allowed us to derive a range-resolved lidar ratio for this tropopause cirrus cloud layer, with a layer average value of 32.18 ± 6.73 sr. We demonstrated that the combination of COBALD and the Boulder Counter can be used to estimate a range-resolved lidar ratio for tropopause cirrus cloud layers, which may prove useful in validating lidar retrievals of the extinction coefficient. Our estimate can be improved with a higher ice particle sampling rate at more size channels, especially at smaller radii. We also investigated the formation mechanism of this layer using back trajectories, satellite, and ERA5 re-analysis data. Back trajectories from the balloon measurement site and their intersection with convective clouds observed in Himawari-8 brightness temperature images suggest that the layer was influenced by a category 3 typhoon named Hato. Both Himawari-8 and CALIOP observations showed that Hato injected ice crystals into the lower stratosphere, causing a hydration patch as revealed by ERA5 WVMR and confirmed by the MLS WVMR observations. These moist plumes were seen to be advected by the ASM anticyclonic flow towards Hyderabad. Moreover, along the back trajectories there were quasi-periodic cold and warm anomalies near the CPT associated with updrafts and downdrafts, respectively. CATS and CALIOP observations showed the presence of tropopause cirrus clouds in these cold anomalies. These perturbations are most likely caused by the waves induced by typhoon Hato. Signatures of such perturbations were also noticed in the temperature and wind profiles obtained from our balloon observations. Through this case study, we conclude that the overshooting clouds in typhoons can cause direct stratospheric hydration during the ASM in addition to the usual overshooting convective systems. Following the ASM anticyclonic flow, this stratospheric hydration can get advected several thousand kilometres and may subsequently lead to tropopause cirrus clouds upon cooling caused by typhoon induced waves.

The occurrence frequency of tropopause-penetrating deep convective clouds (Aumann et al., 2018) and the intensity of tropical cyclones are expected to increase in a warmer climate (Stocker et al., 2013; Emanuel, 2005), which in turn are likely to increase the occurrence of ice injections with consequences for the stratospheric composition, thin tropopause cirrus clouds, and further feedbacks on the global climate (Dessler et al., 2016; Solomon et al., 2010). A recent modelling study (Smith et al., 2022) found that the impact of convective hydration in response to increased CO₂ depends not only on the frequency and penetration altitude of convective overshooting into the stratosphere but also on large-scale temperatures in the TTL. In this context, the occurrence of frequent deep convection and presence of cold temperatures in the TTL over the ASM anticyclone region are important for convective hydration and tropopause cirrus cloud formation. In future, simultaneous measurements of temperature, water vapour, and the microphysical properties of cirrus clouds using quasi-isentropic balloon flights within the

coldest regions of the ASM anticyclone will be planned to obtain a detailed understanding of the impacts of overshooting convection on the formation of tropical tropopause cirrus clouds.

Appendix A: Data and methods

A1 Compact Optical Backscatter Aerosol Detector (COBALD)

COBALD consists of two high-power (500 mW) light-emitting diodes (LEDs) which emit light at 455 nm (blue) and 940 nm (near-infrared) wavelengths. A silicon photodiode placed between the two LEDs collects the light backscattered from the airborne particles (molecules, aerosols, and clouds). When connected with an iMet radiosonde, the instrument monitors backscatter signals at both wavelengths and adds backscatter to the radiosonde pressure, temperature, relative humidity, wind speed, and wind direction data at a rate of 1 Hz to the receiving ground station. The concurrent profiles of pressure, temperature, and backscattered counts are used to calculate the backscatter ratio (BSR), the ratio of the total (molecules and particulate) backscatter coefficient ($\beta = \beta_p + \beta_m$) to the molecular backscatter coefficient (β_m) at a given wavelength. The BSR is calibrated in a region of the atmosphere minimally affected by aerosols and adjusted with a reference density signal derived from pressure and temperature. The absolute error associated with BSR is within 5 %, while its precision is better than 1 % in the UTLS region (Vernier et al., 2015).

