



## Evaluating F2-region long-term trends using the International Reference Ionosphere (IRI) model: is this a feasible approximation for experimental trends?

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**Abstract.** The International Reference Ionosphere (IRI) is a widely used empirical ionospheric model based on observations from a worldwide network of ionospheric stations. Therefore, it would be reasonable to expect it to capture long-term changes in key ionospheric parameters, such as  $foF2$  and  $hmF2$  linked to trend forcings like greenhouse gas increasing concentrations and the Earth's magnetic field secular variation. Despite the numerous reported trends in  $foF2$  and  $hmF2$  derived from experimental data and model results, there are inconsistencies that require continuous refinement of trend estimation methods and regular data updates. This ongoing effort is crucial to address the difficulties posed by the weak signal-to-noise ratio characteristic of ionospheric long-term trends. Furthermore, the experimental verification of these trends remains challenging, primarily due to time and spatial coverage limitations of measured data series. Achieving these needs for accurate detection of long-term trends requires extensive global coverage and high resolution of ionospheric measurements together with long enough periods spanning multiple solar cycles to properly filter out variations of shorter terms than the sought trend. Considering these challenges, IRI-modeled  $foF2$  and  $hmF2$  parameters offer a valuable alternative for assessing trends and obtaining a first approximation of a plausible global picture representative of experimental trends. This work presents these global trend patterns, considering the period 1960–2022 using the IRI-Plas 2020 version, which are consistent with other model predictions. While IRI explicitly takes into account the Earth's magnetic field variations, the increase in the concentrations of greenhouse gases appears indirectly through the Ionospheric Global index (IG) which is derived from ionospheric measurements. F2-region trends induced by the first mechanism should be important only around the magnetic equator at the longitudinal range with the strongest displacement, and it should be negligible out of this region. Conversely, trends induced by the greenhouse effect, which are the controversial ones, should be dominant away from the geomagnetic equator and should globally average to negative values in both cases, i.e.,  $foF2$  and  $hmF2$ . Effectively, these negative global means are verified by trends based on IRI-Plas, even though not for the correct reasons in the  $hmF2$  case. In addition, a verification was performed for more localized  $foF2$  trend values, considering data from nine mid-latitude stations, and a reasonable level of agreement was observed. It is concluded that the IRI model can be a valuable tool for obtaining preliminary approximations of the Earth's magnetic-field-induced long-term changes in  $foF2$  and  $hmF2$ , as well as of experimental trends only in the  $foF2$  case. The latter does not hold for  $hmF2$ , even if the trends obtained are close to the expected values.

## 1 Introduction

The International Reference Ionosphere (IRI) (Bilitza et al., 2022) is an empirical model based on observations from diverse sources. Therefore, it is reasonable to expect it to reflect, to some extent, the long-term trends observed in key ionospheric parameters such as the F2 region critical frequency,  $f_oF2$ , and the electron density peak height,  $hmF2$ . These trends, in timescales of decades to a century, are theoretically expected as a consequence of trends in certain ionospheric forcings, such as the increasing greenhouse gas concentrations and the Earth's magnetic field secular variation, among others (Laštovička, 2017, 2021a).

There are countless  $f_oF2$  and  $hmF2$  reported trends based on experimental data, which, combined with model results, led to a global scenario of trends with the main forcing being the increasing greenhouse gas concentrations over the last decades (Laštovička, 2017, 2021a). However, several inconsistencies remain to date that require a permanent update of data and refinement of the trend estimation methods in order to improve the signal-to-noise ratio that is extremely weak in the case of ionospheric long-term trends. Additionally, experimental verification is still far from being achieved mainly due to two reasons: the limited time span and sparse spatial coverage of measured data. The time length should cover at least two complete solar cycles in order to efficiently filter out this variability that is essential for detecting long-term trends. Moreover, the ionosphere presents other challenges that need extensive series in order to properly identify and analyze trends. Regarding the spatial coverage, it should be global and with enough resolution so as to detect other forcings interfering with the expected trends whose intensity depends on location. This is the case, for example, with Earth's magnetic field secular variation effect on the ionosphere, which seems more prominent close to the geomagnetic equator in some longitudinal ranges (Cnossen, 2020; Elias et al., 2022). Given the difficulty of achieving these two requirements, we found it useful to evaluate trends from IRI modeled  $f_oF2$  and  $hmF2$  parameters and to analyze their usefulness as a reliable approximation of experimental trends.

This research initially focuses on presenting the trends spanning the entire planet. These trends are derived for  $f_oF2$  and  $hmF2$ , which are among the most significant ionospheric parameters (Cander, 2019). They are calculated following the same methodology applied to experimental data involving the simplest solar activity filtering approach. Furthermore, a comparative analysis is conducted between the trend values obtained from the IRI model and experimental trends in order to assess their accuracy. The continued refinement and updating of ionospheric trend estimation methods from data and models, together with data collection efforts, are essential for improving our understanding of the underlying factors driving long-term changes in ionospheric parameters and

their potential impacts on the diverse systems affected, such as communication and navigation systems.

This study is structured as follows. Sect. 2 provides an overview of the IRI model used. Sections 3 and 4 outline the methodology to derive global trends from IRI and to make a comparative analysis between these trends and experimental data of nine selected stations, respectively. The results are presented in Sect. 5, followed by a comparison with trends derived from a general circulation model in Sect. 6, and the discussions and conclusions are given in Sects. 7 and 8.

## 2 On some aspects of the IRI model

The IRI model is an observation-based climatological standard model of the ionosphere that is widely used for several purposes, including the prediction of ionospheric behavior useful for communication and global positioning systems (Gulyaeva and Bilitza, 2012). The model is designed to provide vertical profiles of the main ionospheric parameters for any location over the globe, hours, seasons, and levels of solar activity, representing monthly mean conditions based on experimental evidence. Even though the improvement of the IRI representation of ionospheric parameters, including those selected in this study, still remains a challenge for the IRI project, and despite its empirical nature and the potential for ongoing improvements, we choose to examine its suitability in estimating F2-region long-term trends.

Since its first edition in 1969, the IRI model has been steadily improved with newer data and with better mathematical descriptions of global and temporal variation patterns. A large number of independent studies have validated the IRI model in comparisons with direct and indirect ionospheric measurements not used in the model development (Gulyaeva and Bilitza, 2012; Bilitza et al., 2022).

In this study, we used an IRI adaptation, IRI-Plas, that has been modified to include the plasmasphere, extending the model up to 20 000 km (Gulyaeva et al., 2011). While traditional IRI versions use a given solar activity proxy, such as the Ionospheric Global index (IG) for  $f_oF2$ , to estimate variations in ionospheric parameters associated to the solar activity quasi-decadal cycle, IRI-Plas allows for selecting between eight different solar proxies, among them the MgII index (core-to-wing ratio derived from the magnesium II doublet at 280 nm). Since we chose this solar activity proxy for the filtering step before trend estimation, we decided to use this IRI version.

The IRI-Plas model from IZMIRAN (Institute of Terrestrial Magnetism, Ionosphere and Radiowave Propagation of the Russian Academy of Sciences) (Moscow, Russia) was used, which is available at <https://www.izmiran.ru/ionosphere/weather/> (last access: 12 September 2023).

According to IRI general specifications, long-term variations linked to changes in the geomagnetic field are ex-

pected since IRI uses the International Geomagnetic Reference Field (IGRF) model to specify magnetic poles and the equator, as well as the modified dip latitude, which is an input for  $foF2$  and  $hmF2$  interpolation procedures. Thus, trends due to the magnetic field changes, which are stronger near to the geomagnetic poles and equator, may arise from IRI mathematical interpolation coefficients, which ultimately depend on magnetic inclination. Since these changes are extremely small away from the geomagnetic equator, trends observed in other regions could be attributed to additional sources.

