



*Supplement of*

## **Evolution of turbulent kinetic energy during the entire sandstorm process**

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## **The validation of adaptive segmented stationary method**

To validate the reliability of this method, a synthetic signal (as shown in Fig. S1a) is analyzed by using the adaptive segmented stationary method. The synthetic fluctuating signal consists of 7 components. Components 1–4 are amplitude variation signals with frequencies of 10 Hz, 5 Hz, 0.5 Hz and 0.1 Hz, respectively. Component 5 is an oscillation attenuation signal with a frequency of 1 Hz, Component 6 is a time-varying average, and component 7 is a real turbulent fluctuating signal provided in Wang and Zheng (2016). To enhance the non-stationarity of the synthetic signal, the amplitude of components 1-5 is increased monotonously with time. The IST of the signal is 79% (still 49% after removing the time-varying mean extracted by the EMD), indicating that the signal is non-stationary and cannot be directly analyzed by traditional statistical analysis methods. Therefore, the signal is segmented by employing the adaptive segmented stationary method, as shown in Fig. S1(b). The non-stationary signal is divided into three segments, and then the time-varying characteristics of the signal can be obtained by analyzing these three segments. Fig. S1 (c, d) shows the power spectra density and pre-multiplied spectra of these quasi-stationary signals and compared with the entire (unsegmented) turbulence fluctuating signal spectrum. It is seen in Fig. S1 (c) and S1 (d) that both the power spectra and the pre-multiplied spectra of the synthetic signal agree well with the spectrum of the turbulent fluctuating signal, except for a few extra peaks caused by the additional sinusoidal fluctuating signals. In addition, the energy at 1 Hz in the second and third segments decreases significantly because the amplitude of the oscillation attenuation component decreases with time. This indicates that the adaptive segmented stationary method can accurately exhibit the various frequency components contained in the signal and the time-varying characteristics when processing non-stationary signals. Therefore, the adaptive segmented stationary method is reliable to deal with non-stationary signals, and thus can be used to process the sandstorms data.

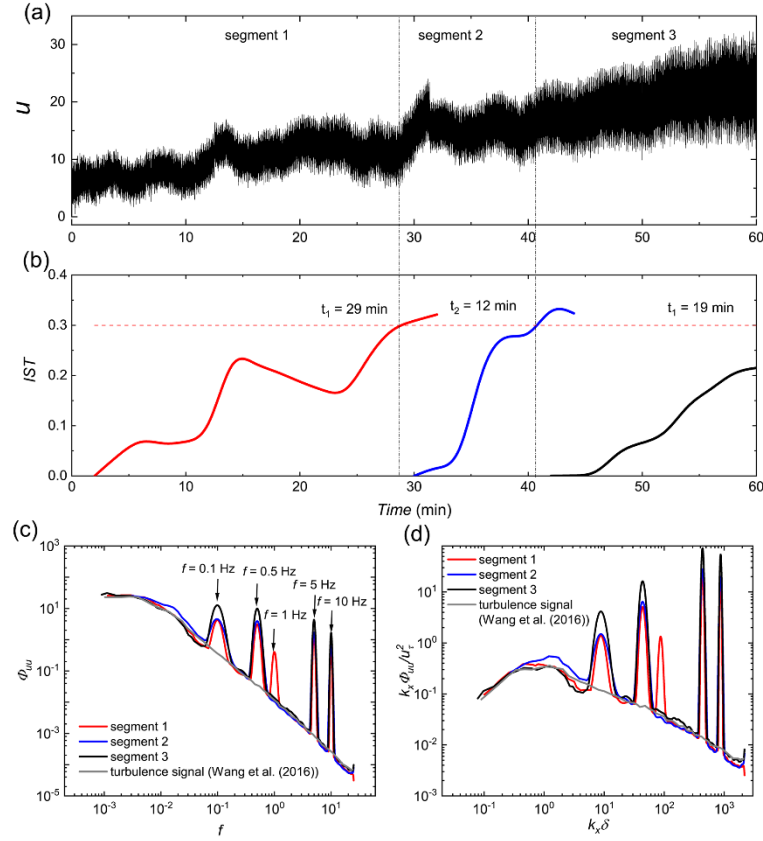


Figure S1. (a) Synthetic signal; (b) the change of the IST of the signal with time; (c) the power spectra density and (d) the pre-multiplied spectra of the synthetic signal. The gray line is the result of the entire (unsegmented) turbulent fluctuating signal contained in the synthetic signal. The red, blue and black lines represent the first, second and third segments of the signal, respectively.

