

Diurnal variation and size dependence of the hygroscopicity of organic aerosol at a forest site in Wakayama, Japan: their relationship to CCN concentrations

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Received: 29 August 2018 – Discussion started: 16 October 2018 Revised: 19 March 2019 – Accepted: 25 March 2019 – Published: 6 May 2019

Abstract. Formation of biogenic secondary organic aerosol (BSOA) and its subsequent evolution can modify the hygroscopicity of the organic aerosol component (OA) in the forest atmosphere, and affect the concentrations of cloud condensation nuclei (CCN) there. In this study, size-resolved aerosol hygroscopic growth at 85 % relative humidity and size-resolved aerosol composition were measured using a hygroscopic tandem differential mobility analyzer and an aerosol mass spectrometer, respectively, at a forest site in Wakayama, Japan, in August and September 2015. The hygroscopicity parameter of OA (κ_{org}) presented daily minima in the afternoon hours, and it also showed an increase with the increase in particle dry diameter. The magnitudes of the diurnal variations in κ_{org} for particles with dry diameters of 100 and 300 nm were on average 0.091 and 0.096, respectively, and the difference in κ_{org} between particles with dry diameters of 100 and 300 nm was on average 0.056. The relative contributions of the estimated fresh BSOA and regional OA to total OA could explain 40% of the observed diurnal variations and size dependence of $\kappa_{\rm org}$. The hygroscopicity parameter of fresh BSOA was estimated to range from 0.089 to 0.12 for particles with dry diameters from 100 to 300 nm. Compared with the use of time- and size-resolved κ_{org} , the use of time- and size-averaged $\kappa_{\rm org}$ leads to under- and overestimation of the fractional contribution of OA to CCN number concentrations in the range from -5.0% to 26%. This indicates that the diurnal variations and size dependence of $\kappa_{\rm org}$ strongly affect the overall contribution of OA to CCN concentrations. The fractional contribution of fresh BSOA to CCN number concentrations could reach 0.28 during the period of intensive BSOA formation. The aging of the fresh BSOA, if it occurs, increases the estimated contribution of BSOA to CCN number concentrations by 52 %–84 %.

1 Introduction

The hygroscopicity (ability to absorb water) of organic aerosol (OA) components is governed by their chemical composition, and has two important roles in the atmosphere. It influences light scattering by aerosols (Titos et al., 2016) and affects the ability of aerosols to work as cloud condensation nuclei (CCN; McFiggans et al., 2006). The hygroscopicity of OA may also influence the aqueous chemistry in aerosols and cloud droplets, which provide a potentially important pathway for the formation of secondary organic aerosols (SOA; McNeill, 2015). Nevertheless, the hygroscopicity of OA is not well characterized in terms of its temporal and spatial variations, size dependence, and its relationship with the chemical composition of OA, given that OA is a complex mixture of a number of compounds.

Studies on the hygroscopicity parameter κ of OA (κ_{org} ; Petters and Kreidenweis, 2007) in different locations have