A key aspect in the present study is how IRI determines F2 parameters for a given location. To begin,  $foF2$  is obtained from CCIR (Consultative Committee on International Radio) maps that are based on a procedure of numerical mapping of a set of coefficients (CCIR Atlas of Ionospheric Characteristics, 1991) determined from a fitting to observed monthly median  $foF2$  data from a worldwide network of ionosonde stations ( $\sim 150$  in total). From these coefficient maps, IRI reproduces the diurnal, seasonal, and solar activity variation of  $foF2$  in terms of latitude and longitude through Fourier time series. First, there is a set of functions in terms of geographic coordinates and the modified dip latitude used to describe the variation of the Fourier coefficients for a given number of harmonics defining the diurnal variation. Then, the seasonal variation is taken into account through a set of these coefficients (988 in total) for every month of the year. Finally, the solar activity dependence is considered by having all these monthly coefficients for two different activity levels:  $IG_{12} = 0$  and  $IG_{12} = 100$ . From a linear fit between these two extremes, the harmonic coefficients for any solar activity level can be estimated. IG was originally computed using 13 globally distributed ionosonde stations that included two of the nine stations here analyzed: Kokubunji and Canberra (Liu et al., 1983). The distribution of these stations was a compromise between good global coverage and reliable long-operating-period ionosonde stations. Due to station closures and data unavailability, the number of stations used in IG has decreased to four but still includes the two stations used in the present study (Brown et al., 2018). Therefore, this proxy, being obtained from ionospheric measurements, involves  $foF2$  variations not covered by a solar index.

Specifically, when a given solar proxy is selected among the eight IRI-Plas options, it is automatically converted to other related indices used by the different modules' procedures (Gulyaeva et al., 2018). In this way,  $foF2$  interannual variation is determined by  $IG_{12}$ , since this index finally defines the CCIR coefficient values.

In the case of  $hmF2$ , we consider the default option, which corresponds to the AMTB-2013 model (standing for Altadill-Magdaleno-Torta-Blanch) (Altadill et al., 2013). This model is based on quiet ionosphere data from 26 digisondes collected between 1998 and 2006. The monthly averages of the global  $hmF2$  variations are represented by spherical harmonics including modified dip latitude and longitude for two selected levels of  $Rz_{12}$  (0 and 100, as in the case of IG). The

interannual variation of  $hmF2$  is defined then by  $Rz_{12}$  since, for a given date,  $hmF2$  is obtained from a linear fit of the spherical harmonic coefficients between  $Rz_{12} = 0$  and 100 particularized for the corresponding  $Rz_{12}$  value. The same procedure is applied in the cases of the other two options for  $hmF2$  modeling. Thus, the proxy used in this case, unlike the  $foF2$  case, is only reflecting solar activity variability. Nevertheless, we include its long-term trend analysis considering that the correlation between IG and  $Rz$  is higher than 0.99 and that, for a given location and hour,  $foF2$  and  $hmF2$  interannual variation highly correlates. Moreover, IG correlates the highest with  $Rz$  exceeding 0.99 throughout the period 1960–2022. The linear correlation between IG and MgII, F10.7 and Lyman- $\alpha$ , for example, are 0.975, 0.985 and 0.970, respectively.

### 3 Methodology to assess F2-region trends and spatial variation patterns based on IRI

To assess  $foF2$  and  $hmF2$  trends, monthly values were obtained first from IRI-Plas. This model was run over a  $5^\circ \times 10^\circ$  latitude–longitude grid, covering  $90^\circ$  N to  $90^\circ$  S and  $180^\circ$  E to  $180^\circ$  W, throughout the period 1960–2022, specifically at 00:00 and 12:00 LT, with the following inputs: (1) MgII as the solar activity proxy, (2) CCIR maps for  $foF2$ , (3) storm model off, and (4) AMTB-2013 model for  $hmF2$ . Considering just one day in the month or assessing the monthly median from all its daily values should give similar results due to the fact that the IRI model presents a smooth variation at daily timescale. Therefore, we considered  $foF2$  and  $hmF2$  values for the 15th day of each month as equivalent to the monthly median. Selecting other days or estimating all daily values within a month to assess the true median does not significantly affect the final results, as is discussed later in the Discussion section.

A total of  $37 \times 37 = 1369$  series were obtained for  $foF2$  and  $hmF2$ . In each case, annual mean series were constructed, together with series for each of the 12 months (that is, 13 series per grid point and per local time considered), all covering the period 1960–2022, which implies 63 points per series.

The  $foF2$  and  $hmF2$  filtering was made in the usual way by estimating the residuals from a linear regression with MgII as the solar extreme UV (EUV) proxy (Laštovička, 2021b, c), according to

$$X_{\text{residual}} = X_{\text{IRI}} - (A \cdot \text{MgII} + B), \quad (1)$$

where  $X_{\text{IRI}}$  is the IRI modeled  $foF2$  or  $hmF2$  data, and  $A$  and  $B$  are the least square parameters of the linear regression between  $X_{\text{IRI}}$  and MgII. The linear trend was assessed from the linear regression between these residuals and time; that is,

$$X_{\text{residual}} = \alpha t + \beta, \quad (2)$$

where  $t$  is in years, and  $\alpha$  is the desired trend [ $\text{MHz yr}^{-1}$ ] for  $foF2$  or [ $\text{km yr}^{-1}$ ] for  $hmF2$ . We will then have one  $\alpha$  value

for each grid point for the annual and for the 12 monthly series. Global means were also calculated in each case using a cosine (latitude) weighting.

The selection of MgII as the solar proxy input for IRI-Plas (and to filter  $foF2$  and  $hmF2$  variability linked to solar activity) is based on recent studies which recommend the use of this index as a solar proxy for  $foF2$  trend estimations (Laštovička, 2021b, c; de Haro Barbas et al., 2021). We assume it is also the most adequate in the case of  $hmF2$ .

The MgII index was obtained from the University of Bremen. It is freely available at <http://www.iup.uni-bremen.de/UVSAT/Datasets/MgII> (last access: 12 September 2023) (Viereck et al., 2010; Snow et al., 2014). The extended time series was considered in order to cover the period before 1978.

To determine trends induced by Earth's magnetic field secular variation only, we also run IRI-Plas for fixed solar activity conditions by keeping  $R_z$  constant at a mean level while running the years from 1960 to 2022. Trends were assessed directly through Eq. (2). A previous filtering is not needed since the only  $foF2$  and  $hmF2$  time variations generated by the model are those linked to the slow changes of the modified dip at each location.

#### 4 Methodology to evaluate the agreement between trends based on IRI and true experimental trends

Only  $foF2$  was considered in the comparison between IRI and experimental trend values. In order to assess the level of agreement between model and data, nine stations were chosen, which are listed in Table 1. Trends were evaluated using Eq. (1) to filter the solar activity effect and Eq. (2) to estimate trends in two ways: using the monthly median data, which will be called experimental trends ( $\alpha_{\text{exp}}$ ), and the IRI-Plas model output, which will be called IRI trends ( $\alpha_{\text{IRI}}$ ).

The following metrics commonly used in data-model prediction comparisons (Willmott and Matsuura, 2005; Chicco et al., 2021) were considered to compare IRI to experimental trends: the mean relative error (MRE) and the mean absolute error (MAE). Their equations are

$$\text{MRE} = \frac{1}{n} \sum \frac{(\alpha_{\text{IRI}} - \alpha_{\text{exp}})}{\alpha_{\text{exp}}}, \quad (3)$$

$$\text{MAE} = \frac{1}{n} \sum |\alpha_{\text{IRI}} - \alpha_{\text{exp}}|. \quad (4)$$

These parameters were assessed to determine overall IRI performance and also for each station separately. In the first case, summation is carried over the nine stations considering the annual mean series for 12:00 and 00:00 LT. In the second case, summation is carried out for each station over the 12 months.

