



Reduced effective radiative forcing from cloud–aerosol interactions (ERF_{aci}) with improved treatment of early aerosol growth in an Earth system model

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Abstract. Historically, aerosols of anthropogenic origin have offset some of the warming from increased atmospheric greenhouse gas concentrations. The strength of this negative aerosol forcing, however, is highly uncertain – especially the part originating from cloud–aerosol interactions. An important part of this uncertainty originates from our lack of knowledge about pre-industrial aerosols and how many of these would have acted as cloud condensation nuclei (CCN). In order to simulate CCN concentrations in models, we must adequately model secondary aerosols, including new particle formation (NPF) and early growth, which contributes a large part of atmospheric CCN. In this study, we investigate the effective radiative forcing (ERF) from cloud–aerosol interactions (ERF_{aci}) with an improved treatment of early particle growth, as presented in Blichner et al. (2021). We compare the improved scheme to the default scheme, OsloAero, which are both embedded in the atmospheric component of the Norwegian Earth System Model v2 (NorESM2). The improved scheme, OsloAeroSec, includes a sectional scheme that treats the growth of particles from 5–39.6 nm in diameter, which thereafter inputs the particles to the smallest mode in the pre-existing modal aerosol scheme. The default scheme parameterizes the growth of particles from nucleation up to the smallest mode, a process that can take several hours. The explicit treatment of early growth in OsloAeroSec, on the other hand, captures the changes in atmospheric conditions during this growth time in terms of air mass mixing, transport, and condensation and coagulation.

We find that the ERF_{aci} with the sectional scheme is -1.16 W m^{-2} , which is 0.13 W m^{-2} weaker compared to the

default scheme. This reduction originates from OsloAeroSec producing more particles than the default scheme in pristine, low-aerosol-concentration areas and fewer NPF particles in high-aerosol areas. We find, perhaps surprisingly, that NPF inhibits cloud droplet activation in polluted and/or high-aerosol-concentration regions because the NPF particles increase the condensation sink and reduce the growth of the larger particles which may otherwise activate. This means that in these high-aerosol regions, the model with the lowest NPF – OsloAeroSec – will have the highest cloud droplet activation and thus more reflective clouds. In pristine and/or low-aerosol regions, however, NPF enhances cloud droplet activation because the NPF particles themselves tend to activate.

Lastly, we find that sulfate emissions in the present-day simulations increase the hygroscopicity of secondary aerosols compared to pre-industrial simulations. This makes NPF particles more relevant for cloud droplet activation in the present day than the pre-industrial atmosphere because increased hygroscopicity means they can activate at smaller sizes.

1 Introduction

Since pre-industrial times, humans have significantly shaped our climate by emitting greenhouse gases to the atmosphere. However, the warming induced from these emissions has been masked by the cooling effects of anthropogenic emissions of aerosols and their precursors (Myhre et al., 2013).

This cooling is highly uncertain and dominates the spread in estimates of radiative forcing and observationally based estimates of climate sensitivity (Myhre et al., 2013).

The present-day atmospheric aerosol state is challenging to fully characterize due to its fast-changing nature, making point observations hard to generalize. The pre-industrial atmosphere, however, is even more challenging since we cannot rely on direct observations, and it is thus only accessible by putting our best knowledge of aerosol processes and sources into models. The pre-industrial atmospheric state is furthermore very important for estimating the cooling by aerosol–cloud interactions (Carslaw et al., 2013) because the cloud albedo is more sensitive to perturbations in a “cleaner” atmosphere (Carslaw et al., 2013; Twomey, 1991). There are two main reasons for this. Firstly, cloud droplets form around cloud condensation nuclei (CCN) when the air mass is cooled, normally through adiabatic lifting. The number of particles that will act as CCN and form cloud droplets is dependent on the maximum achieved supersaturation during cloud formation and how many particles can activate at this supersaturation, which is dependent on size and hygroscopicity. If there are many large CCN, then these will activate “early” during cloud formation and constitute a water vapor sink, which limits the maximum supersaturation and therefore the number of CCN that can activate. We will refer to this effect as supersaturation adjustment. Secondly, cloud albedo A increases with a change in cloud droplet number concentration (CDNC) roughly as $dA/dCDNC = A(1 - A)/(3CDNC)$, which entails a lower increase in albedo with a higher baseline CDNC (Twomey, 1991; Carslaw et al., 2013). Therefore, an initial state with higher CCN concentrations will be less sensitive to CCN perturbations than an initial state with lower CCN concentrations (Twomey, 1959; Bellouin et al., 2020; Carslaw et al., 2013).

One important, but poorly understood, process for adequately simulating the pre-industrial atmosphere is new particle formation (NPF), i.e., the formation and growth of new particles in the atmosphere which can grow to act as CCN. Roughly speaking, the efficiency of NPF – i.e., how “many” particles are formed per available condensate – in the pre-industrial atmosphere will determine if the secondary aerosol mass is distributed as very few, very large particles or many smaller particles. Especially in a clean atmosphere, this can play a large role for CCN and CDNC concentrations. Over recent years, the understanding of the drivers of NPF has increased significantly due to improved instrumentation and extensive research (Kerminen et al., 2018; Lee et al., 2019). However, adequately capturing NPF in climate models is difficult due to the requirement for computational efficiency combined with the fine scale of the governing processes, in addition to incomplete scientific understanding of the mechanisms involved (Kerminen et al., 2018; Lee et al., 2019).

NPF starts with the formation of a cluster of molecules, which must then activate with respect to the condensing atmospheric vapors and grow into larger sizes (~ 10 nm in

diameter) (Kerminen et al., 2018; Semeniuk and Dastoor, 2018). Due to the Kelvin effect, few gases have low enough volatility to participate in the very first stages of NPF, while as the particles grow, more gases contribute (Semeniuk and Dastoor, 2018). During this growth, the particles are subject to coagulation with larger particles, which constitutes a loss in number concentration (Kerminen et al., 2018). The coagulation sink is approximately proportional to $1/d_p^m$, where d_p is the particle diameter and m is a parameter dependent on the background aerosol concentrations (typically 1.6–1.8) (Lehtinen et al., 2007). It is therefore important for successful NPF that the growth rate (GR) is high enough for the particles to quickly grow to larger sizes when the coagulation sink is lower (Lehtinen et al., 2007). Both Lee et al. (2013) and Olenius and Riipinen (2017) show that omitting explicit modeling of this early aerosol growth and rather parameterizing the survival percentage of particles (e.g., Kerminen and Kulmala, 2002; Lehtinen et al., 2007) leads to significant overestimation of particles. This is mainly because these parameterizations assume steady-state conditions during growth, i.e., the growth rate and coagulation sink are constant, and changes in chemistry, mixing or emissions cannot be taken into account. This assumption is usually not appropriate, especially since growth can take many hours or even days.

The importance of adequately capturing NPF in modeling the pre-industrial atmosphere is illustrated in a study by Gordon et al. (2016), which shows a major reduction (27 %) in estimated forcing from cloud albedo change when including a nucleation pathway from pure biogenic organics. NPF is subject to several constraints which would indicate more efficient NPF in the pre-industrial atmosphere compared to the present day. Firstly, since the pre-existing aerosol concentrations and thus condensation sink will be lower, the gas-phase precursor concentrations are higher *per emission* than in the present-day atmosphere. In other words, if an aerosol precursor species were to have the same emissions and production in the pre-industrial and present-day atmosphere, the pre-industrial atmosphere would have higher gas-phase concentrations because the condensation sink would be lower. Secondly, the coagulation sink of the clusters and newly formed particles is smaller in a cleaner atmosphere (Carslaw et al., 2013; Gordon et al., 2017).

To better capture the early growth of particles from formation to CCN sizes, we have previously implemented a sectional scheme in the aerosol scheme, OsloAero, of the Norwegian Earth System Model (Blichner et al., 2021). We refer to the aerosol scheme with the sectional scheme as OsloAeroSec. OsloAeroSec includes five bins and two condensing species (H_2SO_4 and low-volatility organics) and treats only the growth and loss of particles from formation at 5 nm up to the pre-existing modal aerosol scheme at 39.6 nm diameter, in which climate (cloud–radiation) interactions are considered. See Fig. 1 for an illustration of the scheme. This work was motivated by (1) the smallest mode in the aerosol scheme, OsloAero6, which is quite large (number median di-

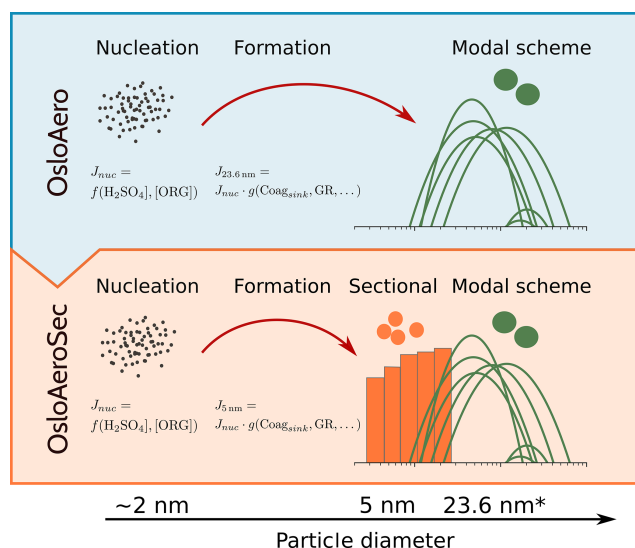


Figure 1. Illustration of changes from OsloAero to OsloAeroSec. In both versions, the nucleation rate is calculated at around 2 nm followed by a calculation of the formation rate (the particles surviving) at 5 and 23.6 nm in OsloAeroSec and OsloAero, respectively, with Lehtinen et al. (2007). In OsloAero, these particles are inserted directly into the modal scheme, while in OsloAeroSec, the particles are inserted into the sectional scheme wherein they can be affected by growth and coagulation over time and space. Finally, the particles in the sectional scheme are moved from the last bin of the sectional scheme to the modal scheme. * 23.6 nm is the number median diameter of the mode the particles from the sectional scheme are moved to, but particles are actually grown to the volume median diameter before they are moved to the modal scheme in order to conserve mass. From Blichner et al. (2021).

ameter 23.6 nm), meaning that growth up to 23.6 nm is parameterized. As mentioned above, this has been shown to lead to significant overestimates of particle formation (Lee et al., 2013; Olenius and Riipinen, 2017). (2) A sectional scheme explicitly grows the particles and does not a priori assume a shape to the size distribution. In this way it is more physically realistic than including, e.g., a nucleation mode. Additionally, the sectional scheme allows for differentiating which organic vapors can contribute to growth from 5 nm upwards compared to from nucleation up to 5 nm.

Our results presented in Blichner et al. (2021) show considerable improvement in the representation of CCN size particles (> 50 nm) compared to observations and significantly reduces the frequent high bias in the original model. This goes in line with Olenius and Riipinen (2017) and Lee et al. (2013). On the other hand, the sectional scheme shows an increase in particle number concentrations in remote areas like the polar regions and the free troposphere.

Motivated by both the improvement to the aerosol scheme and the spatial difference in aerosol formation from the original scheme (remote versus polluted), we investigate the implications of the growth treatment in OsloAeroSec for the

pre-industrial and present-day atmosphere, respectively, especially for the estimated cooling from aerosol–cloud interactions since pre-industrial times.