presented different characteristics of temporal variation and size dependence. Based on a year-long observation under supersaturated water vapor conditions (SUPS) at a downwind site of Manaus in central Amazonia, Thalman et al. (2017) reported that $\kappa_{\rm org}$ presented the lowest value of ~ 0.1 in September and the highest value of ~ 0.15 in December, and that the ranges of the diurnal variations in $\kappa_{\rm org}$ were 0.10 to 0.16 (night to day) and 0.08 to 0.14 (night to day) under the influence of local biomass-burning air masses during the dry season and urban-pollution air masses during the wet season, respectively. Bougiatioti et al. (2016) reported diurnal variation in κ_{org} in the range 0.09–0.18 (day to night) under SUPS for particles influenced by biomass burning in the eastern Mediterranean. Deng et al. (2018) reported diurnal variation ranges of κ_{org} of 0.09 to 0.30 (day to night) and 0.16 to 0.24 (day to night) on days with and without evident new particle formation (NPF), respectively, under a SUPS (0.23 % water vapor supersaturation (SS) condition) in a forest in Wakayama, Japan. Different from the above studies, at a rural site in the southeastern United States, a small diurnal variation in $\kappa_{\rm org}$ (~ 0.13 to ~ 0.17, night to day) under SUPS was observed (Cerully et al., 2015). With respect to the size dependence of $\kappa_{\rm org}$, airborne studies over United States, Canada, Pacific Ocean, and the Gulf of Mexico for a variety of air mass types presented a decrease in κ_{org} (from 0.13 to 0.06) with an increase in the particle modal diameter (from 130 to \sim 210 nm) under subsaturated water vapor conditions (SUBS) (Shingler et al., 2016). By contrast, ground-based observations in the city of Nagoya in Japan under SUBS presented relatively low κ_{org} in small particles (0.12–0.15; 60 and 100 nm) and high κ_{org} in large particles (0.17-0.22; 200 and 359 nm) (Kawana et al., 2016). For aerosols found in forest areas, Deng et al. (2018) reported a difference of 0.03 in mean $\kappa_{\rm org}$ between sub-100 nm particles ($\kappa_{\rm org}$ was 0.19) and ~ 150 nm particles ($\kappa_{\rm org}$ was 0.22), whereas Thalman et al. (2017) did not identify any size dependence of $\kappa_{\rm org}$ for 94–171 nm particles. The hygroscopicity of laboratory-generated model SOA was also reported to be size-dependent. Frosch et al. (2011) reported that the κ of α -pinene SOA at SUPS at 100 nm (~ 0.12) was ~ 0.06 higher than at 200 nm. Zhao et al. (2015) reported that the κ values of model SOA at SUPS produced by different precursors at 50, 100, and 200 nm were ~ 0.17 , ~ 0.11 , and ~ 0.07 , respectively. Tritscher et al. (2011) also found that small (50 nm) α -pinene SOA particles had a higher κ than large ones (150 nm) at SUBS, although the difference was small (0.03). Frosch et al. (2013) reported that the κ of β caryophyllene SOA decreased with the increase in SS, which can be interpreted as the increase in κ with the increase in particle diameter, and that the maximum of the difference was about 0.1.

The variations in κ_{org} observed in the aforementioned studies may have great influence on the prediction of CCN. Based on global climate modeling simulations, Liu and Wang (2010) reported that CCN concentration would change within 40 % by changing the κ of SOA by ± 50 % (from 0.14) to 0.07 or 0.21). Rastak et al. (2017) reported that the difference in the aerosol radiative effects between κ_{org} of 0.05 and 0.15 was -1.02 W m^{-2} , the order of which is the same as that of the overall climate forcing effect of anthropogenic aerosol during the industrial period. Based on CCN closure studies, Wang et al. (2008) reported that, for above-cloud aerosols with high volume fractions of OA, while the CCN number concentration closure could be achieved using κ_{org} of 0.12, the use of κ_{org} of 0.25 led to overestimation of the CCN number concentration by 50 %. Mei et al. (2013b) reported that the increase in κ_{org} from 0.08 to 0.13 led to a 30 % increase in the calculated CCN number concentration and that the increase from 0.03 to 0.18 doubled the concentration. It is therefore important to study the temporal variation and size dependence of κ_{org} in more locations where OA dominates the aerosol chemical composition, to characterize the $\kappa_{\rm org}$ values and to represent $\kappa_{\rm org}$ appropriately in model predictions of CCN number concentrations.