MRE measures the average bias of IRI trends that are overestimating or underestimating the experimental trends depending on its sign: positive or negative, respectively. It gives

similar information to the percentage bias, and its optimal value is 0. MAE is a scale-dependent measure of deviation that corresponds to the IRI trend deviations from experimental ones. The optimal value of MAE is 0, indicating that both trends are identical.

Monthly median  $foF2$  data from the ionospheric stations were obtained as follows. Japanese and Australian station data are available from the National Institute of Information and Communications Technology, Japan ([https://wdc.nict.go.jp/IONO/index\\_E.html](https://wdc.nict.go.jp/IONO/index_E.html), last access: 12 September 2023), and the World Data Centre (WDC) for Space Weather, Australia (<https://downloads.sws.bom.gov.au/wdc/iondata/au/>, last access: 12 September 2023), respectively. Both databases contain monthly medians updated to 2022. European station monthly medians up to 2009 were obtained from the Damboldt and Suessmann database (Damboldt and Suessmann, 2012) (<https://downloads.sws.bom.gov.au/wdc/iondata/medians/>, last access: 12 September 2023). In the case of Juliusruh, the period was updated until 2022 with monthly medians available from [https://www.ionosonde.iap-kborn.de/mon\\_fof2.htm](https://www.ionosonde.iap-kborn.de/mon_fof2.htm) (last access: 12 September 2023). In the case of Boulder and Rome, the period was updated with data from the Lowell GIRO Data Center (LGDC) (Reinisch and Galkin, 2011). The  $foF2$  from the Digital Ionogram Data Base (DIDBase) at LGDC has a frequency of 5 min. In order to obtain the monthly medians, we first selected data with autoscaling confidence score (CS) greater than 70 %, and then we estimated for each month the hourly median. In the case of these two stations, it was checked that the last 2 years available from the Damboldt and Suessmann database had a reasonable coincidence (within 95 %) with the data obtained from the other two sources.

## 5 Results

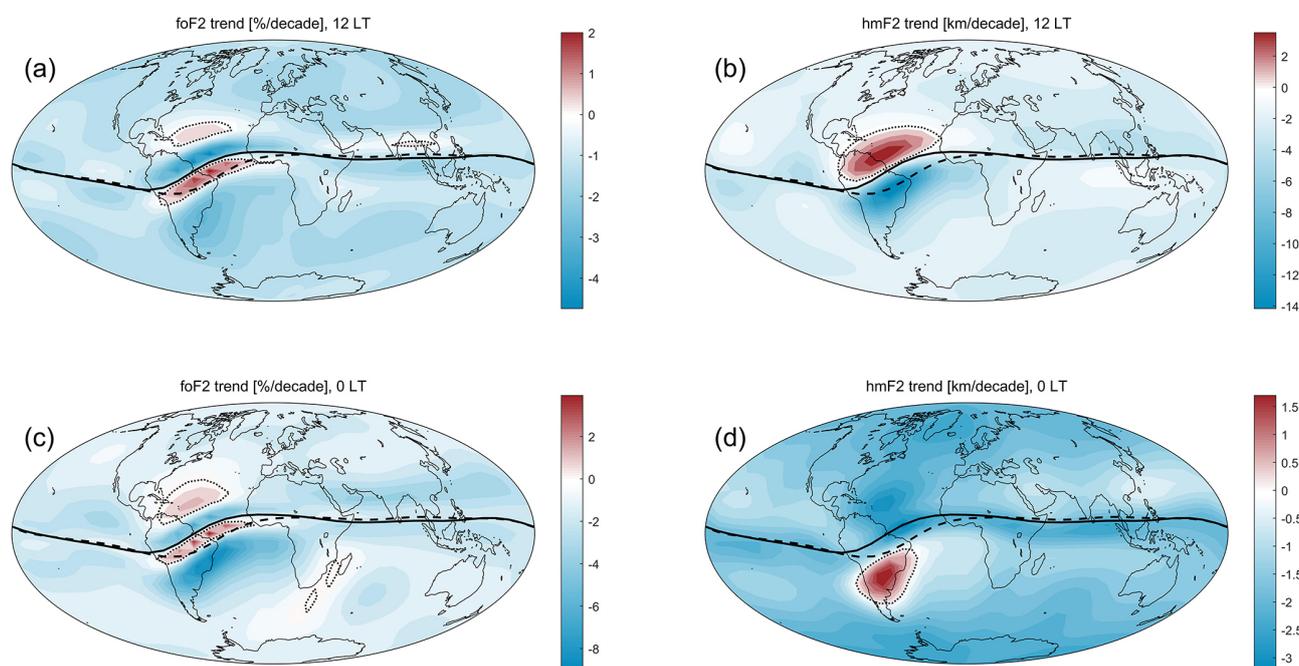
### 5.1 $foF2$ and $hmF2$ trends based on the IRI model, as well as spatial variation patterns

Figure 1 shows  $foF2$  and  $hmF2$  trend values for 12:00 and 00:00 LT. The geomagnetic equator is also plotted for years 1960 and 2022. The  $foF2$  trends are plotted in percent, which were estimated by dividing  $\alpha$  into  $foF2$  means along the complete period (1960–2022) at each grid point. In addition to overall negative trends in all cases, it can be noticed that the strongest trends occur in the region of the greatest geomagnetic equator displacement.

The global mean trends in each case are listed in Table 2, together with the mean values of F2-region parameters, to which the peak electron density,  $NmF2$ , was added in order to make some comparisons with other published results in the next section. Trends are listed in absolute and percentage values. The squared correlation coefficient,  $r^2$ , of each parameter and MgII is also listed to indicate the quality of the fit to each regression model given by Eq. (1).

**Table 1.** Geographic coordinates and geomagnetic latitude of the nine ionospheric stations analyzed to determine IRI  $f_oF2$  trends accuracy.

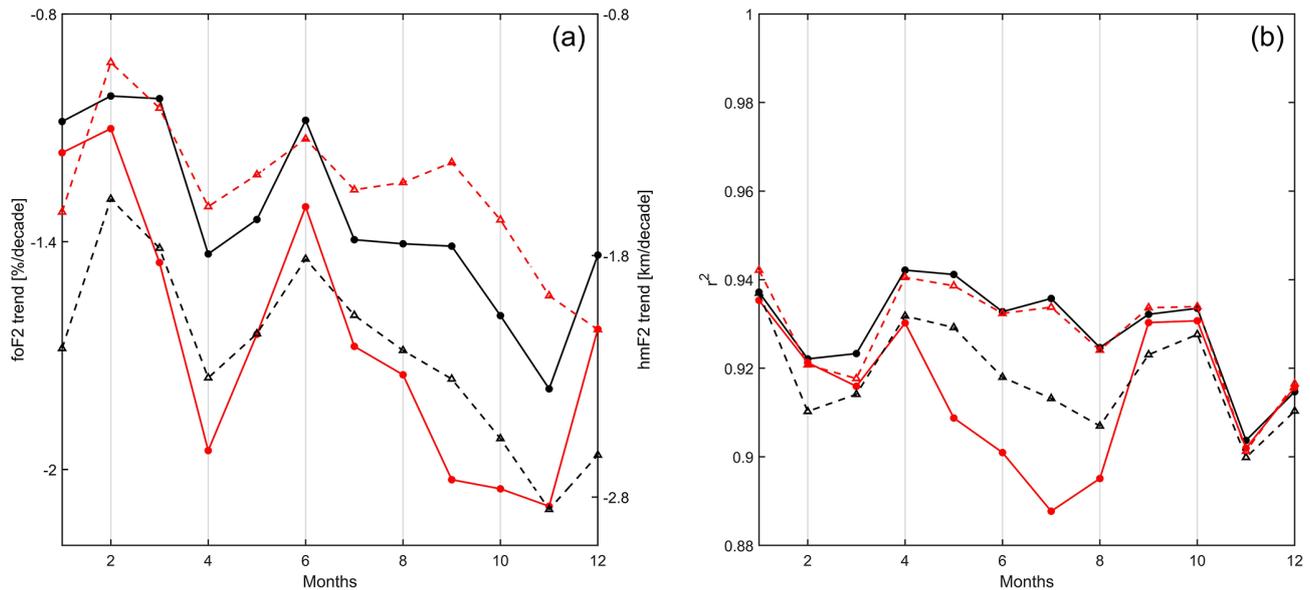
Station	Geographic latitude [°]	Geographic longitude [°]	Geomagnetic latitude [°]
Okinawa	26.31	127.59	17.28
Wakkanai	45.25	141.40	37.06
Kokubunji	35.71	139.49	27.28
Canberra	−35.17	149.08	−41.74
Townsville	−19.16	146.48	−26.21
Hobart	−42.53	147.19	−49.22
Juliusruh	54.60	13.40	53.98
Boulder	40.13	−105.23	47.65
Rome	41.54	12.29	41.49