The cooling effect is commonly quantified by the radiative forcing (RF) or effective radiative forcing (ERF), which are measures of the change in the net radiation into the atmosphere with the addition of a climate forcing agent. RF is, according to the International Panel on Climate Change’s Assessment Report 5 (IPCC AR5) (Boucher et al., 2013) definition, the change in net downwards radiative flux at the tropopause from perturbing the forcing agent, keeping the state variables in the troposphere fixed but allowing the stratosphere to adjust. However, the ERF is in general considered a better indicator of induced surface temperature change because of so-called “rapid adjustments” in the atmospheric column, which may offset or augment the temperature change from the RF, depending on the forcing agent (Bellouin et al., 2020). In this paper, we therefore use the ERF definition as introduced in IPCC AR5, namely the change in top-of-the-atmosphere downwards net flux while allowing adjustments in clouds, temperature, humidity and so on in the atmospheric column but keeping the sea surface temperature fixed.

2 Model description

The Norwegian Earth System Model v2 (NorESM2) (Seland et al., 2020a; Bentsen et al., 2013; Kirkevåg et al., 2013; Iversen et al., 2013) is developed with a basis in the Community Earth System Model (CESM) (Danabasoglu et al., 2020). Firstly, the ocean component, which is not active in these runs since we use fixed sea surface temperature (fSST), is replaced by Bergen Layered Ocean Model (BLOM) (Seland et al., 2020a). Secondly, the atmospheric component, CAM6-Nor, differs from the Community Atmosphere Model v6 (CAM6) in CESM in that its aerosol scheme is replaced by OsloAero6 (Kirkevåg et al., 2018), which we describe briefly below.

In this study we investigate the sensitivities of our sectional scheme for early growth, which was newly implemented into OsloAero6 by Blichner et al. (2021). Both the original aerosol scheme, referred to as OsloAero, and our version with the sectional scheme implemented, referred to as OsloAeroSec, are described in depth in Blichner et al. (2021). We will therefore only give a brief description of the aerosol scheme here.

All runs are done with CAM6-Nor coupled with the Community Land Model v5 (CLM5) in BGC (biogeochemistry) mode and prognostic crop (Lawrence et al., 2019), with prescribed sea ice and sea surface temperatures.

In the following, we start by describing CAM6-Nor in general with the default aerosol scheme, OsloAero, before describing the changes introduced in OsloAeroSec.

2.1 CAM6-Nor

As mentioned earlier, CAM6-Nor shares many characteristics with CAM6 (Bogenschütz et al., 2018), while the aerosol scheme was exchanged for OsloAero, as described below in Sect. 2.1.1. The cloud macrophysics are treated with The Cloud Layers Unified by Binormals (CLUBB; Bogenschütz et al., 2013) model. The microphysics for stratiform and shallow convection clouds are the two-moment bulk from Gettelman and Morrison (2015) (MG2), while the deep convection microphysics are treated with a simplified single-moment representation based on Zhang and McFarlane (1995). The cloud activation of aerosols is done with Abdul-Razzak and Ghan (2000). See Bogenschütz et al. (2018) for more details about the clouds.

2.1.1 OsloAero

OsloAero is often referred to as a “production-tagged” aerosol module, meaning that the model to a large extent keeps track of the processes that each tracer has gone through (e.g., coagulation, condensation). A key difference to other aerosol modules is that it divides the tracers into “process” tracers and “background” tracers. The idea is that the background tracers decide the number concentration, while the process tracers modify the initial size distribution and chemical composition with a look-up table approach (Bentsen et al., 2013; Kirkevåg et al., 2018, 2013; Iversen et al., 2013; Seland et al., 2020a). The background tracers form initial lognormal modes, but after the process tracers are applied, the distribution of the resulting “mixtures” is not necessarily lognormal anymore. This distribution is then used for the optical properties and cloud activation.

The chemistry scheme in NorESM uses the preprocessor MOZART (Emmons et al., 2010) to produce a simplified scheme for sulfur and organic species. The oxidants considered in the model are the hydroxyl radical (OH), ozone (O₃), the nitrate radical (NO₃) and hydroperoxyl (HO₂). These are read from file and interpolated from the monthly mean. The chemistry scheme treats the oxidation of sulfur dioxide (SO₂), dimethyl sulfide (DMS), isoprene and monoterpenes. For a more detailed discussion of the chemistry see Karset et al. (2018), and for a complete overview of reactions and reaction rates, see in particular Table 2 therein.

The aerosol scheme contains three condensing tracers: H₂SO₄ and two organic species, namely SOAG_{LV} and SOAG_{SV}. The H₂SO₄ is produced through oxidation or emitted directly into the atmosphere. The two organic tracers are produced through oxidation of monoterpene and isoprene; each reaction has a certain yield of SOAG_{LV} and SOAG_{SV}. The reactions of isoprene with OH, O₃ and NO₃ all yield 5 % SOAG_{SV}, while monoterpene + OH and monoterpene + NO₃ yield 15 % SOAG_{SV}. Finally, monoterpene reacting with monoterpene + O₃ yields 15 % SOAG_{LV}, thus being the only reaction yielding SOAG_{LV}. The yields used here are simi-

lar to those used in other global models (see, e.g., Tsigaridis et al., 2014; Sporre et al., 2020; Dentener et al., 2006). All these yields are subject to substantial uncertainty (Shrivastava et al., 2017) – see, e.g., Sporre et al. (2020) for an extensive discussion on the sensitivities to these choices.

During condensation these are all treated as non-volatile, but we separate between SOAG_{LV} and SOAG_{SV} because only SOAG_{LV} is considered low-volatility enough to contribute to NPF. In fact, only 50 % of the SOAG_{LV} in each time step is assumed to be low-volatility enough to contribute to nucleation, and we will refer to this fraction of the SOAG_{LV} as extremely low-volatility organic compounds (ELVOCs).

New particle formation is parameterized by using an intermediate concentration of H₂SO₄ and ELVOCs in each time step to calculate a nucleation rate followed by a calculation of how many particles survive growth up to the background mode keeping the particles from NPF (23.6 nm in number median diameter).

The nucleation rate is calculated using Vehkamäki et al. (2002) for binary sulfuric acid–water nucleation and Eq. (18) from Paasonen et al. (2010) to represent boundary layer nucleation.

This survival of particles from nucleation at $d_{\text{nuc}} \approx 2$ nm, the NPF mode, is parameterized (number median diameter $d_{\text{mode}} = 23.6$ nm) by Lehtinen et al. (2007):

$$J_{d_{\text{mode}}} = J_{\text{nuc}} \exp\left(-\gamma d_{\text{nuc}} \frac{\text{CoagS}(d_{\text{nuc}})}{\text{GR}}\right), \quad (1)$$

where $J_{d_{\text{mode}}}$ is the formation rate at d_{mode} , d_{nuc} is the diameter of the nucleated particle, $\text{CoagS}(d_{\text{nuc}})$ is the coagulation sink of the particles [h^{-1}], GR is the growth rate [nm h^{-1}] of the particle (from H₂SO₄ and ELVOCs; calculated using Eq. 21 from Kerminen and Kulmala, 2002), and γ is a function of d_{form} and d_{nuc} :

$$\gamma = \frac{1}{m+1} \left[\left(\frac{d_{\text{form}}}{d_{\text{nuc}}} \right)^{(m+1)} - 1 \right], \quad m = -1.6. \quad (2)$$

2.1.2 OsloAeroSec

We have implemented a sectional scheme for modeling the growth of particles from nucleation up to the mode which keeps the NPF particles in NorESM (number median diameter 23.6 nm). The scheme is described in detail in Blichner et al. (2021). The scheme contains five bin sizes set according to a discrete geometric distribution (Jacobson, 2005, sec.13.3) and two condensing vapors: H₂SO₄ and SOAG_{LV}. The condensation of these species is treated as non-volatile, and after condensation the particles are “grown” (moved) to adjacent bins according to a quasi-stationary structure (Jacobson, 1997, 2005). Coagulation is accounted for both between particles in the sectional scheme and with particles in the modal scheme. When two particles in the sectional scheme coagulate, this contributes to growing the particles, while if they coagulate with particles in the modal scheme,

their mass is added to a process tracer in OsloAero (see Blichner et al., 2021, for more details).

The sectional scheme starts at 5 nm and extends to 39.6 nm, at which the particles are transferred to the NPF mode in the pre-existing aerosol scheme. The sectional scheme extends to the volume median diameter (39.6 nm) rather than the number median diameter (23.6 nm) in order to preserve both number and mass during the transfer between the schemes.

The boundary layer nucleation parameterization has been updated from Paasonen et al. (2010) to Riccobono et al. (2014) and is now

$$J_{\text{nuc}} = A_3[\text{H}_2\text{SO}_4]^2[\text{ELVOC}], \quad (3)$$

where $A_3 = 3.27 \times 10^{-21} \text{ cm}^6 \text{ s}^{-1}$.

Finally, in this version of the model, we have also added improvements to the diurnal variation of the oxidant concentrations, as described below.

2.2 Chemistry: changes to oxidant diurnal variation

The oxidant concentration in CAM6-Nor is read from prescribed 3D monthly mean fields (Seland et al., 2020a) with a diurnal cycle superimposed on OH, HO₂ and NO₃. In the case of OH, this is basically a step function based on before vs. after sunrise, which in turn leads to a step function in the H₂SO₄ concentration and an unrealistic NPF diurnal cycle. In OsloAeroSec, we therefore implemented a simple sine shape on the daily variation in OH to improve the realism of NPF.

2.3 Model versions

In the Results section we compare three different model versions: OsloAero_{def}, OsloAero_{imp} and OsloAeroSec. The first, OsloAero_{def}, is the default model used, e.g., in the CMIP6 simulations, as described in Sect. 2.1.1 above. The second version, OsloAero_{imp}, is the default model but with the same changes to the nucleation scheme and the oxidant diurnal variation as used in OsloAeroSec. The third is with the sectional scheme, OsloAeroSec, as described in Sect. 2.1.2 and by Blichner et al. (2021). This is summarized in Table 1. The motivation for including all these model versions is to be able to distinguish the effect of the sectional scheme from that of the changes in nucleation and oxidants.

3 Simulation setup

All simulations are performed with NorESM2 release 2.0.1 with 1.9° (latitude) × 2.5° (longitude) resolution with 32 height levels from the surface to ~2.2 hPa in hybrid sigma coordinates. The time step is 0.5 h. We use a configuration with an active atmosphere (CAM6-Nor; Seland et al., 2020a) and land component (CLM5-BGC; Lawrence et al., 2019),

while sea ice and sea surface temperatures are read from file. We use the fixed SST method combined with nudging to estimate effective radiative forcing (ERF) from aerosol–cloud interaction, ERF_{aci}, and ERF from aerosol–radiation interactions, ERF_{ari} (Hansen et al., 2005; Forster et al., 2016). This means that we use prescribed SST and sea ice and perturb the anthropogenic aerosol emissions.

We use nudging against model-produced meteorology to constrain the natural variability (Koopman et al., 2012; Zhang et al., 2014; Forster et al., 2016), nudging the horizontal wind components (U , V) and surface pressure with a relaxation time of 6 h (as described in Karset, 2020, Sect. 4.1). Only nudging U , V and surface pressure is preferable over nudging more variables (temperature, humidity, energy fluxes, surface drag) because it allows for rapid adjustments, which should be included in ERF_{aci}. See Karset (2020, Sect. 4.1) for a discussion.

In addition, we use the method proposed by Karset et al. (2018) to estimate the effective radiative forcing; i.e., we use not only the anthropogenic aerosol emissions but also the oxidants from the present-day atmosphere.

To produce the meteorology, we first ran a 7-year simulation (plus 2 years discarded as spin-up), MMET₁₈₅₀, with the default model, OsloAero_{def}. This was done with standard CMIP6 pre-industrial (here meaning 1850) forcing and emissions.

Two simulations were performed with each model version:

- PI – pre-industrial (1850) simulation nudged to MMET₁₈₅₀ and
- PD – simulation with aerosol emissions and oxidant fields from the “present day” (2014) nudged to pre-industrial meteorology (MMET₁₈₅₀).