The temporal variation and size dependence of κ_{org} of ambient aerosol is reported to relate to variations in the chemical composition of OA, which can result from the mixing of aerosols of different origins, formation of SOA, and aging processes (Cerully et al., 2015; Bougiatioti et al., 2016; Shingler et al., 2016; Thalman et al., 2017; Deng et al., 2018). The size-dependent chemical composition of model SOA has been explained by the size-dependent contributions of different organic vapors to particle growth (e.g., Winkler et al., 2012; Ehn et al., 2014; Zhao et al., 2015, 2016). The sizedependent κ of model SOA has also been explained from the viewpoint of size-dependent chemical composition (Zhao et al., 2015; Frosch et al., 2013) and other factors: the dependence of water activity on particle size, the dependence of the surface tension on the solution concentration, and the evaporation of semivolatile SOA under high-SS conditions (Frosch et al., 2011, 2013; Zhao et al., 2015). In recent studies, the variation in κ_{org} was explained by the variation in OA subcomponents derived from positive matrix factorization (PMF) analysis of OA mass spectra (Cerully et al., 2015; Bougiatioti et al., 2016). From these studies it is reported that the daily variation in $\kappa_{\rm org}$ could be well explained by the daily variation in the contributions of the retrieved PMF factors to $\kappa_{\rm org}$.

In the forest atmosphere, the oxidation of biogenic volatile organic compounds (BVOC) emitted by vegetation can produce substantial amounts of biogenic secondary organic aerosols (BSOA; Tunved et al., 2006; Pöschl et al., 2010; Han et al., 2014). BSOA is reported to contribute to the growth of newly formed particles in forests (e.g., Han et al., 2013; Yu et al., 2014; Zhou et al., 2015). BSOA may also condense on preexisting background particles or particles transported with inflowing air masses (e.g., Cerully et al., 2015; Thalman et al., 2017). Moreover, BSOA is subject to aging processes that include photochemical oxidation and aqueous-phase reactions that must depend on ambient meteorological condi-

tions (e.g., Han et al., 2014; Thalman et al., 2017). Such processes could result in time- and size-dependent variation in the chemical composition of OA and thus time- and size-dependent κ_{org} in the forest atmosphere. However, the characteristics of the temporal variations and size dependence of κ_{org} , and their relationships to the atmospheric processes of BSOA in forest environments, are not well understood.

We performed field observation at a forest site in Wakayama, Japan, in August and September 2015 and characterized the diurnal variations and size dependence of $\kappa_{\rm org}$. The variations and dependence were interpreted based upon the size-resolved chemical composition of OA from the viewpoint of BSOA formation. Furthermore, the influence of these variations on the fractional contribution of OA and BSOA to the CCN concentration was assessed. Previous observational studies at the site indicated that BSOA formation was intensive and that aging occurred after formation (Han et al., 2014; Deng et al., 2018). It was also observed that $\kappa_{\rm org}$ was time- and size-dependent and that the contribution of OA and BSOA to CCN number concentrations could be substantial (Deng et al., 2018). This work is an extension of previous studies on the hygroscopicity and CCN activity of aerosols, and the contributions of OA and BSOA to the CCN concentration, in the same forest (Kawana et al., 2017; Deng et al., 2018). It is intended to clarify the diurnal variation and size dependence of the hygroscopicity of OA and their influence on the contributions of OA and BSOA to CCN.

2 Field observation

The field observation was performed at Wakayama Forest Research Station, Kyoto University (34.06° N, 135.52° E, about 500 m above sea level), located in the central part of the Kii Peninsula. The observation site is about 70 km south of Osaka (2.7 million inhabitants) and 60 km northwest of the North Pacific. Both coniferous trees (such as *Cryptomeria japonica*, *Chamaecyparis obtuse*, and *Pinus densiflora*) and broad-leaf trees (such as *Quercus serrata* and *Quercus crispula*) are distributed on the Kii Peninsula (Okumura, 2009). The study period was from 14:30 Japan Standard Time (JST) on 31 August to 06:00 JST on 22 September 2015.

The hygroscopic growth at 85 % relative humidity (RH), number–size distributions, and size-resolved chemical composition of ambient aerosols were measured using a hygroscopicity tandem differential mobility analyzer (HTDMA), a scanning mobility particle sizer (SMPS), and a highresolution time-of-flight aerosol mass spectrometer (AMS), respectively. Ambient air was aspirated from an inlet about 7.5 m above the ground. The air was transferred through a $PM_{2.5}$ cyclone (URG) installed at the lower end of the 10.4 m stainless-steel inlet tubing (1/2 in. OD) and introduced to the instrument room at a flow rate of 16.7 L min⁻¹. A manifold combined with an assistant pump (ULVAC, DA30S) was used to split the air flow, and the sample flow for the instrument system composed of the HTDMA, SMPS, and AMS was $0.9 \,\mathrm{L\,min^{-1}}$. The sample flow upstream of the AMS was dried with two diffusion driers containing silica gel. The sample flow to the HTDMA and SMPS was dried with three diffusion driers, in series, two with silica gel and one with molecular sieves.