**Figure 1.** Trends of  $f_oF2$  (a, c) and  $hmF2$  (b, d) at 12:00 LT (a, b) and 00:00 LT (c, d) throughout the period 1960–2022 assessed with the IRI outputs, which were previously filtered using Eq. (1). Note that trends are indicated per decade and that  $f_oF2$  trends are in percent. Enhanced dashed and solid lines indicate the magnetic equator position in 1960 and in 2022, respectively. Dotted lines indicate zero trend.

Trends assessed for each month have a similar spatial pattern as the annual trends shown in Fig. 1, even though they are not identical. Figure 2 (left panel) shows the global mean trend values from January to December. Weaker global trends are noticed in February and in June. Something to notice is the decrease of  $r^2$  of the fit to filter solar activity, shown in Fig. 2 (right panel). All values are lower than the annual case. This would be due to the variation of  $f_oF2$  and  $hmF2$  associated with solar activity being efficiently described by the 12-month moving average of a solar proxy. When analyzing the time series corresponding to each month separately, considering the unsmoothed monthly values lower  $r^2$  because the inter-monthly variation is not eliminated. As an additional comment, in general, when considering solar EUV proxies, they are all more alike when the time

series compared consist of annual means rather than monthly or daily means. This is because, at these shorter timescales, each time series conserves distinct variability patterns that are erased when annual or 12-month running means are used.

Figure 3 shows trends for IRI-Plas run keeping  $R_z$  constant at a mean level ( $R_z=70$ ). These trends result, then, from the Earth's magnetic field secular variation only since they reflect the modified dip changes at each location. From a comparison with Fig. 1 trend values and spatial patterns, two things become clear: (1) the positive and negative trend spatial configuration is due to the magnetic field variation, and (2) the overall negative trends, away from the region with the pronounced geomagnetic field equator displacement throughout the period considered, are not due to the magnetic field effect. Global mean trends in the case of Fig. 3 are



**Figure 2.** Global mean values of the linear trends (a) and the squared correlation coefficient ( $r^2$ ) between each parameter and MgII (b) of  $foF2$  (solid line with circles) and  $hmF2$  (dashed lines with triangles) at 12:00 (in black) and 00:00 LT (in red).

**Table 2.** Global mean values, using a cosine (latitude) weighting, of F2-region ionospheric parameters, squared correlation coefficient ( $r^2$ ) of each parameter and MgII, linear trends of filtered parameters indicated in units per decade, and the same trends in percentage per decade.

	Mean	$r^2$	$\alpha$	$\alpha$ [% per decade]
$foF2(12:00\text{ LT})$	7.78 MHz	0.967	−0.10 MHz per decade	−1.31
$foF2(00:00\text{ LT})$	4.87 MHz	0.962	−0.08 MHz per decade	−1.62
$NmF2(12:00\text{ LT})$	$8.05 \times 10^5\text{ cm}^{-3}$	0.970	$-2.03 \times 10^4\text{ cm}^{-3}$ per decade	−2.57
$NmF2(00:00\text{ LT})$	$3.18 \times 10^5\text{ cm}^{-3}$	0.963	$-1.15 \times 10^4\text{ cm}^{-3}$ per decade	−3.17
$hmF2(12:00\text{ LT})$	303.1 km	0.959	−2.16 km per decade	−0.72
$hmF2(00:00\text{ LT})$	323.0 km	0.971	−1.50 km per decade	−0.47

−0.0004 MHz per decade and −0.086 km per decade. In percentage they become −0.0006 % per decade and −0.023 % per decade, respectively. When comparing these values with those listed in Table 2 for 12:00 LT, it could be said that the global mean trend driven by the Earth’s magnetic field, despite being relatively strong at some regions, averages essentially to zero. The  $foF2$  and  $hmF2$  global means in this case are 7.93 MHz and 308.6 km, similar to the Table 2 values.

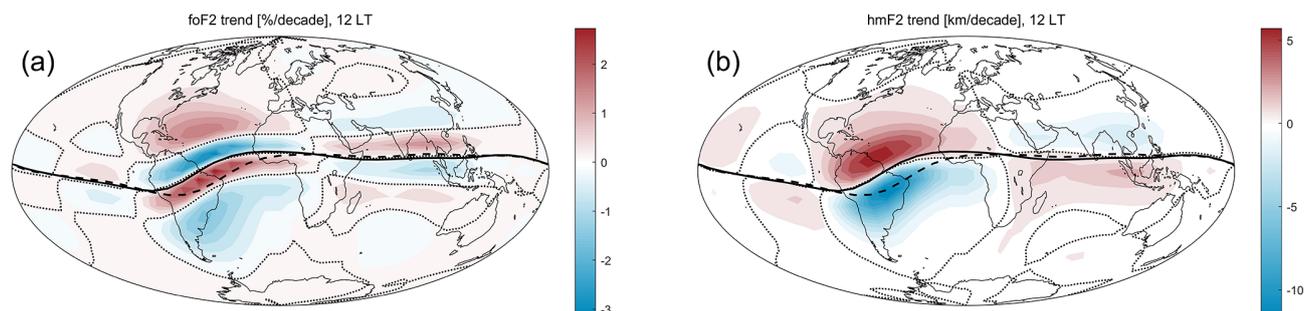
## 5.2 Agreement between IRI and experimental trends for selected stations

Figures 4 and 5 show experimental and IRI  $foF2$  trends for each of the nine stations, at 12:00 and 00:00 LT, respectively, in terms of months. Error bars are estimated as  $1\sigma$ . Generally good agreement can be noticed, which is evinced by MAE and MRE values listed in Table 3, in particular for the 12:00 LT case. Annual experimental and IRI trends are listed in Table 4.