These are the simulations used to calculate the ERF and which are analyzed in the Results section. Emissions of aerosol and precursors for both the present and pre-industrial are from Hoesly et al. (2018) and van Marle et al. (2017). Oxidant fields are as described in Seland et al. (2020a) from Danabasoglu et al. (2020).

The PI simulations were all initialized from a 2-year simulation with the OsloAero_{def} model version with pre-industrial conditions and free meteorology (SPINUP_PI). Similarly, the PD simulations were all initialized from a 2-year simulation with the OsloAero_{def} model version with free meteorology and pre-industrial conditions but present-day aerosol emissions and oxidant fields (SPINUP_PD). MMET_PI, SPINUP_PI and SPINUP_PD were all initialized from a 30-year simulation with PI configuration.

Table 3 summarizes the model simulations, and Table 2 summarizes the abbreviations for the model versions and configurations.

Table 1. Model version overview.

Simulation	Nucleation parameterization	Oxidant treatment	Early growth treatment
OsloAeroSec	$A_3[\text{H}_2\text{SO}_4]^2 \times [\text{ELVOC}]^a$	Improved diurnal variation	Lehtinen et al. (2007) + sectional scheme
OsloAero _{imp}	$A_3[\text{H}_2\text{SO}_4]^2 \times [\text{ELVOC}]^a$	Improved diurnal variation	Lehtinen et al. (2007)
OsloAero _{def}	$A_1[\text{H}_2\text{SO}_4] + A_2[\text{ELVOC}]^b$	Default diurnal variation	Lehtinen et al. (2007)

$$A_1 = 6.1 \times 10^{-7} \text{ s}^{-1}; A_2 = 3.9 \times 10^{-8} \text{ s}^{-1}; A_3 = 3.27 \times 10^{-21} \text{ cm}^6 \text{ s}^{-1}; {}^a \text{ Riccobono et al. (2014). } {}^b \text{ Paasonen et al. (2010).}$$

Table 2. Abbreviations for model configurations and versions.

	Abbreviation	Description
Forcing configuration	PI	Pre-industrial (1850) run with pre-industrial aerosol emissions and oxidants
	PD	Pre-industrial (1850) run with anthropogenic emissions and oxidant fields from present day (2014)
Model versions	OsloAero _{def}	Run with OsloAero _{def}
	OsloAero _{imp}	Run with OsloAero _{imp}
	OsloAeroSec	Run with OsloAeroSec

4 Terminology

Because we are comparing model versions with and without the sectional scheme, we will only discuss particle number concentrations of particles in the modal OsloAero part of the scheme: that is, excluding the ones still in the sectional scheme. This gives us an apples-to-apples comparison with the original model version. We will use N_a to refer to total aerosol concentration, excluding the particles in the sectional scheme, and N_{NPF} for the subset of these particles originating from NPF. Furthermore, we use $N_{d_1-d_2}$ to refer to the particles with a diameter larger than d_1 but smaller than d_2 . These definitions are summarized in Table 4.

We will use the term *NPF efficiency* or *the efficiency of NPF* to describe model-to-model differences in how many NPF particles are produced with the same emissions (PI or PD). If model versions A and B are both run with the same setup (e.g., pre-industrial emissions) and model A produces more NPF particles than model B, we will say that A has higher NPF efficiency than B.

We use the Ghan (2013) method for calculating ERF_{aci} and ERF_{ari}, meaning that we output the net radiation at the top of the atmosphere, F , and in addition output calls to the radiation scheme with clean (no aerosols), F_{clean} , and clean and clear (no aerosol, no clouds), $F_{\text{clean,clear}}$. Thus, the direct aerosol radiative effect is $\text{DIR}_{\text{Ghan}} = F - F_{\text{clean}}$ and the cloud radiative effect is $\text{CRE} = F_{\text{clean}} - F_{\text{clean,clear}}$. It follows further that $\text{ERF}_{\text{ari}} = \Delta \text{DIR}_{\text{Ghan}} = \Delta(F - F_{\text{clean}})$ and $\text{ERF}_{\text{aci}} = \Delta \text{CRE} = \Delta(F_{\text{clean}} - F_{\text{clean,clear}})$, where Δ signifies the difference between PD and PI.

5 Results and discussion

We will start by presenting globally averaged ERF_{ari} and ERF_{aci} in the model versions and how these relate to PI to PD changes in globally averaged aerosol and cloud properties (Sect. 5.1). Next, in Sect. 5.2, we present a series of hypotheses for the differences in ERF_{ari} and ERF_{aci} between the model versions, which we will use to analyze the results.

In Sect. 5.3, we discuss the PI-to-PD changes on a regional level before discussing the PI and PD simulations separately in Sect. 5.4 and 5.5. We discuss all model versions for which this is helpful to understand the results, but we otherwise focus on OsloAeroSec versus OsloAero_{def} because OsloAero_{def} is the version used in CMIP6.

5.1 Global averages: aerosol number and ERF

5.1.1 Aerosol number

In general, the sectional scheme produces more particles than the original scheme in very pristine environments while producing fewer in areas with high aerosol concentrations (Blichner et al., 2021). This is reflected in the globally averaged profiles of NPF particles, N_{NPF} , for each model version shown in Fig. 2. In the PD simulations, OsloAeroSec mostly has lower N_{NPF} concentrations than the other model versions, surpassing OsloAero_{imp} only above ~ 650 hPa. However, in the cleaner PI atmosphere, OsloAeroSec has N_{NPF} concentrations closer to, or even higher than, the other two schemes. OsloAeroSec has higher N_{NPF} concentrations above ~ 850 and ~ 700 hPa compared to OsloAero_{imp} and OsloAero_{def}, respectively. Close to the surface, where aerosol concentrations in general are higher, OsloAeroSec

Table 3. Description of runs. See Table 2 for abbreviations.

	Simulation name	Model version	Forcing conf.	Initialized from	Meteorology	Years
Meteorology	MMET_PI	OsloAero _{def}	PI	^a	Free meteorology	1–8
Spin-up runs	SPINUP_PI	OsloAero _{def}	PI	^a	Free meteorology	1–2
	SPINUP_PD	OsloAero _{def}	PD	^a	Free meteorology	1–2
PI runs	OsloAero _{def} _PI	OsloAero _{def}	PI	SPINUP_PI	Nudged MMET_PI	(3)4–8 ^b
	OsloAero _{imp} _PI	OsloAero _{imp}	PI	SPINUP_PI	Nudged MMET_PI	(3)4–8 ^b
	OsloAeroSec_PI	OsloAeroSec	PI	SPINUP_PI	Nudged MMET_PI	(3)4–8 ^b
PD runs	OsloAero _{def} _PD	OsloAero _{def}	PD	SPINUP_PD	Nudged MMET_PI	(3)4–8 ^b
	OsloAero _{imp} _PD	OsloAero _{imp}	PD	SPINUP_PD	Nudged MMET_PI	(3)4–8 ^b
	OsloAeroSec_PD	OsloAeroSec	PD	SPINUP_PD	Nudged MMET_PI	(3)4–8 ^b

^a 30-year run with PI emissions. ^b Year 3 is discarded as spin-up, and years 4 to 8 are used in the analysis.

Table 4. Model variable definitions.

Variable name	Definition
N_a	Number of particles excluding those in the sectional scheme
N_{NPF}	Number of particles from NPF excluding those in the sectional scheme
$N_{d_1-d_2}$	Number of particles with diameter d such that $d_1 \leq d \leq d_2$
N_{d_1}	Number of particles with diameter d such that $d_1 \leq d$

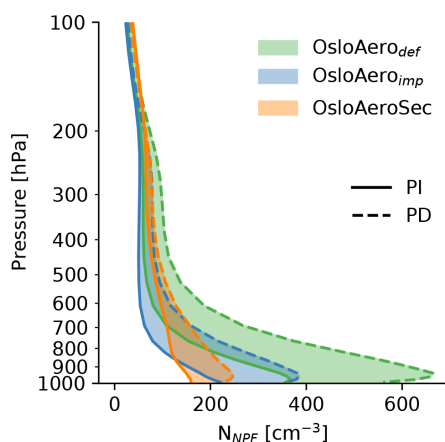


Figure 2. Globally averaged concentration of aerosols from NPF. The solid lines show the concentration in the PI simulation, while the dashed lines show the concentration in the PD. The shading signifies the change in each model. Note that the inter-annual variability in the globally averaged N_{NPF} within each simulation is very low (see Fig. S8 in the Supplement).

has lower N_{NPF} than the other two models, even in the PI simulation.

As we shall explain more in-depth later, these changes in NPF in clean remote versus higher-aerosol-concentration areas are important for ERF_{aci} because the NPF particles are more likely to activate in pristine regions, while they may even act to suppress activation in more polluted regions.

Furthermore, note that even though OsloAero_{imp} is the same as OsloAeroSec, excluding the sectional scheme, the profile is qualitatively different: OsloAeroSec has fewer particles close to the ground and much more further up in the PI atmosphere; see Sect. 5.6.

5.1.2 ERF

The globally averaged ERF_{aci} is significantly influenced by the introduction of the sectional scheme, as seen in Fig. 3, showing total, shortwave and longwave components of ERF_{aci} and ERF_{ari} . ERF_{aci} in OsloAeroSec is significantly ($p < 0.01$) lower than both OsloAero_{def} and OsloAero_{imp}, as found using a two-tailed paired Student's t test on the globally averaged monthly output. The ERF_{aci} is 0.13 W m^{-2} weaker in OsloAeroSec compared to OsloAero_{def}. The ERF_{aci} with OsloAero_{imp} and OsloAero_{def} is roughly the same (difference of 0.01 W m^{-2}). Also, the total radiative effect from aerosols, $ERF_{aci+ari}$, is $\sim 0.1 \text{ W m}^{-2}$ lower in OsloAeroSec compared to both OsloAero_{def} and OsloAero_{imp}. One can further see in Fig. 3 that the difference in the ERF_{aci} between OsloAeroSec and OsloAero_{def} is completely caused by the difference in the SW forcing. Moreover, even though OsloAero_{imp} has roughly the same ERF_{aci} as OsloAero_{def} (not significantly different with $p < 0.05$) it has a significant strengthening ($p < 0.01$) of the forcing in both the SW and LW component that end up canceling each other out in the total forcing. Lastly, the direct effective aerosol forcing, ERF_{ari} , is also shown in Fig. 3, and the direct effect

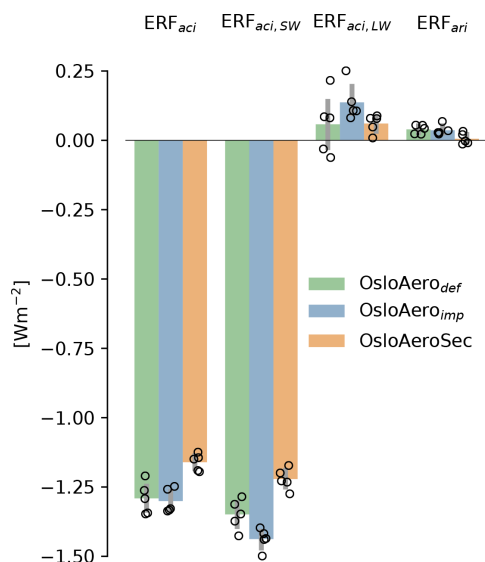


Figure 3. Globally averaged effective radiative forcing (ERF) from aerosols. ERF_{aci} is the ERF from aerosol–cloud interaction, $ERF_{aci,SW}$ and $ERF_{aci,LW}$ are the shortwave and longwave component of ERF_{aci} , and ERF_{ari} is the ERF from aerosol–radiation interaction alone. All are computed in accordance with Ghan (2013). The circles are the averages for each individual year in the 5-year simulations, and the gray bar indicates the 95 % confidence interval of the mean.

is slightly closer to zero with OsloAeroSec than OsloAero_{def} and OsloAero_{imp} ($\sim -0.03 \text{ W m}^{-2}$ smaller than OsloAero_{def} and OsloAero_{imp}, significant with $p < 0.01$). It may seem surprising that both OsloAero_{def} and OsloAero_{imp} have positive ERF_{ari} . Note that we are using Ghan (2013) to calculate ERF_{ari} and that other methods may give a slightly different result. Smith et al. (2020) show comparisons of different estimates of the ERF_{ari} for CMIP6 models and find values similar to ours for NorESM with the Ghan (2013) method, while e.g., the approximate partial radiative perturbation (APRP) method gave a negative ERF_{ari} for the same simulations. The difference between OsloAeroSec and the default model likely originates from OsloAeroSec producing fewer particles than OsloAero_{def} in the PD simulation and thus allowing the remaining particles to grow larger and thus scatter radiation more efficiently (Blichner et al., 2021).