In the HTDMA, the dried aerosol (RH < 1.8 %) was passed through the first differential mobility analyzer (DMA1; 3081, TSI), where the aerosol was classified, and quasimonodisperse particles of 30, 50, 70, 100, 200, 300, and 360 nm in diameter (d_{drv}) were obtained. The setting for the classification was fixed for 5 min at each diameter. In each hour, the sequential diameter setting of DMA1 was 30, 50, 70, 100, 200, 360, 30, 50, 100, 200, 300, and 360 nm. During 05:50-05:54 and 17:50-17:54 JST, the setting of DMA1 was for system performance check. The classified aerosol was passed through a Nafion humidifier (MD-110-24S-4, Perma Pure) where it was humidified to $\sim 85\%$ RH. The aerosol was then introduced to a second DMA (DMA2; 3081, TSI) coupled to a condensation particle counter (CPC, 3775, TSI), which was operated by scanning the voltage of DMA2. For both DMA1 and DMA2, the aerosol flow rate was $0.3 \,\mathrm{L\,min^{-1}}$ and the sheath-to-sample flow ratio was 10:1. The residence time of the monodisperse particles from the outlet of the humidifier to the inlet of DMA2, where the RH was considered to be $\sim 85\%$, was approximately 11 s. The sheath air flow of DMA2 was also humidified using another Nafion humidifier (PD-100T-12MSS, Perma Pure). The RH (temperature) measured (HMP237, Vaisala) at the inlets of aerosol flow and sheath flow to DMA2 were $85.0\pm0.2\%$ (20.3±0.5° C) and $85.0\pm0.2\%$ (20.4±0.5° C), respectively, and that at the outlet of the sheath flow of DMA2 was $86.0 \pm 0.3 \%$ (20.3 $\pm 0.6^{\circ}$ C). For analysis of the particle hygroscopic growth, RH of 85 % was applied. The SMPS for the measurement of aerosol number-size distributions was composed of a third DMA (DMA3; 3080, 3081, TSI) and a CPC 3772 (TSI). The aerosol flow (RH < 1.6%)was $0.3 \,\mathrm{L\,min^{-1}}$ and the sheath to aerosol flow ratio was 10:1. At the inlet of CPC 3772, the sample flow was diluted to 1 L min⁻¹ with purified dry air, which was generated using a compressor (RD-45-N, IAC) and an air dryer (QD 30-50, IAC). The aerosol number-size distributions were measured for a dry diameter range of 13.8–749.9 nm every 5 min. The performances of the three DMAs were assessed using standard size PSL particles before and after the observation (Sect. S1 in the Supplement). Furthermore, an aqueous solution of ammonium sulfate (AS) (99.999 % purity, Sigma-Aldrich) was nebulized and the generated aerosols were dried and introduced to the HTDMA, to assess the consistency of the sizing of the two DMAs under dry condition, and to validate the RH setting before the observation (Sect. S2). The setup and calibration procedures of the AMS were the same as those for the observations in 2014 (Deng et al., 2018). The V-mode (MS and PToF modes) data were analyzed using

the Igor high-resolution data analysis package (PIKA1.20Q, Igor 6.37) to obtain the bulk and size-resolved mass concentrations of the chemical components (sulfate (SO₄), ammonium (NH₄), nitrate (NO₃), chloride (Chl), and OA), and the atomic ratios of O to C (O : C ratio) and H to C (H : C ratio) for organics. In addition, high-resolution bulk OA mass spectra observed in V-mode were subjected to PMF analysis (Paatero and Tapper, 1994; Ulbrich et al., 2009) (Sects. 3.3, S3). The RH of the sample flow was lower than 0.5 %.