## Spectral method

The spectral method is a useful and important tool in data analysis. For a signal  $x(t)$ , the  $k$  th-order spectrum (or  $k-1$  spectrum) can be obtained by a  $(k-1)$ -dimensional Fourier transform of the  $k$  th-order cumulant, which is given as

$$B_{k,x}(f_1, f_2, \dots, f_{k-1}) = \sum_{m_1=-\infty}^{\infty} \dots \sum_{m_{k-1}=-\infty}^{\infty} c_{k,x}(\tau_1, \tau_2, \dots, \tau_{k-1}) \exp \left[ -j \sum_{i=1}^{k-1} f_i \tau_i \right] \quad (S1)$$

where  $c_{k,x}(\tau_1, \tau_2, \dots, \tau_{k-1})$  is the  $k$  th-order cumulant of  $x(t)$ ,  $\tau_i$  ( $i=1, 2, \dots, k-1$ ) represents the temporal lead/lag,  $f_i$  is the frequency. In particular, when  $k=2$ , Equation (S1) represents the power spectrum,

$$P_{2,x}(f) = \sum_{m_1=-\infty}^{\infty} c_{2,x}(\tau) \exp[-jf\tau] = \langle X(f) X^*(f) \rangle \quad (\text{S2})$$

where  $\langle \bullet \rangle$  represents the expected value,  $X(f)$  is the Fourier transform of a segment of the time series record, and the asterisk denotes the complex conjugate. The power spectrum represents the energy distribution at different frequencies. For studies on turbulence, the frequency is usually converted to the length scale by using Taylor's hypothesis of frozen turbulence, that is,  $\lambda = \bar{u} / f$ , where  $\bar{u}$  is the velocity taken as the local mean. In addition, the pre-multiplied spectrum is obtained by multiplying the power spectral density by the frequency (or wavenumber).

When  $k=3$  in Equation (S1), we can obtain the bispectrum,

$$B_{3,x}(f_1, f_2) = \sum_{m_1=-\infty}^{\infty} \cdots \sum_{m_2=-\infty}^{\infty} c_{3,x}(\tau_1, \tau_2) \exp[-j(f_1\tau_1 + f_2\tau_2)] \quad (\text{S3})$$

$$= \langle X(f_1) X(f_2) X^*(f_1 + f_2) \rangle$$

which provides a measure of the quadratic interaction exhibited by the given triad of waves at frequencies of  $f_1$ ,  $f_2$  and  $f_3$  that satisfy the resonance wave matching requirement  $f_3 = f_1 + f_2$  (Hajj et al., 1997; Liu et al., 2019). And a nonzero bispectrum value means that phase coupling may occur between  $f_1$ ,  $f_2$  and  $f_3$ . As example, Fig. S2 shows the absolute bispectrum of the high frequency band of the streamwise velocity fluctuations in the sandstorm, where Fig. S2 (a) is the three-dimensional perspective and Fig. S2 (b) is the color contour map. It is seen that the bispectrum exists multiple peaks at different positions, indicating a strong quadratic phase coupling between the corresponding frequencies. In addition, the unphysical negative frequencies are derived from the process of bispectrum calculation. Thus, only the positive frequency component, that is, the first quadrant part, will be shown in the following study.

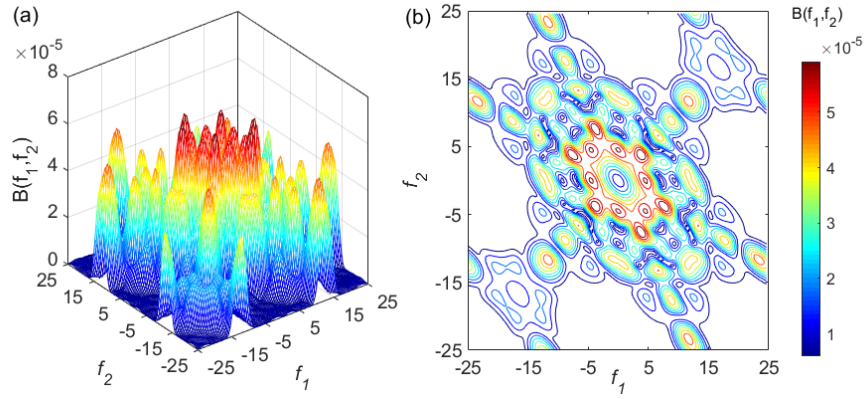


Figure S2. Bispectrum of the high frequency band of the streamwise velocity fluctuations in the sandstorm, (a) three-dimensional view, (b) colour contour map