As discussed in the Introduction, ERF_{aci} depends on both the increase in CCN between PI and PD and on the number of CCN in the PI base state. The fewer CCN there are in the base state, the larger the impact of a given increase in CCN will be because the clouds are more susceptible. As OsloAeroSec has much lower particle number concentrations than OsloAero_{def} in the PI, we might expect OsloAeroSec to have fewer CCN and lower CDNC as well as weaker (less negative) $NCRE_{Ghan}$ in the PI. In this case OsloAeroSec would have clouds that are more susceptible to change from PI to PD than OsloAero_{def}. The opposite

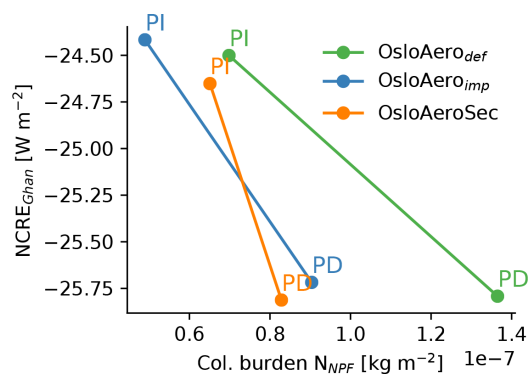


Figure 4. Globally averaged aerosol values of $NCRE_{Ghan}$ (y axis) and column burden of NPF particles (x axis) for the pre-industrial (PI) and present-day (PD) atmosphere. The circles show each annual average and are included to indicate the variability.

is in fact the case, as can be seen in Fig. 4, which relates the column burden of N_{NPF} particle mass (which, due to the technical setup of OsloAero, is proportional to the number) to the net cloud radiative effect ($NCRE_{Ghan}$). While the column burden of N_{NPF} is lower in OsloAeroSec compared to OsloAero_{def}, the $NCRE_{Ghan}$ is stronger (more negative). On the other hand, OsloAero_{imp} has the lowest column burden of N_{NPF} and the weakest $NCRE_{Ghan}$, and it thus follows the logic that a cleaner atmosphere gives a less negative (weaker) $NCRE_{Ghan}$. In the PD simulations, OsloAeroSec has the lowest column burden of N_{NPF} of all the models and approximately the same $NCRE_{Ghan}$ as OsloAero_{def}, while OsloAero_{imp} has a less negative $NCRE_{Ghan}$ than the other two. Since $ERF_{aci} = NCRE_{Ghan,PD} - NCRE_{Ghan,PI}$, it is clear from Fig. 4 that most of the difference between the schemes originates in different $NCRE_{Ghan}$ in the PI simulations: -0.15 and -0.24 W m^{-2} compared to OsloAero_{def} and OsloAero_{imp}, respectively. The difference in the PD simulations partially compensates for this but is considerably smaller: -0.02 and -0.1 W m^{-2} compared to OsloAero_{def} and OsloAero_{imp}, respectively. Furthermore and maybe surprisingly, this plot shows that the change in $NCRE_{Ghan}$ per change in column burden N_{NPF} (i.e., the slope of the line in Fig. 4) is much more negative for OsloAeroSec than for the other two model versions.

5.2 Reasons for differences in ERF_{aci}

From what we have seen so far, it is first of all clear that changes in the PI $NCRE_{Ghan}$ dominate the difference in ERF_{aci} between the models; i.e., the spread in modeled $NCRE_{Ghan}$ between the models is larger in PI than in PD. Secondly, we have seen that, at least in globally averaged properties, more efficient NPF, meaning more particles with the same emissions, does not necessarily lead to a stronger negative $NCRE_{Ghan}$. To explain the somewhat nonintuitive relationship between particle number and $NCRE_{Ghan}$, we

must also consider their geographical distributions with respect to where the NPF particles are likely to activate in clouds and contribute to CDNC. In this section we first outline some important processes and then lay out some hypotheses for the difference in $NCRE_{Ghan}$ with OsloAeroSec compared to the other versions. These will serve to facilitate the rest of the results and discussion.

The cloud droplet activation of particles and resulting CDNC depend on the following factors: (1) the maximum achieved supersaturation (S_{max}) together with the hygroscopicity of the particles decides the activation diameter of each mode, (2) S_{max} depends on the updraft velocity but is also influenced by supersaturation adjustment due to the uptake of water vapor from large(r) particles which activate early during lifting, and finally, (3) the absolute number of particles in each mode which are larger than the activation diameter and thus activate.

Furthermore, note that the number of particles from NPF is strongly negatively correlated with the number median diameter of the modes in the size distribution for both the NPF mode and the larger modes. This is because the total available surface area is larger when there are more NPF particles, which means the available condensate is distributed to more numerous but smaller particles. This leads, as we will show, to NPF inhibiting cloud droplet activation in many regions in the model.

Figure 5 illustrates the effect of changing the NPF efficiency on CDNC in two different environments. For simplicity, let us assume that we are comparing two models with different NPF efficiency: model A with high NPF efficiency and model B with low NPF efficiency. As noted above, model A will have more numerous but smaller particles (A1 and A2 in Fig. 5), while model B will have fewer but larger particles (B1 and B2 in Fig. 5). Furthermore, we will consider two different environments. Environment 1 has a small activation diameter because, e.g., there are few large particles (no early activation) or the updraft is strong (A1 and B1 in Fig. 5). Environment 2 has a large activation diameter because, e.g., it has high emissions of large primary particles which activate early and limit the maximum supersaturation (A2 and B2 in Fig. 5). In this simplification we assume that the activation diameter does not change between model A and B. This is not strictly true, but it is a good assumption because the inter-model changes in S_{max} (Fig. S21 in the Supplement) and hygroscopicity (Fig. S26 in the Supplement) are small and do not dominate the response in terms of CDNC.

We start with environment 1 wherein the activation diameter is small (e.g., Antarctica). This is illustrated by the two size distributions, A1 and B1, on the top in Fig. 5. In this environment model A (high NPF efficiency, A1) will result in higher cloud droplet activation and higher CDNC than model B (low NPF efficiency, B1). This is because a considerable fraction of the small NPF-mode particles activate, and thus the decrease in the size of the larger particles does not matter.

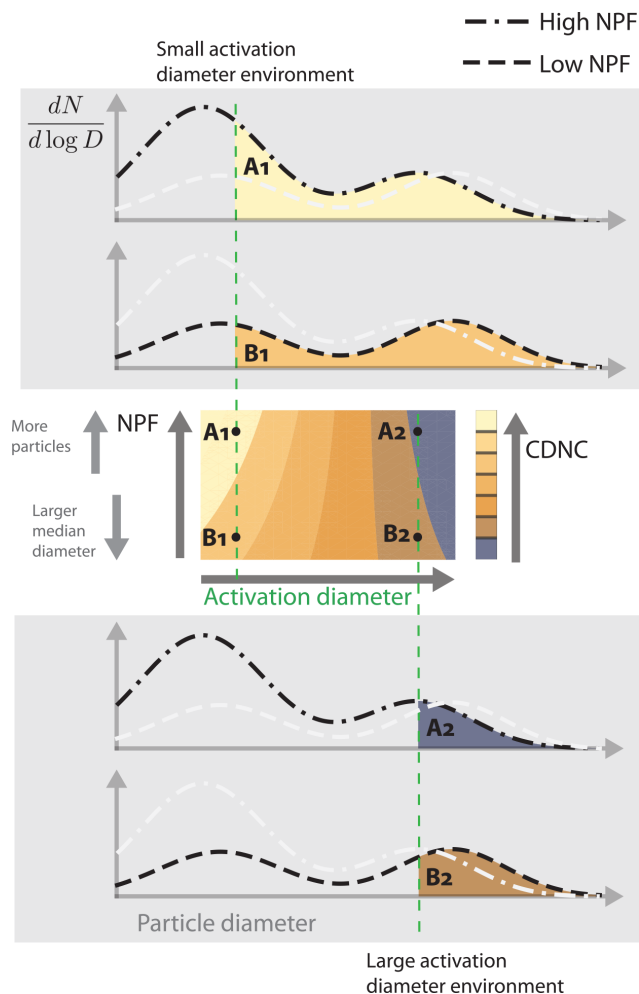


Figure 5. Schematic illustrating the influence of NPF on cloud droplet activation and CDNC. The top and bottom panel shows what happens to activation in two different environments (1 and 2) and for two models: one model with high NPF efficiency (A) and one with low NPF efficiency (B). Let us first consider environment 1 (top panels): here the activation diameter is small (either due to strong updrafts, few large particles or high hygroscopicity), and particles all the way down to the mode holding the NPF particles (\sim Aitken mode) activate. In this environment model A will activate more particles than model B and have higher CDNC. Next let us consider environment 2 (bottom panels): here the activation diameter is large (due to weak updrafts, supersaturation adjustment due to larger particles or hygroscopicity), and only the largest particles activate. Here model B will activate more particles than model A because the size of the larger particles is what dominates.

Next we consider environment 2 wherein the activation diameter is large (e.g., a polluted area like China). This is illustrated by the two size distributions, A2 and B2, at the bottom of Fig. 5. In this environment model A with high NPF efficiency (A2) will result in lower cloud droplet activation and lower CDNC than model B with a low NPF efficiency (B2). This is because the change in the diameter of the larger par-

ticles is the only thing which matters for activation, since the smaller particles will not activate anyway.

In this simplified thought example, we can say that in environment 1 (small activation diameter), NPF enhances cloud droplet activation, while in environment 2 (large activation diameter) NPF inhibits cloud droplet activation.

With all this in mind, we can lay out some plausible hypotheses that might contribute to a weaker ERF_{aci} in OsloAeroSec compared to the other model versions.

1. *Smaller $\Delta_{PD-PI}N_a$.* The difference in ERF_{aci} is due to a smaller change in number concentration between PI and PD in OsloAeroSec than the other model versions.
2. *Higher N_a in PI.* OsloAeroSec produces more particles under PI conditions, and therefore the clouds are less susceptible to increased anthropogenic emissions.
3. *Higher activation in PI.* The number of particles that actually act as CCN and activate is higher with OsloAeroSec than the other model versions in the PI simulations, leading to a higher baseline CDNC. This is due to
 - (a) more efficient NPF in remote regions where NPF enhances activation and
 - (b) less efficient NPF in regions where NPF inhibits activation (only larger particles activate).
4. *Lower activation in PD.* The number of particles that actually act as CCN and activate is lower with OsloAeroSec than the other model versions with PD emissions, leading to a weaker ERF_{aci} . This is
 - (a) due to lower NPF efficiency in regions where NPF enhances activation.

Hypothesis 2 has already been partly disproven because in terms of global averages, OsloAero_{def} has higher particle number concentrations than OsloAeroSec all the way up to approximately 700 hPa (with most of the liquid clouds being below this level).