A single-wavelength particle soot absorption photometer (1 λ -PSAP, 567 nm, Radiance Research Inc.), with a thermodenuder maintained at 300 °C, was deployed to obtain the mass concentration of sub-micrometer black carbon (BC; Kondo et al., 2009; Deng et al., 2018). The mixing ratios of target gaseous species, NO–NO₂–NO_x, CO, CO₂, and O₃, were monitored using commercial instruments (APNA-370, Horiba, for NO–NO₂–NO_x; model 48ij, Thermo Fisher Scientific, for CO; LI-820, LI-COR, for CO₂; model 49ij, Thermo Fisher Scientific, for O₃). Meteorological data were collected (Kyoto University, 2017). Air temperature, RH (HMP-155, Vaisala), precipitation (RH-5E, IKEDA-KEIKI), and solar radiation (CMP3-L, Campbell) were used in this study.

All the observation data except meteorological data were screened to eliminate data that might have been under the strong influence of local anthropogenic emissions, for example those from vehicles. This was performed by omitting data with spikes in the number concentrations of aerosols from their size distribution data, and in the mass concentration of BC (Sect. S4).

3 Data analysis

3.1 Hygroscopicity of ambient aerosols

The hygroscopic growth factor of aerosol particles, g_f , was defined as the ratio of the particle wet diameter (d_{wet} , 85 % RH) to the corresponding dry diameter (d_{dry}). The distributions of g_f for specific d_{dry} ($n(g_f)$, i.e., the number distribution of particles as a function of g_f) were retrieved using the Twomey algorithm as presented by Mochida et al. (2010) with consideration of the shape of the transfer functions of the two DMAs. The difference in the processing is that the transfer function and the $n(g_f)$ in this study were analyzed and presented in the fine mode of 1024 diameter bins per decade while 64 bins per decade were used in Mochida et al. (2010). The g_f probability distribution function, g_f -PDF, is the normalized $n(g_f)$. The g_f -PDF in this study is presented in linear scale, which was converted from the original logarithmic scale distribution. The time-resolved mean value of

 $g_{\rm f}$ for respective $d_{\rm dry}$, $g_{\rm f,m}$, was calculated as follows.

$$g_{f,m} = \frac{\sum n (g_f) g_f}{\sum n (g_f)} \quad (0.8 \le g_f \le 2.2 \text{ for } 30 \le d_{dry})$$

$$\le 300 \,\text{nm}, \text{ or } 0.8 \le g_f \le 2.0 \text{ for } d_{dry} = 360 \,\text{nm}) \quad (1)$$

For ambient particles, the time-resolved mean-water-volume equivalent $g_f(g_{f,mw}, i.e.)$, the average of g_f that corresponds to the mean water volume retained by particles of certain d_{dry}) was also calculated using Eq. (2) (Kawana et al., 2016).

$$g_{\rm f,mw} = \left[\frac{\sum n \left(g_{\rm f}\right) \left(g_{\rm f}^3 - 1\right)}{\sum n \left(g_{\rm f}\right)} + 1\right]^{\frac{1}{3}}$$

(0.8 \le g_{\rm f} \le 2.2 for 30 \le d_{\rm dry} \le 300 nm,
or 0.8 \le g_{\rm f} \le 2.0 for d_{\rm dry} = 360 nm) (2)

The hygroscopicity parameter of ambient particles at 85 % RH (κ_t) was calculated following the κ -Köhler theory (Petters and Kreidenweis, 2007).

$$\kappa_t = \left(g_{f,\text{mw}}^3 - 1\right) \left[\frac{\exp\left(\frac{4\sigma M_w}{RT\rho_w d_{\text{wet}}}\right)}{0.85} - 1\right],\tag{3}$$

where σ is the surface tension at the solution-air interface; $M_{\rm w}$ and $\rho_{\rm w}$ are the molecular mass and density of pure water, respectively; $d_{\rm wet}$ is the product of $g_{\rm f,mw}$ and $d_{\rm dry}$; R is the universal gas constant; and T is the temperature in kelvin. In this study, the mean temperature at the inlets of aerosol flow and sheath flow of DMA2, weighted by their flow rates, was applied as T (294 K) and the surface tension of pure water at this temperature (Vargaftik et al., 1983) was used as σ in Eq. (3). Because κ_t was calculated from $g_{\rm f,mw}$, the aerosol mixing state was not considered in the analysis of κ in this study.