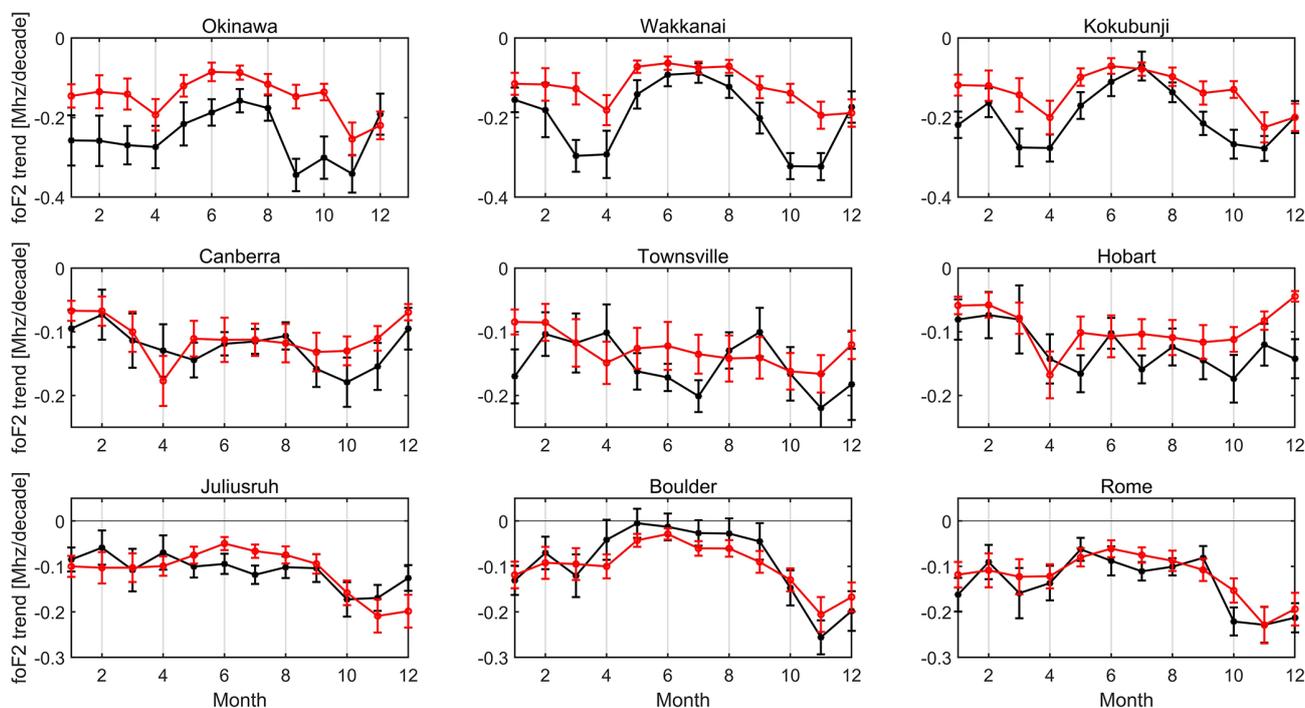
The cases with large MRE values correspond to those stations and local times that have an experimental trend value very close to zero. Since this value appears in the denominator of MRE (see Eq. 3), even a small difference in the numerator leads to a large MRE. However, we can re-estimate MRE values excluding experimental trends equal to zero within the error. Specifically, in the 12:00 LT case, these would correspond to experimental trend values for Boulder in May. In the 00:00 LT case, these would correspond to Kokubunji in February and December, Townsville in June, Juliusruh in February, and Boulder in September and October. By doing so, the MRE decreases, as indicated by the values presented within round brackets in Table 4.

## 6 Comparison with a general circulation model

The Whole Atmosphere Community Climate Model eXtension (WACCM-X) has been run to assess trends in the upper atmosphere (Solomon et al., 2018; Cnossen, 2020) and some results can be analyzed comparatively with the mean global



**Figure 3.** Trends of  $foF2$  (a) and  $hmF2$  (b) at 12:00 LT throughout the period 1960–2022 assessed with the IRI outputs using Eq. (1), without previously filtering. Note that trends are indicated per decade, and  $foF2$  trends are in percent. Enhanced dashed and solid lines indicate the magnetic equator position in 1960 and in 2022, respectively. Dotted lines indicate zero trend.

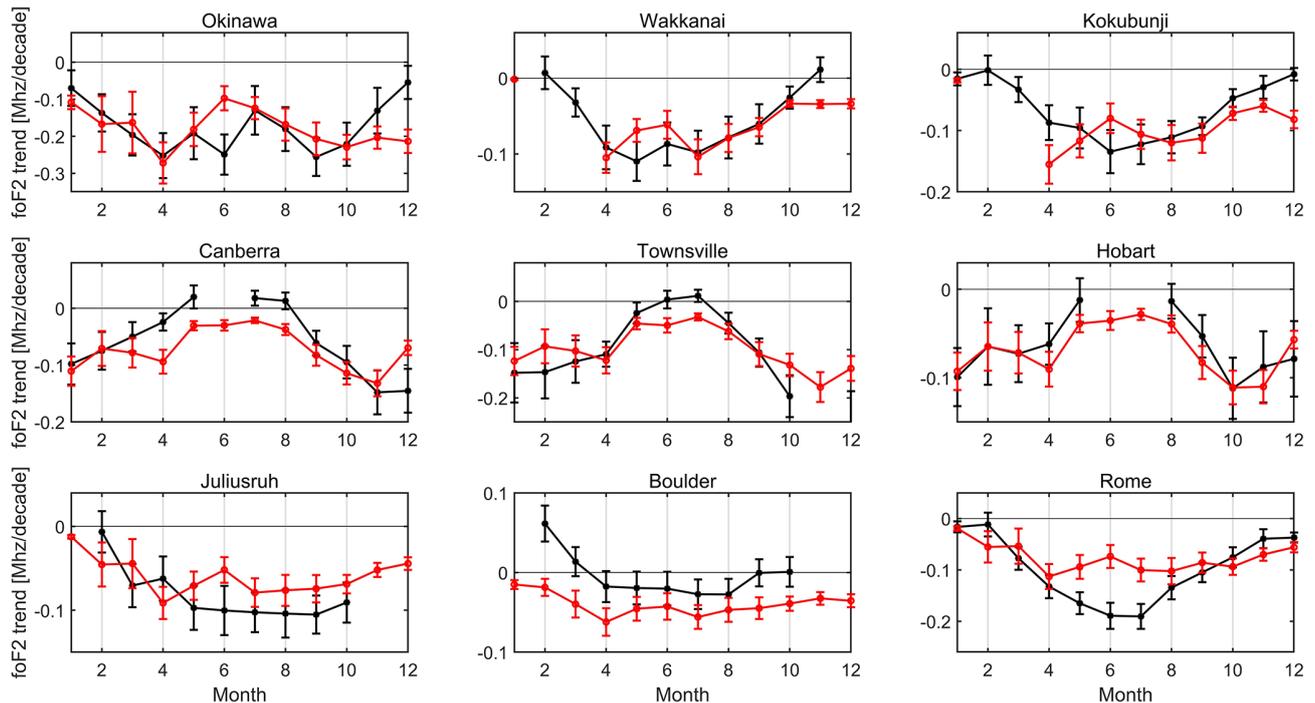


**Figure 4.** Monthly variation of  $foF2$  trends [MHz per decade], at 12:00 LT, estimated with experimental data (black) and with the IRI-Plas model (red). Error bars correspond to 1 standard deviation.

trends here obtained with IRI-Plas, as well as the spatial variation of the trends.

WACCM-X is a general circulation and complex model with high-resolution modeling capabilities, which incorporates a comprehensive set of physics processes to estimate a more realistic representation of the atmospheric (and ionospheric) status, including chemical, dynamical, and radiative processes. This model is coupled with several Earth systems, making it easy to analyze the weighting of any change in trends, e.g., the increase of a particular component in atmospheric composition. The trend results obtained by Solomon et al. (2018) and by Cnossen (2020) with WACCM-X that can be compared with those of IRI-Plas are listed in Table 5.

In the case of Solomon et al. (2018), global mean values are presented considering only minimum solar activity level and solar quiet conditions, with which no filtering is needed before the trend assessment. The period considered is 1972–2005, and there is no local time consideration, so we will assume that their values could be compared to the mean of our 12:00 and 00:00 LT values. Their trends are weaker than assessed with IRI-Plas, even if we reassess trends considering 1972–2005 instead of 1960–2022. In both cases, trends are negative, but the  $NmF2$  trend they obtain is around half the IRI-Plas trend, as can be deduced from Table 5. In addition to trend values,  $NmF2$  and  $hmF2$  mean global values estimated by Solomon et al. (2018) can be compared to IRI-Plas output averages considering only years around solar minimum



**Figure 5.** Monthly variation of  $f_oF2$  trends [MHz per decade], at 00:00 LT, estimated with experimental data (black) and with the IRI-Plas model (red). Error bars correspond to 1 standard deviation.