To estimate the energy transfer process, the turbulent system can be represented by a nonlinear model, which are describable by a set of source (“input”) signals  $X$  and the response (“output”) signals  $Y$ . The model consists of linear, quadratic and higher order nonlinear elements (Ritz and Powers, 1986; Ritz et al., 1989),

$$Y_f = L_f X_f + \sum Q_f^{f_1, f_2} X_{f_1} X_{f_2} + \sum Q_f^{f_1, f_2, f_3} X_{f_1} X_{f_2} X_{f_3} + \dots + \varepsilon_f \quad (\text{S4})$$

where,  $L_f X_f$  represents the linear transfer  $Q_f^{f_1, f_2} X_{f_1} X_{f_2}$ , represents the three wave process,  $Q_f^{f_1, f_2, f_3} X_{f_1} X_{f_2} X_{f_3}$  represents the four wave process and  $\varepsilon_f$  is error term. Generally, the four wave process and the higher-order processes are much weaker than the three wave process (Ritz and Powers, 1986; Ritz et al., 1989), and thus Equation (S4) is simplified as

$$Y_f = L_f X_f + \sum_{\substack{f_1 > f_2 \\ f = f_1 + f_2}} Q_f^{f_1, f_2} X_{f_1} X_{f_2} + \varepsilon_f \quad (\text{S5})$$

By multiplying Eq. (S5) with the complex conjugate of the input signal  $X_f$ , neglecting the error term and expressing the process in terms of time derivatives, we can obtain the wave kinetic equation (Ritz et al., 1989; Cziegler et al., 2013):

$$\frac{\partial P_f}{\partial t} = \gamma_f P_f + \sum_{f_1, f_2} T_f(f_1, f_2) \quad (\text{S6})$$

where  $P_f = \langle Y_f X_f^* \rangle$  is the power spectrum,  $\frac{\partial P_f}{\partial t}$  is the time change of the spectrum of frequency  $f$ ,  $\gamma_f$  is the growth rate at the frequency  $f$ ;  $Q_f^{f_1, f_2}$  is the nonlinear transfer function, and  $T_f(f_1, f_2) = |Q_f^{f_1, f_2}| \langle X_{f_1} X_{f_2} X_f^* \rangle = |Q_f^{f_1, f_2}| B(f_1, f_2)$  is the nonlinear energy transfer function of the system, which represents a net flow of energy into or away from a given frequency  $f$ . Since we are concentrating on the effects of the nonlinear energy transfer in the process of sandstorm, the nonlinear transfer term under more scrutiny is  $T_f(f_1, f_2)$ . Eq.(S6) makes it clear that if  $T_f(f_1, f_2)$  is positive (negative), the components at “target” frequency  $f_2$  are gaining (losing) power and the components at the “source” frequency  $f_1$  are losing (gaining) it via the convection represented in the  $X_{f_2}$  term (Ritz et al., 1989; Cziegler et al., 2013).

For a turbulent system, where many waves interact with one frequency, the contribution of any given three components will be weakened, but the sum of all possible three components interactions can result in a significant change of power at this frequency (Ritz et al., 1989). Thus, by integrating the bispectrum along one of the frequency directions, we can obtain the integral bispectrum, which is defined as

$$AIB(f_1) = \int B(f_1, f_2) df_2 \quad (\text{S7})$$

where  $B(f_1, f_2)$  is the bispectrum of the signal. The integral bispectrum represents the total nonlinear interactions by all of the other frequency components on a certain frequency. It should be emphasized that  $f_1$  and  $f_2$  have the same meaning in the auto-bispectrum. Eq. (S6) and (S7) infer that when the integral bispectrum is positive (negative), the components at frequency  $f_2$  are gaining (losing) power from the interaction with all of the possible components.

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