3.2 Hygroscopicity of OA

The hygroscopicity parameter of organics, κ_{org} , was calculated using Eq. (4) assuming the volume additivity of water retained by different aerosol components (Petters and Kreidenweis, 2007).

 $\kappa_t = \varepsilon_{\rm org} \kappa_{\rm org} + \varepsilon_{\rm inorgsalt} \kappa_{\rm inorgsalt} + \varepsilon_{\rm BC} \kappa_{\rm BC}$

$$=\varepsilon_{\rm org}\kappa_{\rm org} + \sum_{i=1}^{3}\varepsilon_{\rm i}\kappa_{\rm i} + \varepsilon_{\rm BC}\kappa_{\rm BC} \tag{4}$$

Here, κ_t is the hygroscopicity parameter of ambient aerosol at 85 % RH calculated using Eq. (3), while κ_{org} , $\kappa_{\text{inorgsalt}}$, and κ_{BC} are the hygroscopicity parameters of OA, inorganic salts, and BC, respectively. The volume fractions of OA, inorganic salts, and BC are ε_{org} , $\varepsilon_{\text{inorgsalt}}$, and ε_{BC} , respectively, and ε_i and κ_i are the volume fraction and hygroscopicity parameter of the inorganic salts: ammonium nitrate (AN), sulfuric acid (SA), ammonium hydrogen sulfate (AHS), letovicite (LET), and ammonium sulfate (AS). The ε_{org} , ε_i , and ε_{BC} were calculated based on the size-resolved mass concentrations of organics, sulfate, nitrate, and ammonium from the AMS, and the sub-micrometer BC mass concentrations from the PSAP. BC was assumed to have the same mass-size distribution as OA. The aerosol particles were assumed to be spherical and without voids. PToF mode data in vacuum aerodynamic diameter (d_{va}) ranges that were ~ 1.0 (0.98–0.99) to 2.0 times that of $d_{\rm dry}$, corresponding to the particle density of ~ 1.0 (0.98-0.99) to $2.0 \,\mathrm{g}\,\mathrm{cm}^{-3}$, were adopted. More details about the calculations of the size-resolved $\varepsilon_{\rm org}$, ε_i , and $\varepsilon_{\rm BC}$ are presented in Sect. S5. The derivation of κ_i was based on the online Extended AIM Aerosol Thermodynamics Model II (E-AIM II, http://www.aim.env.uea.ac.uk/aim/kohler/input kohler.html, last access: 1 August 2018; Clegg et al., 1998; Wexler and Clegg, 2002) as presented in Sect. S6 and Table S3 in the Supplement. The κ of BC was assumed to be zero. Because of the low signal intensity of the PToF data in the sub-100 nm d_{va} range (Sect. S7), the κ_{org} was only derived for particles with $d_{\rm dry}$ of 100, 200, 300, and 360 nm. Furthermore, to assess the influence of the choice of the d_{va} range on the derivation of $\kappa_{\rm org}$, the derived $\kappa_{\rm org}$ for particles with $d_{\rm dry}$ of 100 nm using the chemical composition in the d_{va} range 98–197 nm was compared with that using the chemical composition in the d_{va} range 69–138 nm (Fig. S4 in the Supplement). The result indicates that κ_{org} was not sensitive to change in the selected d_{va} range when ε_{org} was greater than 40 %. Note that, although the volume additivity assumption between organics and inorganics may not necessarily hold (Vaishya et al., 2013; the κ_{org} derived in the manner in this study represents the perturbation of κ_t as a result of the presence of organics), the inverse linear correlation between κ_t and $\varepsilon_{\rm org}$ (correlation coefficients: -0.45 to -0.83; Fig. S5) suggests that the additivity holds well for the aerosols studied.