**Table 3.** The  $f_oF2$  trends assessed with experimental data and with the IRI-Plas model, considering annual mean data series at 12:00 and 00:00 LT. The last rows present the MAE and MRE between these trends carried over the nine stations. MRE\* corresponds to MRE without the stations of highest relative error; that is, Okinawa in the 12:00 LT case and Boulder in the 00:00 LT case.

Station	$\alpha$ [MHz per decade], 12:00 LT		$\alpha$ [MHz per decade], 00:00 LT	
	Experimental	IRI	Experimental	IRI
Okinawa	$-0.30 \pm 0.07$	$-0.14 \pm 0.03$	$-0.18 \pm 0.07$	$-0.18 \pm 0.03$
Wakkanai	$-0.18 \pm 0.05$	$-0.12 \pm 0.03$	$-0.04 \pm 0.02$	$-0.06 \pm 0.01$
Kokubunji	$-0.20 \pm 0.03$	$-0.13 \pm 0.03$	$-0.07 \pm 0.02$	$-0.09 \pm 0.02$
Canberra	$-0.12 \pm 0.04$	$-0.11 \pm 0.02$	$-0.05 \pm 0.03$	$-0.07 \pm 0.02$
Townsville	$-0.16 \pm 0.05$	$-0.13 \pm 0.03$	$-0.09 \pm 0.06$	$-0.10 \pm 0.02$
Hobart	$-0.13 \pm 0.04$	$-0.10 \pm 0.02$	$-0.05 \pm 0.03$	$-0.07 \pm 0.02$
Juliusruh	$-0.11 \pm 0.03$	$-0.11 \pm 0.02$	$-0.06 \pm 0.03$	$-0.06 \pm 0.01$
Boulder	$-0.08 \pm 0.06$	$-0.09 \pm 0.03$	$0.01 \pm 0.02$	$-0.04 \pm 0.01$
Rome	$-0.13 \pm 0.04$	$-0.12 \pm 0.03$	$-0.10 \pm 0.01$	$-0.08 \pm 0.02$
MAE	0.04 MHz per decade		0.03 MHz per decade	
MRE	$-0.19$ (−19 %)		$-0.79$ (−79 %)	
MRE*	$-0.15$ (−15 %)		$-0.16$ (−16 %)	

activity levels out of the 1960–2022 period. In this case, the results are similar for  $NmF2$ , but for  $hmF2$  their mean value is lower than that obtained with IRI-Plas.

Crossen (2020) presents the global mean values, as in the previous case, and the spatial pattern variation. Our trend estimation methodology is similar to “Model 1”, which includes terms for only solar activity and a long-term trend, with two differences: F10.7 is used instead of MgII as solar

proxy, and the trend term is included in a multiple regression together with the solar activity term. The differences due to methodologies are not expected to be significant (Laštovička et al., 2006). Absolute values of trends in this case are slightly higher than in Solomon et al. (2018) but again lower than those of IRI-Plas, with the greatest difference in the  $NmF2$  trend case, as can be noticed from Table 5. The squared correlation coefficient, that indicates the quality of

**Table 4.** MAE and MRE of *foF2* trends assessed with experimental data and with the IRI-Plas model, considering monthly data series at 12:00 and 00:00 LT, for each station. MRE values between round brackets correspond to estimations excluding experimental trends equal to zero within the error.

Station	$\alpha$ , 12:00 LT		$\alpha$ , 00:00 LT	
	MAE [MHz per decade]	MRE	MAE [MHz per decade]	MRE
Okinawa	0.10	−0.39	0.05	0.27
Wakkanai	0.08	−0.35	0.02	−0.52
Kokubunji	0.06	−0.29	0.03	1.19 (0.28)
Canberra	0.02	−0.11	0.04	−0.48
Townsville	0.04	−0.10	0.04	−1.67 (−0.39)
Hobart	0.04	−0.23	0.02	0.02
Juliusruh	0.03	0.04	0.03	0.48 (−0.21)
Boulder	0.03	1.07 (0.51)	0.04	6.26 (0.21)
Rome	0.03	−0.08	0.04	0.26

**Table 5.** Comparison between IRI-Plas and WACCM-X results from Solomon et al. (2018) and Cnossen (2020). All values correspond to global means throughout the period analyzed in each case, with the exception of *NmF2* and *hmF2* mean values which correspond to global means along solar minimum activity level periods.

	IRI-Plas (12:00 LT)	IRI-Plas (00:00 LT)	Solomon et al. (2018)	Cnossen (2020)
<i>NmF2</i> trend [% per decade]	$-2.6 \pm 0.8$	$-3 \pm 2$	−1.2	$-1.6 \pm 0.3$
<i>hmF2</i> trend [km per decade]	$-2 \pm 1$	$-1.5 \pm 0.5$	−1.3 km	$-1.5 \pm 0.1$
$r^2$ ( <i>NmF2</i> , MgII)	0.97	0.96	–	0.95
$r^2$ ( <i>hmF2</i> , MgII)	0.96	0.97	–	0.94
<i>NmF2</i> mean [ $\text{cm}^{-3}$ ]	$2.14 \times 10^5$	$1.39 \times 10^5$	$1.74 \times 10^5$	–
<i>hmF2</i> mean [km]	302.3	269.5	259.6	–

the fit to each regression model given by Eq. (1), is similar in all the cases.

Is important to remark that the trends reported by Solomon et al. (2018) may have resulted in lower values because they run the simulation with constant low solar activity. This would have neglected part of the trend that may be induced by the solar EUV flux negative trend along the last minima periods and which we consider partly responsible for the overall negative trends observed in measured ionospheric data.

The spatial variation pattern can also be compared. In the case of *hmF2*, the Cnossen (2020) spatial pattern is consistent with the IRI-Plas trends at 12:00 LT, with overall negative trends and a positive patch above the geomagnetic equator between Africa and South America. This would be in agreement with the trend expected from the northward geomagnetic equator secular displacement, which is strongest in this region, and assuming that the equatorial ionization anomaly (EIA) pattern of *hmF2* moves along with this displacement. The highest decreasing trends are observed consequently below this positive patch, in response to the northward movement of *hmF2* highest values. With respect to trend values, the strongest positive trends seem similar, around 2 km per

decade. However, the highest negative trends are greatest in the IRI-Plas case, reaching values of 12 km per decade at noontime, while in the Cnossen (2020) case this value corresponds to 5 km per decade.

In the case of *hmF2* trends at 00:00 LT, *hmF2* presents a trough, even though not as well defined as the crest during daytime hours. The displacement of this trough attached to the geomagnetic field northward displacement induces an effect inverse to that during noon. That is, a positive trend patch appears below it, with the strongest negative trends above.

A similar situation occurs with the *foF2* trend spatial pattern. In order to compare the trend values in percent with those of *NmF2*, they should be multiplied by 2. This means that our strongest negative trend is again the highest. The spatial pattern here has alternating bands of positive and negative trend values aligned with the EIA, which can be explained in terms of the EIA displacement following the geomagnetic equator. Between  $\sim 60^\circ$  W and  $0^\circ$  in longitude, the equator shift is the greatest and northward, so this is the region where the strongest alternating trends are noticed (Elias et al., 2022). This longitudinal extension is narrower than in the Cnossen (2020) case, who detect it between  $\sim 60^\circ$  W and  $\sim 20^\circ$  E. A notorious difference is that between the ini-

tial and the final position of the geomagnetic equator in this longitudinal range, Cnossen (2020) detects a negative trend band, while in our case a positive band is observed.