5.3 Pre-industrial to present-day changes

We start by considering Hypothesis 1 and how the PI-to-PD change looks on a regional level in OsloAeroSec versus OsloAero_{def}.

This is shown in Fig. 6 where the first row is the change between PD and PI (Δ_{PD-PI}) for OsloAeroSec and the two subsequent rows show the difference to this first quantity, Δ_{PD-PI} , between the model versions ($\Delta_{PD-PI}(\text{OsloAeroSec})$ minus $\Delta_{PD-PI}(\text{OsloAero}_{imp})$ and $\Delta_{PD-PI}(\text{OsloAeroSec})$ minus $\Delta_{PD-PI}(\text{OsloAero}_{def})$, denoted $\Delta\Delta_{PD-PI}$). The first column, showing the near-surface averaged N_{NPF} , shows that, as expected, most of the PI-to-PD change happens in the Northern Hemisphere. This is consistent with the major anthropogenic emission sources being located here. Over ocean

regions in the Southern Hemisphere, there is even a small decrease in NPF particles in many places. Comparing to OsloAero_{def} (row 3) we see that OsloAeroSec has a smaller increase in N_{NPF} from PI to PD, except in the South Pacific and over the Amazon. High-pollution areas over land especially stand out as strongly negative. Note that the first column in Fig. S15 in the Supplement shows the same but for zonal averages and underlines the fact that $\Delta_{PD-PI}N_{NPF}$ is higher in OsloAero_{def} than OsloAeroSec all through the atmospheric column.

The second column shows the change in cloud droplet number concentration at cloud top (CDNC(CT)). Again the first row shows $\Delta_{PD-PI}CDNC(CT)$, which, as expected, shows an increase – in particular in the Northern Hemisphere. Comparing OsloAeroSec to OsloAero_{def} (row 3), however, the first thing that stands out is that, somewhat surprisingly, $\Delta\Delta_{PD-PI}CDNC(CT)$ is positive over polluted regions, meaning that the PI-to-PD increase in CDNC(CT) is stronger with OsloAeroSec than with OsloAero_{def}, in spite of N_{NPF} increasing less with OsloAeroSec. In other words, in these regions we are in the bottom panel of Fig. 5, where more particles are added with OsloAero_{def} than OsloAeroSec, but fewer of these extra particles are activating into cloud droplets. Meanwhile, in more remote regions, like the North Pacific and the Arctic, we are in the top panel of Fig. 5 and CDNC(CT) increases less with OsloAeroSec than OsloAero_{def}, following the more expected logic that a smaller increase in particle number leads to a smaller increase in cloud droplets from PI to PD.

Finally, the last column shows the ERF_{aci} . Here we see (first row, c) that the ERF_{aci} is strongly negative over the North Pacific as well as over China and India. The difference in ERF_{aci} between the models shows that the remote Pacific dominates in making ERF_{aci} more strongly negative in OsloAero_{def} than in OsloAeroSec. Even though the increase in CDNC(CT) from PI to PD is stronger in polluted regions with OsloAeroSec, these regions seem to have reached saturation with respect to changing albedo, and the ERF_{aci} changes little between the model versions.

To summarize with regard to Hypothesis 1: the change in particle number between PI and PD is indeed smaller with OsloAeroSec than the other model versions, but this can only explain the change in CDNC in remote regions (North Pacific, Siberia, etc.). Furthermore, as mentioned earlier, we need to consider the influence of the baseline aerosol state in PI and not just the change between PI and PD.

5.4 The pre-industrial atmosphere: model-to-model differences

To consider Hypothesis 3, “higher activation in PI”, we now consider differences between OsloAeroSec and the default model versions in the PI separately from PD (covered in the next section).

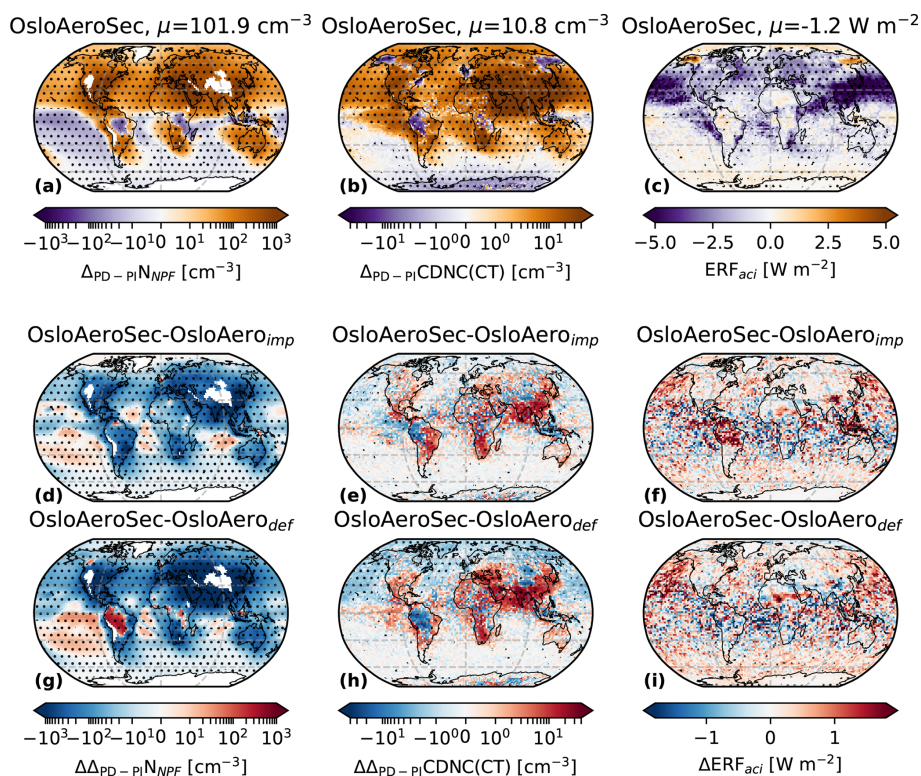


Figure 6. Annual average values of near-surface N_{NPF} concentrations (a, d and g), cloud-top droplet number concentration (CDNC(CT), b, e and h) and $NCRE_{Ghan}$ (c, f and i). The top panel shows the PD – PI for OsloAeroSec, while the second and third rows show the change in this value (PD – PI) from OsloAero_{imp} (second row) and OsloAero_{def} (third row) to OsloAeroSec. The N_{NPF} values are averaged up to 850 hPa and weighted by the pressure thickness of each grid cell. Dots are included in the plots to indicate where the difference between the two models is significant with a two-tailed paired Student's t test with 95 % confidence.

Figure 7 shows the near-surface concentration of N_{NPF} in the PI simulation (left column) for OsloAeroSec (Fig. 7a) and the relative difference in this value between the model versions (Fig. 7b and c). We see that compared to OsloAero_{def}, N_{NPF} is lower in OsloAeroSec almost everywhere in PI. However, as seen in Fig. 8c, showing the zonally averaged difference, this decrease with OsloAeroSec is mostly confined to the near-surface areas. The decrease in N_{NPF} with OsloAeroSec near the surface switches to an increase higher up in the atmosphere.

5.4.1 Cloud properties

OsloAeroSec has a higher cloud droplet number concentration at cloud top (CDNC(CT)) than OsloAero_{def} in most of the PI atmosphere, as can be seen in Fig. 9a. This is despite the fact that OsloAeroSec has lower N_{NPF} concentrations in most near-surface areas compared to OsloAero_{def}. We must therefore investigate what happens to the size distribution rather than just the absolute number. Figure 9c, e and g show the OsloAeroSec-to-OsloAero_{def} difference in number concentrations of N_{100} , N_{150} and N_{200} . The N_{100} concentration (Fig. 9c) is lower in OsloAeroSec than OsloAero_{def} most places in the PI atmosphere, while N_{150} (Fig. 9e) and N_{200}

(Fig. 9g) are higher. This follows the mechanism explained in Sect. 5.2 that lower NPF efficiency in OsloAeroSec leads to fewer but larger particles. The higher concentrations in OsloAeroSec of, e.g., N_{200} come from the modes shifting to higher median diameters when the number of NPF particles is lower. There is also a good correspondence between the difference in N_{150} and/or N_{200} and the difference in CDNC in most areas in the atmosphere. Note, for example, the Amazon area, where much lower concentrations of N_{100} (and NPF efficiency) are associated with much higher concentrations of N_{200} but not N_{150} . That the CDNC is higher here tells us that the activation diameter is probably usually between 150 and 200 nm. Additionally, the supersaturation is higher in OsloAeroSec due to fewer particles that compete for the water vapor (see Fig. S20 in the Supplement), which has a small positive impact on the number of particles which activate.

To further investigate these relationships between changes in N_d and CDNC in the PI simulations, we compute the correlation between $\Delta CDNC$ and ΔN_d , where Δ signifies the difference between OsloAeroSec and OsloAero_{def}. First we compute the correlation between $\Delta CDNC$ and ΔN_{NPF} over time and longitude, as shown in Fig. 10c. This reveals

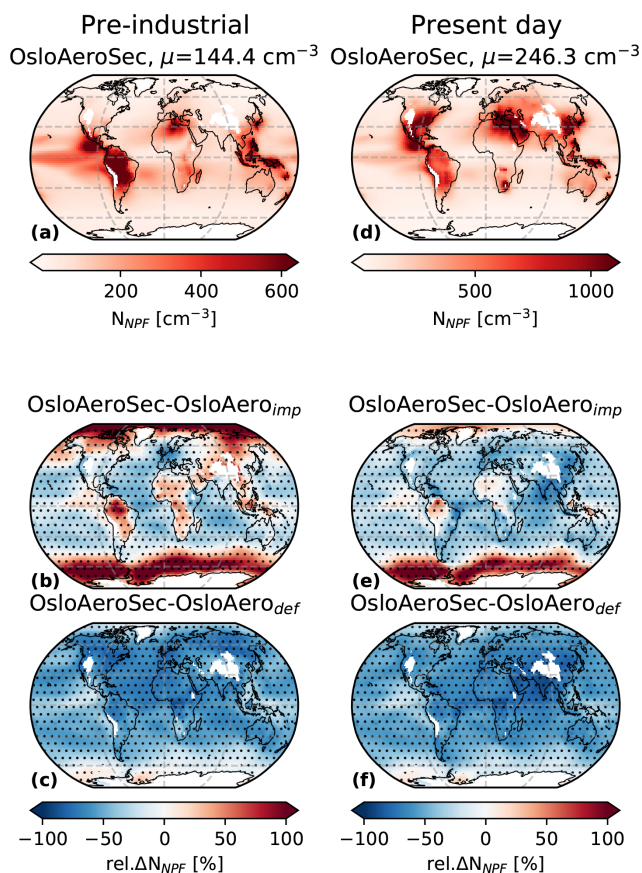


Figure 7. (a, d) Annual average near-surface N_{NPF} concentrations for OsloAeroSec for PI (a) and PD (d). Rows 2–3: the relative difference of OsloAeroSec to OsloAero_{imp} (b, e) and OsloAero_{def} (c, f) for PI (b, c) and PD (e, f). All values are averaged up to 850 hPa and weighted by the pressure thickness of each grid cell. Dots are included in the plots to indicate where the difference between the two models is significant with a two-tailed paired Student's t test with 95 % confidence.

Table 5. Region overview. These regions are used to create vertical average profiles.

Region name	Latitudes	Longitudes
Global	All	All
Antarctic	60–90° S	180° W–180° E
Pacific S	30–60° N	170° E–120° W
Pacific N	60–30° S	170° E–140° W

that close to the surface, ΔCDNC and ΔN_{NPF} are mostly negatively correlated, indicating that in these areas NPF inhibits activation. In remote regions, like the Southern Ocean or high in the free troposphere, there is a positive correlation between ΔN_{NPF} and ΔCDNC , indicating that we are in an NPF-enhanced activation regime and relevant parts of the NPF-mode particles activate.