A2 Extinction to backscatter coefficient ratio (lidar ratio)

The ratio of extinction to backscatter coefficient, often called the lidar ratio, is used in retrieving the extinction coefficient from the backscatter measurements of elastic backscatter lidars. The retrieved extinction coefficient and hence the layer optical thickness are very sensitive to the magnitude of the lidar ratio used in the extinction retrieval, which depends on the number concentration, size and shape distributions, and refractive index of the particles (Ackermann, 1998). Since microphysical properties of cirrus clouds mainly depend on temperature and RH_i (Heymsfield et al., 2017), it is essential to study the dependence of the lidar ratio on these factors. The variation of the lidar ratio within a cirrus cloud layer can be measured in situ or by using inelastic lidars, such as Raman lidars (Ansmann et al., 1992; Sakai et al., 2003) or high-spectral-resolution lidars (Grund and Eloranta, 1990). In the case of elastic backscatter lidars, for an extinction coefficient retrieval the lidar ratio is usually assumed to be constant (range-independent) inside a cirrus cloud layer. Version 3 CATS lidar cloud products use a constant lidar ratio of 28 sr over land and 32 sr over the oceans for tropical ice clouds (CATS L2O Layer Products Quality Statements Version Release: 3.00, <https://cats.gsfc.nasa.gov/media/docs/>

CATS_QS_L2O_Layer_3.00.pdf, last access: 12 December 2024). This approach of using a constant lidar ratio is unable to account for the natural variability of the actual lidar ratio, leading to erroneous extinction coefficient and cloud optical thickness retrievals for cirrus clouds (Saito et al., 2017). Using multiyear nighttime CALIOP two-way transmission and collocated IIR absorption optical depth at 12.05 μm for single-layer semi-transparent cirrus cloud layers over the ocean, Garnier et al. (2015) derived a temperature-dependent multiple-scattering factor and a corresponding temperature-dependent parameterization of the initial lidar ratio. Following this study, the multiple-scattering factor and lidar ratio in CALIOP V4 algorithm are approximated by a sigmoid function of the centroid temperature of the 532 nm attenuated backscatter coefficient within the cloud layer (Young et al., 2018). According to this sigmoid function, the value of initial lidar ratio for a cirrus cloud layer with centroid temperature between -80 and -90 $^{\circ}\text{C}$ falls in the range of 20–22 sr. The introduction of this approach in the CALIOP V4 algorithm has resulted in the reduction of the relative uncertainty assigned to the initial lidar ratio for semi-transparent ice clouds by 15 % compared to the V3 algorithm (Young et al., 2018).

There are many studies on the lidar ratio of cirrus clouds using lidars measurements (see Table 5 of Voudouri et al., 2020), but such studies using in situ measurements of subvisible cirrus clouds near the tropical tropopause are rare. Our unique balloon measurements provided us with independent in situ measurements of backscatter (from COBALD) and extinction coefficients (from a Boulder Counter) inside a subvisible cirrus layer near the CPT, which allowed us to estimate the vertical profile of the lidar ratio as shown in Table 2.

Data availability. CALIPSO data were obtained from the NASA Langley Research Centre Atmospheric Science Data Center at https://doi.org/10.5067/CALIOP/CALIPSO/LID_L2_05KMCLAY-STANDARD-V4-20 (NASA/LARC/SD/ASDC, 2018a) and https://doi.org/10.5067/CALIOP/CALIPSO/LID_L2_05KMCPRO-STANDARD-V4-20 (NASA/LARC/SD/ASDC, 2018b). GEOS-5 FP wind data are from NASA GMAO, and Himawari-8 data are from the University of Wisconsin–Madison Space Science and Engineering Center (SSEC). Doppler weather radar data were obtained from the India Meteorological Department (IMD). Daily 00:00 UTC radiosonde data were obtained from the University of Wyoming (<https://weather.uwyo.edu/upperair/sounding.html>, University of Wyoming, 2024), GNSS-RO temperature profile data from UCAR-COSMIC (<https://cdaac-www.cosmic.ucar.edu/>, UCAR-CDAAC, 2024), ERA5 reanalysis data from the Copernicus Climate Data Store at <https://doi.org/10.24381/cds.bd0915c6> (Hersbach et al., 2023), Aura MLS data from NASA Earthdata at <https://doi.org/10.5067/Aura/MLS/DATA2019> (Read and Livesey, 2015) and <https://doi.org/10.5067/Aura/MLS/DATA2009> (Lambert et al., 2015), and CATS data from NASA-GSFC (<https://cats.gsfc.nasa.gov/>, last access: 28 October 2024). The

typhoon Hato track was obtained from the Japan Meteorological Agency (JMA) Best Track Data at <https://www.jma.go.jp/jma/jma-eng/jma-center/rsmc-hp-pub-cg/besttrack.html> (JMA, 2024). Balloon data used in this study are available at <https://doi.org/10.6084/m9.figshare.27932037.v1> (Li, 2018a). Trajectory data used in this study are available at <https://doi.org/10.6084/m9.figshare.27932082.v1> (Li, 2018b).

Supplement. Please note that the supplement link is a placeholder that will be substituted during the publication. The supplement related to this article is available online at: <https://doi.org/10.5194/acp-24-14209-2024-supplement>.

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