This difference may be caused by a poor resolution in latitude. In order to see the trend bands expected in the region between the initial and final position of the equator, a schematic plot is shown in Fig. 6a of the EIA  $f_oF2$  trough in its initial and its final position in 1960 and 2022, respectively. In this figure, it can be clearly noticed that the region between these positions will present a positive portion followed by a negative one. The first one corresponds to the region with low  $f_oF2$  in 1960, which has now become a region with higher  $f_oF2$  values (since the trough has moved). The second one corresponds to a region of higher  $f_oF2$  in 1960, which now is located under the EIA trough. On average, the geomagnetic equator has displaced  $\sim 5$  to  $10^\circ$  in the region's strongest shift, so for low resolutions, the grid points may coincide with one of either trend bands. This could partly explain the difference between the Cnossen (2020) negative band between the equator positions and our corresponding positive band. Figure 6b shows an enlarged portion of the trend's spatial pattern obtained with IRI-Plas but increasing the latitude resolution to  $1^\circ$ , where a positive and a negative band within the limits of the 1960 and the 2022 equator positions can be noticed.

The spatial pattern linked to the EIA displacement following the geomagnetic equator and clearly isolated in Fig. 3 is expected in IRI-Plas  $f_oF2$  and  $hmF2$  modeling since the model includes a real geomagnetic field. Even though there are very few stations along its location, the IRI model reproduces the EIA pattern through the variation in the magnetic inclination, obtained from IGRF, on which interpolation coefficients depend.

## 7 Discussion

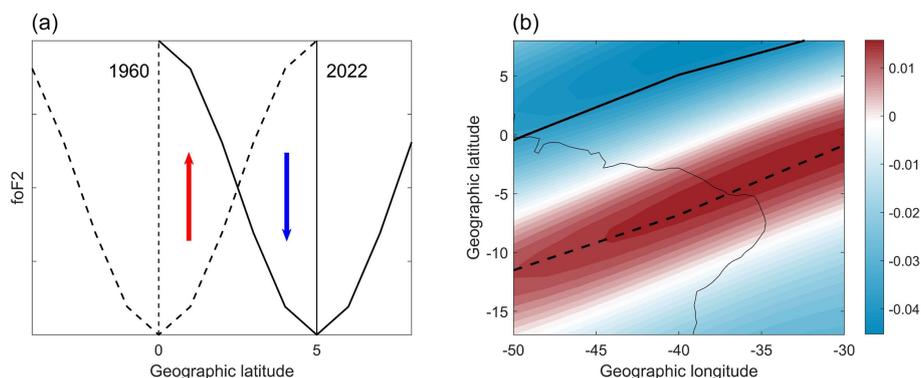
It is worth noting that in very recent IRI studies, the 30 cm solar flux index, F30 (available at <https://spaceweather.cls.fr/services/radioflux/>, last access: 12 September 2023), is identified as the most suitable EUV solar proxy for filtering  $f_oF2$  to subsequently estimate long-term trends, followed by MgII (Laštovička and Burešová, 2023). We also conducted a recent study (Zossi et al., 2023) where we concluded that both F30 and MgII are equally appropriate but without being able to distinguish which one is better of the two. In the present work, we did not consider F30 because IRI-Plas does not have this option. However, we compared the trend values for the nine stations here analyzed, from measurements and IRI-Plas model, considering each of these indices to filter solar activity effect, and, while we did not obtain identical values, they are in strong agreement. This agreement is nearly complete in terms of sign and practically within the error range of the trends in terms of values. Nevertheless, this deserves a detailed comparative analysis and could possibly suggest the

inclusion of F30 as an additional index to the options already available in this model.

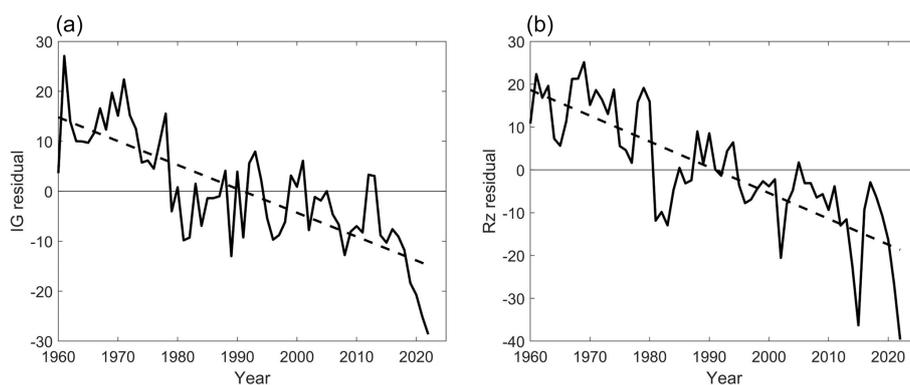
Another important aspect concerns IG and Rz long-term trends. The explained variance of each of these proxies by MgII is  $\sim 95\%$  ( $r^2 \times 100$ ) in both cases (IG vs. MgII and Rz vs. MgII), and if the solar activity effect is filtered from them through the same linear regression as that performed on F2 parameters, negative trends are obtained in both residuals as is shown in Fig. 7. The decreasing trend observed in IG when filtered with MgII is mainly due to the last two solar cycle minima which are much lower than the previous two in the IG case (and also in the Rz case) than in the MgII case. This is due to the fact that the solar EUV flux during the last two minima has been lower than the values indicated by solar proxies. This would also induce a decreasing trend in  $f_oF2$  (and in  $hmF2$ ) which might be connected to the inadequate performance of the proxy in capturing the variations in solar EUV flux (Emmert et al., 2010; Chen et al., 2011; Bruevich and Bruevich, 2019; Elias et al., 2023). However, this is a topic for further research.

Another aspect to discuss is the use of  $f_oF2$  and  $hmF2$  values for a single day of each month as representative of the monthly median. On the one hand, as mentioned in Sect. 3, the IRI model, with the “storm off” option, exhibits a smooth daily variation throughout each month. To analyze the impact of choosing this particular day on the trend instead of the monthly median or mean obtained from all daily values, we assessed annual noon  $f_oF2$  trends for a mid-latitude location ( $20^\circ$  N,  $30^\circ$  E), considering other days of each month (but using the same day for every month and year). We also assessed the trends by considering the median and the mean value of each month. Even though the trend values are not the same, the difference between any of them is around  $\sim 0.006$  MHz per decade, which is smaller than the trends' standard error ( $\sim 0.02$  MHz per decade). As an additional possibility, we considered using a random day in each month. For example, for the year 1960, we used day 12 for January, day 27 for February, day 5 for March, and so on, for the following months and years. From 10 000 random estimations we made, the minimum trend value obtained is  $-0.09$  MHz per decade, and the maximum value is  $-0.13$  MHz per decade. Both include within the error interval ( $\pm 0.02$  MHz per decade) the value of the trend obtained considering day 15 (which is  $-0.0110$  MHz per decade) and that considering the true  $f_oF2$  median (which is  $-0.0111$  MHz per decade). The most probable trend values in this running of 10 000 trend estimations lies between  $-0.111$  and  $-0.109$ , and it again includes the value estimated in this work considering day 15.

An additional topic deserving further research is the global picture easily obtained with IRI of the geomagnetic field secular variation-induced trends. We will not go deeper into this aspect in this work, but we considered it important to mention that the positive and negative trend patches are consistent with the results by Cnossen and Richmond (2012), who ana-



**Figure 6.** (a) Schematic representation of the  $foF2$  latitudinal profile around the EIA trough in 1960 (dashed line), centered in the magnetic equator in 1960 (vertical dashed line at latitude = 0), and in 2022 (solid line), centered in the magnetic equator in 2022 (vertical solid line at latitude = 5). The red arrow indicates the  $foF2$  increase that would be observed in latitudes between 0 and  $\sim 2.5$ , and the blue arrow indicates the decrease between  $\sim 2.5$  and 5. (b) The  $foF2$  trend along 1960–2022 assessed with IRI-Plas (solar activity filtering with MgII) in the region with the largest equator displacement with an increased resolution:  $1^\circ \times 2^\circ$  latitude–longitude grid. Enhanced dashed and solid lines indicate the magnetic equator position in 1960 and in 2022, respectively.