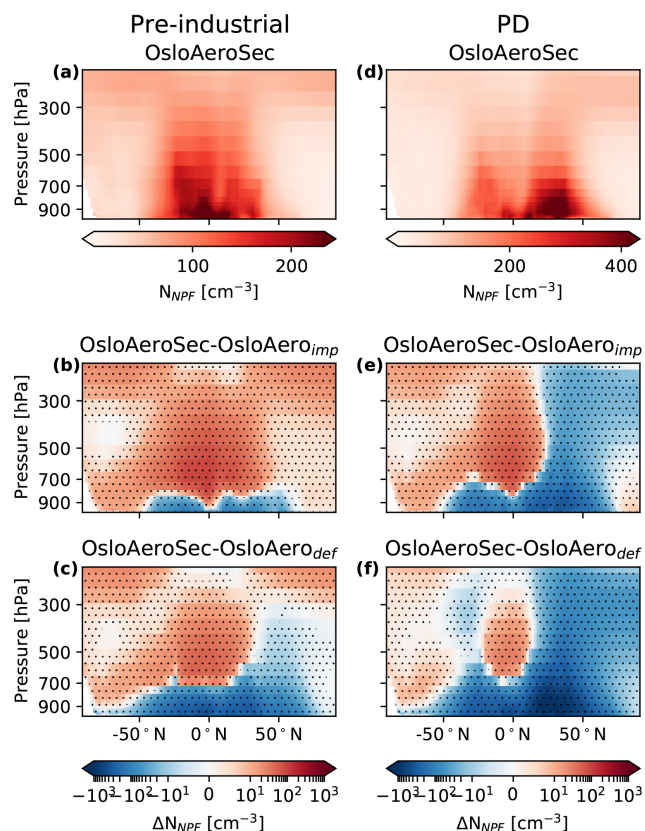


Figure 8. (a, d) Zonally and annually averaged concentrations of N_{NPF} for OsloAeroSec for PI (a) and PD (d). Rows 2–3: the absolute difference of OsloAeroSec to OsloAero_{imp} (b, e) and OsloAero_{def} (c, f) for PI (b, c) and PD (e, f). Dots are included in the plots to indicate where the difference between the two models is significant with a two-tailed paired Student's t test with 95 % confidence.

Second, we compute the correlations between ΔCDNC and ΔN_{50} , ΔN_{100} , ΔN_{150} , ΔN_{200} and ΔN_{250} for different regions (see Table 5 for definitions) at different heights. These relationships for the PI simulations are shown in Fig. 11, column 1. If ΔCDNC clearly correlates with the change in the concentration of particles above some diameter d , N_d , this indicates that these particle sizes are relevant for cloud droplet activation in the region. On the other hand, if there is a negative correlation, this indicates that the particles are too small to activate.

Globally, Fig. 11a shows that CDNC correlates most strongly with N_{200} and N_{250} close to the surface, with an anticorrelation with N_{50} and N_{100} . The sign of the correlations switches at around 600 hPa. In the relatively clean Antarctic (here defined as below 60° S), the correlation is positive with the smaller particles, i.e., N_{50} , throughout the atmosphere. This indicates that NPF enhances activation in Antarctica and that the number of particles dominates rather than the size of the particles. Figure 11e and g show the South and North Pacific and are included because they show opposite sign in

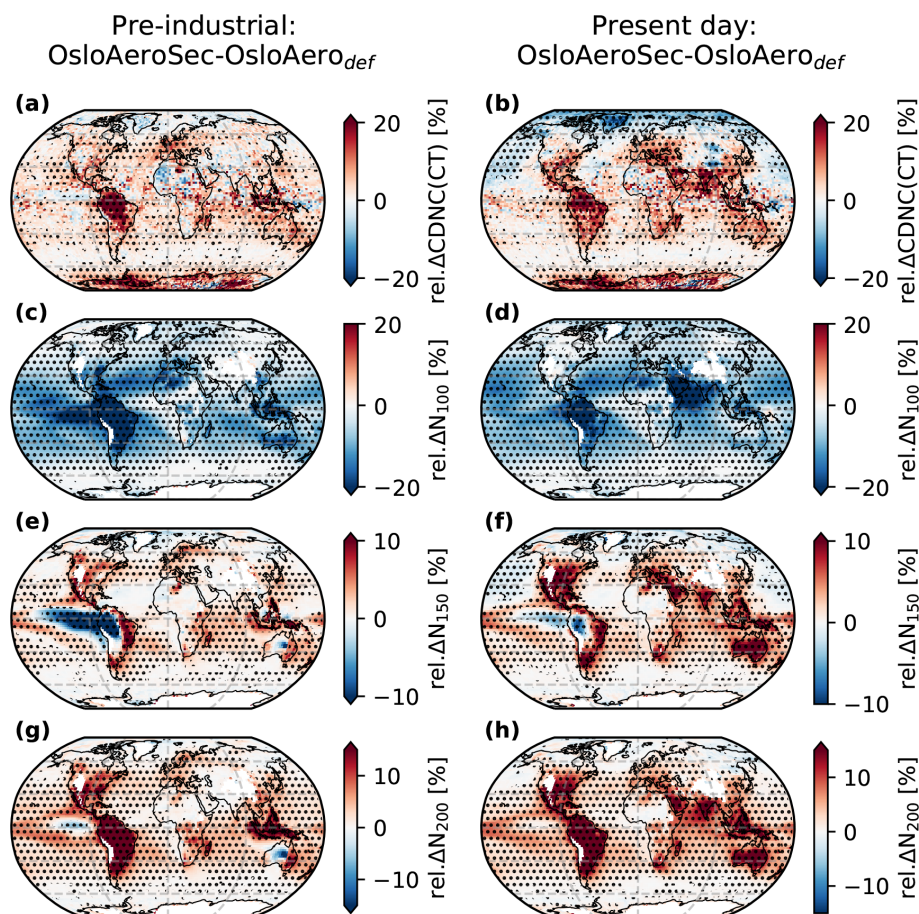


Figure 9. (a, b) Relative difference in annual average cloud-top cloud droplet number concentrations (CDNC(CT)) at cloud top between OsloAeroSec and OsloAero_{def}. Row 2–3: difference in average particle number concentration for particles larger than 100 nm (c, d), 150 nm (e, f) and 200 nm (g, h). The left column shows the difference for the pre-industrial atmosphere, and the right column shows the difference for the present-day atmosphere. The average particle concentrations are calculated by averaging up to 850 hPa and averaging by pressure difference. Dots are included in the plots to indicate where the difference between the two models is significant with a two-tailed paired Student's *t* test with 95 % confidence.

CDNC for the PD simulations; we will discuss them further in the next section. In the PI simulations, however, the South Pacific shows a clear correlation with the larger particles (diameters larger than 150, 200 and 250), while in the North Pacific, the correlation is closer to zero or insignificant.

5.4.2 Summary Hypothesis 3: higher activation in the pre-industrial atmosphere

We do indeed see higher aerosol activation and higher CDNC with OsloAeroSec in the PI simulations. This is due to a combination of two things: (1) in pristine areas, NPF particles are likely to activate and lead to higher CDNC – i.e., NPF enhances activation. In these areas OsloAeroSec in general produces more NPF particles than OsloAero_{def}, and thus CDNC increases. (2) In areas with higher aerosol number concentrations, NPF particles are unlikely to activate and NPF inhibits cloud droplet activation due to reducing the size of the larger

particles. In these regions, OsloAeroSec in general produces fewer NPF particles than OsloAero_{def}, and thus CDNC increases.

5.5 The present-day atmosphere: model-to-model differences

We now move to consider differences in the PD simulations between OsloAeroSec and OsloAero_{def} and will discuss Hypothesis 4, “lower activation in PD”.

With PI emissions, there are large regions, especially at higher altitudes, for which OsloAeroSec produced more NPF particles than the other model versions. With PD emissions, these areas shrink, as the atmosphere becomes less pristine overall. This is seen in Fig. 7d–f (near-surface average) and Fig. 8d–f (zonal average). Furthermore, it is interesting to see the impact of emissions in the Northern Hemisphere versus the Southern Hemisphere in the PD simula-

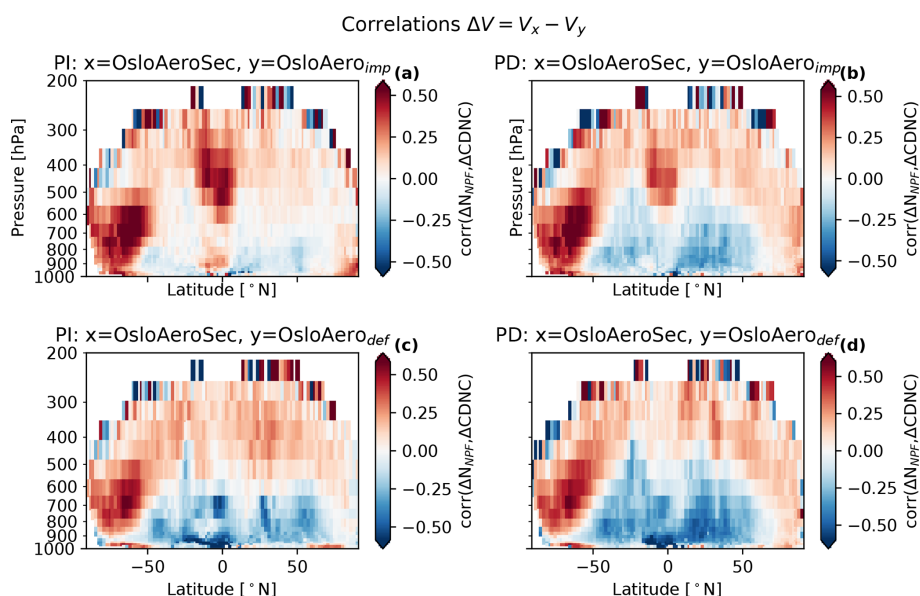


Figure 10. Correlations between the absolute difference in CDNC and the absolute difference in N_{NPF} between the model versions, calculated from monthly mean files over time and longitude. The correlations from the difference between OsloAeroSec and OsloAero_{imp} are shown in panels (a) and (b). The correlations from the difference between OsloAeroSec and OsloAero_{def} are shown in panels (c) and (d). The correlations in the PI simulations are shown to the left, and the ones for the PD simulations are shown to the right.

tions. In the Northern Hemisphere, OsloAeroSec produces fewer particles than the other model versions at most heights and latitudes, while the opposite is the case for the Southern Hemisphere. This is likely due to a combination of much higher emissions and more vertical mixing in the Northern than Southern Hemisphere. In other words, larger parts of the Northern Hemisphere pass into a pollution level regime wherein the sectional scheme produces fewer particles than the others.

5.5.1 Cloud properties

Figure 9b shows the difference in CDNC(CT) between OsloAeroSec and OsloAero_{def} in the PD simulations. The Southern Hemisphere resembles the difference in PI (Fig. 9a) with a widespread increase in CDNC. In the middle to high northern latitudes, on the other hand, CDNC is lower in OsloAeroSec than in OsloAero_{def}, which is opposite to the PI simulations. In these last pristine northern regions, more NPF particles in OsloAero_{def} indeed seem to lead to higher CDNC than in OsloAeroSec.

Let us again consider the model-to-model difference in size distribution. Figure 9d, f and h show ΔN_{100} , ΔN_{150} and ΔN_{200} . Here we see that the pristine Northern Hemisphere ΔCDNC most resembles the change in N_{100} , while in the Southern Hemisphere, ΔCDNC more closely resembles that of the larger particles (N_{150} and N_{200}). Note how the polluted regions in the PD simulations around India and China have higher concentrations of N_{200} and N_{150} in OsloAeroSec than OsloAero_{def} and corresponding higher CDNC. In these pol-

luted regions, NPF in general inhibits cloud droplet activation because the activation diameter is large (bottom panel in Fig. 5). This is because there are many large particles which activate early and act as a sink for water vapor, thus reducing S_{max} and increasing the activation diameter (see Fig. S20b). On the other hand, the decreases in CDNC in OsloAeroSec compared to OsloAero_{def} in the PD northern high latitudes correspond better to the change in the smaller particles, N_{100} and partially N_{150} . This indicates that in these regions NPF enhances cloud droplet activation due to a smaller activation diameter (top panel in Fig. 5). Note that this is different in the PI and PD simulations: in the PD simulations, the CDNC goes down with OsloAeroSec in the northern high latitudes; in the PI it goes up. The reason for this is that the activation diameter depends on both the maximum supersaturation and the hygroscopicity. The hygroscopicity of the particles almost doubles from the PI to the PD due to increased sulfate emissions (see Fig. S26). More hygroscopic particles in the PD simulations can then activate at smaller diameters (given the same S_{max}). The regions where CDNC is enhanced by NPF thus spread in the pristine northern latitudes, favoring cloud droplet activation in OsloAero_{def} over OsloAeroSec. Mark that the difference in hygroscopicity is large between the PI and PD simulations (again, see Fig. S26) but small ($\sim 5\%$) between the different model versions.

It is thus clear that hygroscopicity plays a role, but only in terms of making the effect of NPF particles different in the PI and PD simulations; with PD emissions the NPF particles are more likely to activate. In other words, because hygroscopicity increases in PD, the areas where NPF enhances cloud

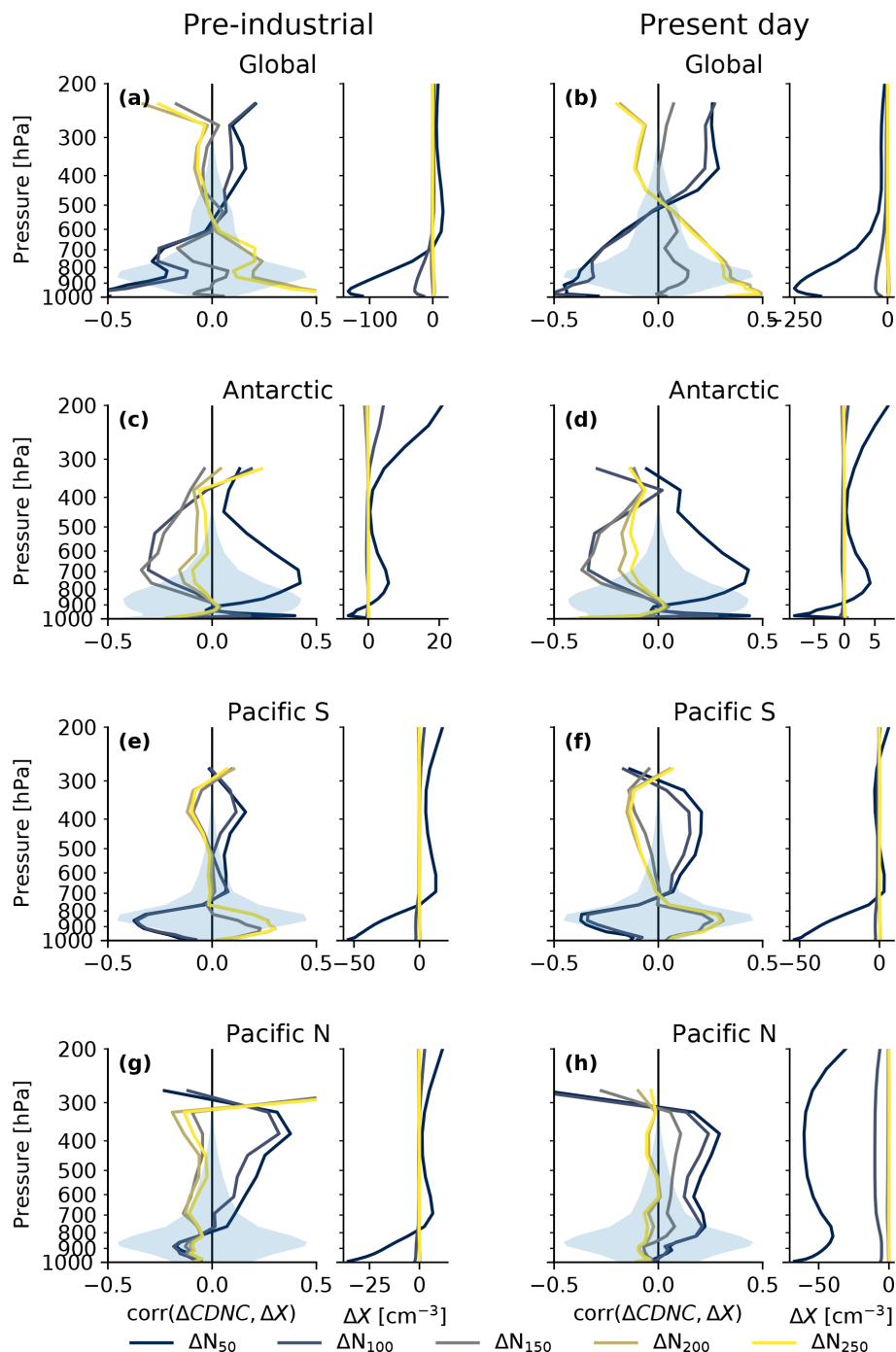


Figure 11. Left side of each panel: correlations by pressure level between the absolute difference between OsloAero_{def} and OsloAeroSec in cloud droplet number concentration ($\Delta CDNC$) and the absolute difference in the number of particles with diameters above 50, 100, 150, 200 and 250 nm for different regions. The blue shading signifies the fractional occurrence of liquid cloud and is included to give an idea of where the aerosols may actually have a noticeable impact on clouds. The right side of each panel shows the change in the aerosol concentration for the relevant region. See Table 5 for definitions of regions.

activation expand in the PD Northern Hemisphere compared to the pre-industrial atmosphere.

Let us again consider the correlations between $\Delta CDNC$ and N_{NPF} , ΔN_{50} , ΔN_{100} , ΔN_{150} , ΔN_{200} and ΔN_{250} for different regions, as shown for the PD atmosphere in Figs. 10 and 11b, d, f and h.

Globally, the correlation of $\Delta CDNC$ with the change in larger particles is more pronounced in the PD than the PI simulations (Figs. 11b and 10d), possibly indicating a stronger supersaturation adjustment (reduced S_{max}), with more polluted PD emission conditions leading to a higher activation diameter.

Furthermore, we investigate the North and South Pacific separately in Fig. 11e–h because these two show opposite sign in the PD simulations: in the North Pacific, OsloAeroSec has lower CDNC than OsloAero_{def}, while in the South Pacific OsloAeroSec has higher CDNC (see Fig. S9b in the Supplement). In the South Pacific (Fig. 11e and f), the CDNC correlates best with the larger particles (diameter above 150 nm) in both PI and PD. In the North Pacific, on the other hand, the correlation is not clear for any particle number in the PI (Fig. 11g) and slightly positive for the smaller particle sizes in PD (Fig. 11h). The likely cause for the difference between the two cases is that (1) the South Pacific has higher concentrations of larger sea salt particles than the North Pacific (not shown), which can limit the maximum supersaturation and thus lead to a higher activation diameter; and (2) as mentioned above, the sulfate emissions are much higher in the PD Northern Hemisphere, leading to more hygroscopic particles and a lower activation diameter. In the South Pacific, we are therefore at the bottom panel of the sketch in Fig. 5, while in the North Pacific, we are more on the top panel. Note again that the hygroscopicity between the model versions with the same emissions (either with PI or PD emissions) changes very little (Fig. S26), which is why we only discuss changes between the PI and PD.

5.5.2 Summary Hypothesis 4: lower activation in the present-day atmosphere

The discussion above shows that regionally, lower cloud droplet activation and CDNC with OsloAeroSec in the PD simulations does indeed play a role in reducing the ERF_{aci} in the pristine high northern latitudes and the North Pacific. Here the CDNC is lower with OsloAeroSec than OsloAero_{def}, and thus OsloAero_{def} has a stronger negative cloud radiative effect in the PD simulations. On the other hand, cloud droplet activation and CDNC in more polluted regions are higher with OsloAeroSec than OsloAero_{def} (see Fig. 9b) in the PD simulations. This, however, does not have as big an impact on radiation (see, e.g., Fig. S16 in the Supplement) firstly because these areas are mostly continental and the cloud radiative effect is larger over dark ocean surfaces (e.g., the North Pacific) and secondly because the CDNC is already high in these regions with

OsloAero_{def}, and thus the clouds are less susceptible to the increase in OsloAeroSec (see the Introduction for a description of this effect). Furthermore, we have found that hygroscopicity changes from PI to PD play a role by reducing the activation diameter and making NPF particles more likely to activate in the PD simulations compared to the PI. This means that the areas where NPF enhances cloud droplet activation expand, and thus there are larger areas where OsloAero_{def} has higher CDNC than OsloAeroSec. Both these factors result in a lower CDNC in the high northern latitudes with OsloAeroSec and a corresponding lower magnitude in $NCRE_{Ghan}$.

5.6 Comparison to OsloAero_{imp}

We have mostly focused on the comparison of OsloAeroSec to OsloAero_{def} in the above section, but there are important points to take away from comparing OsloAeroSec to OsloAero_{imp} as well. Note that OsloAero_{imp} has the same updates to oxidants and nucleation rate as OsloAeroSec, but it does not have the sectional scheme. Also, remember that OsloAero_{imp} has much lower NPF efficiency than OsloAero_{def}, but compared to OsloAeroSec it is more similar; this depends on the region. In general OsloAeroSec produces more NPF particles in pristine regions, while OsloAero_{imp} produces more particles in regions with higher aerosol concentrations.

When comparing only OsloAeroSec and OsloAero_{def}, it is not possible to separate the effect that increased NPF efficiency in remote regions has from decreased NPF efficiency in high-aerosol regions with respect to the ERF_{aci} . It is perhaps tempting to think that the reduction in NPF efficiency alone is responsible for the overall effect and that the increase in NPF efficiency in remote regions is negligible. If so, any scheme which reduced NPF efficiency would have the same effect. The OsloAero_{imp} simulation, however, represents exactly such a scheme which reduces the NPF efficiency compared to OsloAero_{def}, with roughly the same amount as OsloAeroSec, though without the increases in NPF efficiency in remote regions. However, OsloAero_{imp} does not weaken ERF_{aci} like OsloAeroSec does but rather slightly strengthens it.

5.7 Summary of hypotheses

We now summarize and relate the results back to the hypotheses presented in Sect. 5.2.

1. *Smaller $\Delta_{PD-PI}N_a$.* While it is true that N_a increases less from PI to PD with OsloAeroSec than OsloAero_{def} (and OsloAero_{imp}), this can only explain the results in remote regions. Furthermore, OsloAero_{imp} offers as a counter-argument against this hypothesis: it also has a lower PD-to-PI change in N_a ($\Delta_{PD-PI}N_a$) than OsloAero_{def}, but contrary to OsloAeroSec, OsloAero_{imp} has a stronger negative ERF_{aci} than OsloAero_{def}. In

sum, this hypothesis does not explain the differences in ERF_{aci} well.