**Figure 7.** Residuals of the linear regression between annual means IG and MgII (a) and Rz and MgII (b).

lyzed the effect of the Earth’s dipole inclination variation using the Coupled Magnetosphere-Ionosphere-Thermosphere (CMIT) model. The only difference is that they show the effect of a dipole axis increasing its inclination, and the secular variation observed during the last decades is consistent with a dipole aligning with the rotation axis. That is why, in a rough comparison, the sign of the trend’s patches in Fig. 3 are opposite to those of Fig. 7 (lower panels) of Cnossen and Richmond (2012). This is something worth exploring, using IRI, at least as long as the field remains mainly dipolar.

## 8 Conclusions

Considering how the  $foF2$  and  $hmF2$  interannual variation is determined in IRI-Plas (and in other IRI model versions), it can be argued that the overall negative trends are due to the same long-term trend occurring in IG and Rz.

For  $foF2$ , the attribution to external forcings other than the magnetic field is clear since IG carries the information of  $foF2$  measurements. Thus, we can expect that the trends

obtained could be a reasonable approach to experimental trends. In fact, this index includes the variability by other sources affecting the ionospheric F2 layer, like the greenhouse gas concentration increases or other neutral composition changes or dynamical disturbances. It does not include, however, the magnetic field secular variation effect since it averages to almost zero. Hence, the  $foF2$  trends obtained using IRI-Plas model values can be, to a first approximation, attributed to the greenhouse cooling effect plus the secular variation in Earth’s magnetic field. This, of course, assumes that the dominant driver behind the global declining  $foF2$  trend, and of IG, is indeed the greenhouse effect. In addition, to verify this in a more localized spatial scale, we were able to compare  $foF2$  trends, considering annual and monthly series, determined with experimental data from nine mid-latitude stations and with the corresponding IRI-Plas modeled values. We obtained a reasonable agreement, with average differences of  $\sim 20\%$ . Something to argue is that by using mid-latitude stations, we use the stations for which IRI surely works best.

On the other hand, *hmF2* trends result from the Rz overall decreasing trend when it is filtered with MgII (or any other solar EUV proxy). Even though the values are coherent with expected trends due to greenhouse cooling (due to the fact that Rz varies almost identically to IG), we cannot conclude, using IRI, that its long-term lowering is due to greenhouse gas concentration increases. This is due to the coincidence that both *hmF2* from IRI and *hmF2* from measurements and theoretical considerations are forced by a mechanism inducing a downward trend; in the first case, it is the Rz overall downward trend throughout the period considered, while in the latter it would be due to the greenhouse effect. In addition, of course, the downward trend in Rz has nothing to do with the increasing concentrations of greenhouse gases during the last decades. They both just happen to be in the same direction. Despite this, it is considered worthwhile to compare these trends with experimental values as a future task. Smaller Rz values since  $\sim 2001$  have been mentioned some years ago by Lukianova and Mursula (2011) and Mielich and Bremer (2013).

To be able to attribute observed trends in IRI to processes unrelated to solar activity, it would be valuable to consider two potential approaches:

1. The first is interpolation of the CCIR maps with “effective indices” derived from data, as already proposed by Pignalberi et al. (2018). This would be similar to the approach used for IG. By using effective indices based on observed data, time variations not directly tied to solar activity would be accounted for.
2. The second is interpolation from annual CCIR maps, instead of the two maps currently used. This would involve updating the CCIR maps on an annual basis, assimilating the most recent and accurate data; thus, the time variation obtained would not result exclusively from solar activity variability.

Compared to a more theoretically based model, it is important to remark that IRI-Plas is designed for modeling a specific atmospheric sub-region, namely, the ionosphere, whereas WACCM-X is a global circulation model that simulates the entire atmosphere. Thus, an advantage of this model is that, in considering coupling processes among several Earth systems, it allows one to analyze the weighting of any change in trends, e.g., the increase of a particular component in atmospheric composition. Nevertheless, the negative side of general circulation models is the substantial computational resources and time needed to run simulations, being almost exclusive for high-performance computing centers. In the case of IRI, a great advantage is its user-friendly design, allowing it to be run on modest computers consuming few resources with extremely short computational times (on the order of minutes), while still being a reliable tool to get an approximated status of the ionosphere. On the other hand,

this reference model adopts simplified assumptions and parameterizations to represent complex ionospheric processes.

Before summarizing the answer to the question raised by this study’s title, we bring up again some recommendations for future tasks suggested throughout this work: (1) a comparison between *hmF2* trends at specific locations between IRI and ionosonde data, similar to *foF2* analysis; (2) the spatial pattern of IRI trends due only to the Earth’s magnetic field and its comparison with complex models with theoretical approaches; (3) the correct attribution of the general *foF2* and *hmF2* downward trend: the greenhouse affect or a long-term decreasing EUV flux not shown by EUV proxies which are used to filter solar activity effect since Rz was suggested to be discarded for this purpose (Mielich and Bremer, 2013); and (4) to modify IRI effective proxies or coefficient maps in order to determine *foF2* and *hmF2* interannual variations that include external sources other than solar activity only.

In summary, regarding the question set forth in this study’s title, we conclude that the IRI model can be a valuable tool for obtaining preliminary approximations of experimental trends, at least in the case of *foF2*. This is particularly significant given the low spatial density of data and the scarcity of time series with sufficient length to estimate trends. In the case of *hmF2*, there would be an added advantage considering that, while *foF2* is an accurately measured parameter, *hmF2* is often missing or derived from the proxy M(3000)F2 parameter. However, even if *hmF2* trends obtained with IRI-Plas are close to the expected values, they are linked to different drivers.

**Code and data availability.** The IRI-Plas code is freely available from the IZMIRAN Ionospheric Weather server at <https://www.izmiran.ru/ionosphere/weather/grif/SPIM/> (last access: 12 September 2023, Gulyaeva et al., 2011). The IRI 2016 and 2020 versions, also used in this work, were run from the HF (high-frequency) propagation toolbox, PHaRLAP, created by Manuel Cervera, Defence Science and Technology Group, Australia ([manuel.cervera@dst.defence.gov.au](mailto:manuel.cervera@dst.defence.gov.au)), available at <https://www.dst.defence.gov.au/our-technologies/pharlap-provision-high-frequency-raytracing-laboratory-propagation-studies> (last access: 12 September 2023). This toolbox is available by request from its author. The *foF2* data from Japanese and Australian stations are available at [https://wdc.nict.go.jp/IONO/index\\_E.html](https://wdc.nict.go.jp/IONO/index_E.html) (last access: 12 September 2023, WDC for Ionosphere and Space Weather, 2023), and <https://downloads.sws.bom.gov.au/wdc/iondata/au/> (last access: 12 September 2023, Australian Space Weather Forecasting Centre, 2023a), respectively. Boulder and Rome *foF2* data were obtained from <https://downloads.sws.bom.gov.au/wdc/iondata/medians/> (last access: 12 September 2023, Australian Space Weather Forecasting Centre, 2023b) and <https://doi.org/10.5047/eps.2011.03.001> (Reinisch and Galkin, 2011). Juliusruh *foF2* data are available from [https://www.ionosonde.iap-kborn.de/mon\\_fof2.htm](https://www.ionosonde.iap-kborn.de/mon_fof2.htm) (last access: 12 September 2023, Leibniz Institute of Atmospheric Physics at the University of Rostock, 2023). The MgII index is freely available at <http://www.iup.uni-bremen.de/UVSAT/Datasets/MgII>

(last access: 12 September 2023, Viereck et al., 2010; Snow et al., 2014), IG is from [https://www.ukssdc.ac.uk/cgi-bin/wdcd1/secure/geophysical\\_parameters.pl](https://www.ukssdc.ac.uk/cgi-bin/wdcd1/secure/geophysical_parameters.pl) (last access: 12 September 2023, UK Solar System Data Centre, 2023), and Rz is from <https://www.sidc.be/SILSO/datafiles> (last access: 12 September 2023, WDC-SILSO, 2023).

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