2. *Higher N_a in PI.* OsloAeroSec mostly produces fewer particles than OsloAero_{def} in the PI simulations, and this is thus only true in remote regions. This hypothesis can therefore not explain the resulting ERF_{aci} .
3. *Higher cloud droplet activation in PI.* We found that OsloAeroSec has higher CDNC than the other model versions in the PI simulations due to more efficient NPF in remote regions where NPF enhances cloud droplet activation (small activation diameter) and due to less efficient NPF in regions where NPF inhibits cloud droplet activation (large activation diameter). In these last areas, OsloAeroSec indeed has a higher concentration of larger particles than OsloAero_{def} and OsloAero_{imp} due to the condensate being distributed to fewer particles in OsloAeroSec. This hypothesis therefore explains the part of the change in ERF_{aci} originating from the difference in $NCRE_{Ghan}$ in the PI simulations well.
4. *Lower cloud droplet activation in PD.* We found this hypothesis to play an important role in the northern high latitudes, especially the North Pacific, where sulfate emissions are high in the PD simulations. Due to higher hygroscopicity in the PD simulations compared to the PI, the NPF particles are more likely to activate (smaller activation diameter), and thus the number of particles (which is lower in OsloAeroSec) is more important than the particles sizes. This hypothesis is therefore important to explain the changes in the PD simulations.

Additionally, after the analysis of the results, we may add two more explanations.

5. *Hygroscopicity.* As explained for Hypothesis 4 above, the change in hygroscopicity from PI to PD results in larger areas in the northern pristine latitudes having an NPF-enhanced cloud droplet activation regime in the PD simulations compared to the PI. This results in stronger $NCRE_{Ghan}$ with OsloAero_{def} than OsloAeroSec in the PD simulations, which further leads to a stronger ERF_{aci} in OsloAero_{def} than OsloAeroSec.
6. *Regional differences.* The comparison with OsloAero_{def} shows that regional differences in NPF matter significantly. For reasons discussed above, OsloAeroSec gives higher CDNC in the PI simulation in regions with susceptible clouds and large ERF_{aci} , which dominates the global average.

6 Implications and discussion

The results in this paper go in line with previous work, which shows that the ERF_{aci} is sensitive to the PI aerosol charac-

teristics (e.g., Carslaw et al., 2013) and that changes in the NPF parameterization can highly influence ERF_{aci} (e.g., Gordon et al., 2016). However, the reduction in ERF_{aci} found with OsloAeroSec in our simulations is not a result of increased NPF under PI conditions alone. Rather, the increase in CDNC and $NCRE_{Ghan}$ in the PI simulation originates from increased NPF efficiency at locations where the NPF enhances cloud droplet activation and decreased NPF efficiency where NPF inhibits particle activation. Additionally, we find that the modeled increase in hygroscopicity from PI to PD from increased sulfate emissions results in a lower activation diameter, and thus more of the NPF particles contribute to CDNC.

The effect of NPF inhibition on cloud droplet activation was also found by Sullivan et al. (2018); they modeled the NPF effect on clouds over the midwestern USA using WRF-Chem v3.6.1 and using a 20-bin sectional aerosol scheme (Model for Simulating Aerosol Interactions and Chemistry, MOSAIC). As in this study, they find that the growth of larger particles is inhibited by the increased condensation sink from the NPF particles. That fact that the same effect is seen in simulations with a completely differently structured aerosol model shows it to be unlikely that this is an artifact of the OsloAero model. However, their study uses the same activation scheme, Abdul-Razzak and Ghan (2000), and we cannot exclude the possibility that this scheme, for example, overestimates the supersaturation adjustment effect.

It is intrinsically difficult to directly evaluate cloud–aerosol interactions in models versus reality. This is partly because we cannot measure the pre-industrial atmosphere, but also due to the noisy nature of clouds. The evaluation of the model versions used in this study therefore focused on particles in sizes relevant for cloud activation and was primarily done in Blichner et al. (2021). We have added further validation using three datasets (Andreae et al., 2015; Wofsy et al., 2018; Andrade et al., 2015) representing different parts of the atmosphere than the previous comparison to the Supplement to this study (see Sect. S2 in the Supplement). Overall, the sectional scheme shows significant improvement in the representation of particles in the CCN size range, and this indicates that our results for ERF_{aci} here represent an improvement.

As mentioned in Sect. 2, the sectional scheme, OsloAeroSec, has a higher contribution from organics to growth from 5 nm than OsloAero_{def} and OsloAero_{imp} (only ELVOCs in OsloAero). One could argue that this may be the driving factor of all these results, but in fact this is not the case. We did a test run in which organics were treated in the same way in OsloAeroSec as in OsloAero_{def} and OsloAero_{imp}, and the result in terms of particle number changes very little (see Figs. S27 and S28 in the Supplement).

We also investigated the sensitivity of ERF_{aci} to changes in the nucleation rate with both the original model and with the sectional scheme. This investigation is detailed in Sect. S1

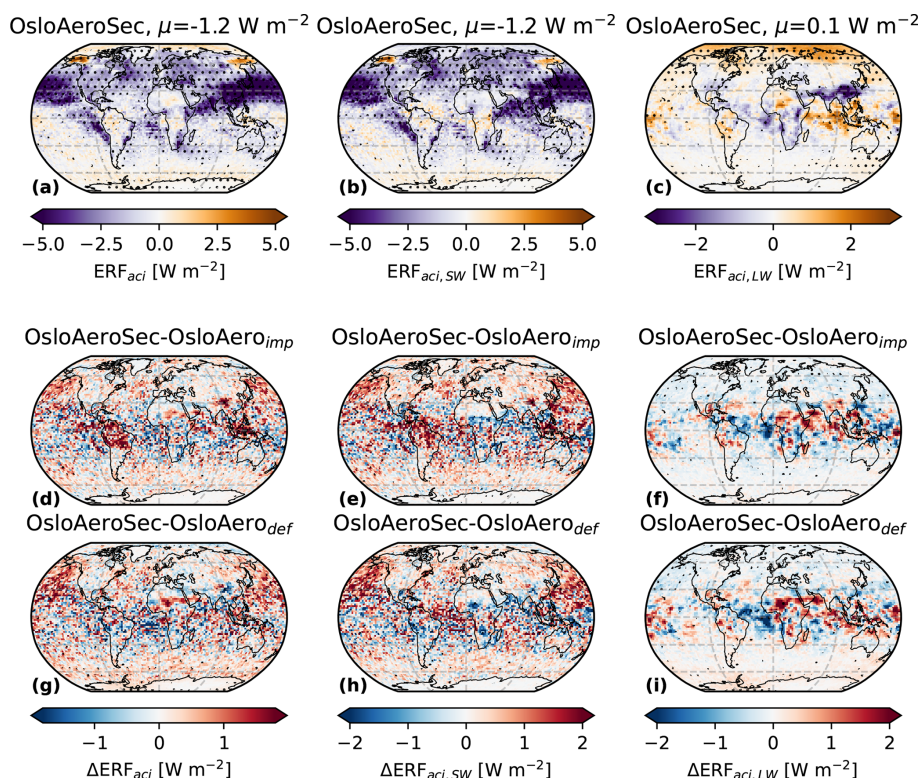


Figure 12. Annual averages of the ERF_{aci} (a, d and g), the shortwave component of ERF_{aci} , $ERF_{aci,SW}$ (b, e and h) and the longwave component of the ERF_{aci} , $ERF_{aci,LW}$ (c, f and i). The top panel shows the absolute values for OsloAeroSec, while the second and third rows show the difference of OsloAero_{imp} minus OsloAeroSec (second row) and OsloAero_{def} minus OsloAeroSec (third row). Dots are included in the plots to indicate where the difference between the two models is significant with a two-tailed paired Student's *t* test with 95 % confidence.

in the Supplement. Overall the results show that the change in ERF_{aci} between the sectional and default model is very resistant to changes in nucleation rate. There are small differences within the OsloAero model versions and within the OsloAeroSec versions based on the nucleation rate, but there are larger differences between the two groups.

Note that we have not discussed CCN concentrations in this discussion. There are two reasons for this: firstly, these are not yet available as standard output for CAM6-Nor. Secondly, the CCN concentrations at a given supersaturation matters only when this supersaturation is actually achieved, so focusing on CDNC gives a more complete picture which is more closely related to the actual climatic impact of the particles in question.

These results also illustrate the importance of adequately representing activation when investigating the effect of NPF on climate and not simply considering CCN at fixed supersaturation, as this will omit not only regional changes in updraft velocities, but also supersaturation adjustment by the aerosol population.

7 Conclusions

In this study, we have shown that including a sectional scheme (OsloAeroSec) for the growth of particles from nucleation up to the original modal scheme reduces the estimated ERF_{aci} by 0.13–0.14 $W m^{-2}$. The reduction originates from higher CDNC and $NCRE_{Ghan}$ in the PI simulation, together with a smaller increase from PI to PD. By comparing model versions with different NPF parameterizations in the pre-industrial and present-day atmosphere, we find that NPF in fact inhibits cloud droplet activation in parts of the atmosphere and leads to lower CDNC due to reducing the growth of the larger primary particles. The overall ERF_{aci} therefore depends on in which regions NPF is high or low in both the PI and the PD simulations. The reduction in ERF_{aci} with OsloAeroSec originates partly from higher NPF efficiency in PI areas where NPF enhances cloud droplet activation and lower NPF efficiency in PI areas where NPF inhibits cloud droplet activation. Furthermore, we find that the increase in sulfate from the PI to the PD simulation increases the hygroscopicity of the particles and thus allows more NPF particles to activate. This expands the areas where NPF enhances cloud droplet activation in the PD simulations, which

also contributes to a weaker ERF_{aci} for OsloAeroSec than OsloAeroDef.

Roughly speaking, we can say that the results in ERF_{aci} originate from OsloAeroSec: adding particles where the NPF particles are likely to act as CCN and removing them where they are unlikely to activate directly and rather act to diminish the size of the other particles.

Overall, this study shows that a more physical representation of the early growth of particles results in a lower ERF_{aci} and that adequately representing early growth on a regional scale is important for estimation of ERF_{aci} .

Code availability. The model code of NorESM2, release 2.0.1, is available at <https://doi.org/10.5281/zenodo.3760870> (Seland et al., 2020b; see Seland et al., 2020a, for details). The code modifications in OsloAeroSec are available at <https://doi.org/10.5281/zenodo.4265057> (Blichner, 2020); see Blichner et al. (2021) for details. The post-processing code for creating the figures in this paper is available at <https://doi.org/10.5281/zenodo.5559027> (Blichner, 2021b).

Data availability. The model output is available for download at <https://doi.org/10.11582/2021.00087> (Blichner, 2021a). Measurement data used in the Supplement are available for download as follows: the Chacaltaya data are available for download from the EBAS database (http://ebas-data.nilu.no/DataSets.aspx?nations=BO068BOL&components=particle_number_size_distribution&matrices=aerosol&fromDate=1970-01-01&toDate=2021-12-31, Andrade and Zaratti, 2021). The data from the ATTO were downloaded from the Laboratory of Atmospheric Physics (LFA, IF-USP) at http://ftp.lfa.if.usp.br/ftp/public/LFA_Processed_Data/T0a_ATTO/Level3/SMPS_2014toNov2020_ATTO_60m_InstTower/ (Laboratory of Atmospheric Physics LFA/IF-USP, 2021). The ATom data were accessed from <https://doi.org/10.3334/ORNLDAAC/1581> (Wofsy et al., 2018) and <https://doi.org/10.3334/ORNLDAAC/1613> (Wofsy and ATom Science Team, 2018).

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Author contributions. SMB did the model code development and performed the simulations with NorESM. SMB did the data analysis and wrote the paper. SMB, MKS and TKB contributed to discussions regarding the experimental design and data analysis. All authors have contributed to the discussions regarding the paper.

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