3.3 PMF analysis of OA mass spectra

To characterize the diurnal variations and size dependence of $\kappa_{\rm org}$, the high-resolution OA bulk mass spectra derived from the V-mode AMS data were subjected to PMF analysis (Sect. S3), followed by derivation of the size-resolved contributions of the PMF factors to the OA mass concentration (Sect. S8). A two-factor PMF solution was adopted, which resolved two oxygenated OA factors: one with a lower atomic O:C ratio (0.47) named less-oxygenated organic aerosol (LOOA), and the other with a higher O:C ratio (0.95) named more-oxygenated organic aerosol (MOOA). The low relative residual (2.6%) for the bulk mass spectra supports the use of the two PMF factors to illustrate the observed OA. Note that the two OA factors resolved here represent two different groups of OA chemical structures, not necessarily two different OA sources (Zhang et al., 2011). The use of a PMF result with more factors could make illustration of the variation in κ_{org} complex and was not adopted. The PToF mode OA mass spectra in 2 h time resolution were attributed to the two PMF factors through multivariable linear regression (Sect. S8). For particles with d_{dry} equal to or larger than 100 nm, the variation in κ_{org} was discussed with regard to the variations in the two PMF factors. Furthermore, the hygroscopicity parameters for the two OA fractions were derived and then used to estimate the hygroscopicity of freshly formed BSOA (Sect. 4.2).

4 Results and discussions

4.1 Overview of the observations

4.1.1 Meteorological conditions, gaseous species, and aerosol chemical composition

During the measurement period with effective data, the mean \pm standard deviation (SD) of the temperature and RH of the ambient air were $18.2 \pm 2.4^{\circ}$ C and 94.2 ± 7.6 %, respectively. Precipitation events occurred intermittently during 1-3, 6-10, and 16-17 September (Fig. S7). Backward air mass trajectories (Fig. S9) generated using NOAA's HYS-PLIT atmospheric transport and dispersion modeling system (Draxler and Hess, 1998) indicate that, apart from on 1 and 17 September, most of the air masses that arrived at the observation site had traveled from the Japan archipelago (and even from the Asian continent) within 5 days, and may have transported aged anthropogenic pollutants to the observation site. The mean \pm SD of the BC concentration during the entire study period was $0.07 \pm 0.06 \,\mu g \,m^{-3}$ (Fig. S7). The mean \pm SD of the mixing ratios of CO, NO, NO₂, NO_x , and O_3 during the entire study period were 164 ± 42 , 0.33 ± 0.12 , 0.56 ± 0.35 , 0.63 ± 0.38 , and 11.5 ± 8.4 ppb, respectively (Fig. S7). The concentration of BC was low, and the mixing ratios of CO and NO_x were modest. A daily maximum of BC appeared in the afternoon hours (Fig. 2), which, however, might have been caused by the charring of OA at the heating temperature of 300° C. The mixing ratios of CO and NO_x tended to be relatively high during 10:00-22:00 JST (Fig. S8), which might have been caused by the transport of anthropogenic pollution to the surface site by enhanced vertical convection in the daytime. The concentration of O₃ was substantial and presented obvious diurnal variation (Fig. S7). On average, O_3 peaked during noon with the solar radiation (Fig. S8), indicating the occurrence of photochemical reactions during the daytime.

The time series of the mass concentrations of aerosol chemical components and aerosol number–size distributions are presented in Fig. 1. Among non-refractory aerosol chemical components derived from the AMS and BC derived from the PSAP (total concentration: $6.2 \pm 4.4 \,\mu g \, m^{-3}$), organics on average accounted for the largest fraction (45.0%, of which LOOA and MOOA accounted for 24.9% and 19.1%,

respectively), followed by sulfate (39.6%) and ammonium (11.6%). The contributions of nitrate, BC, and chloride were minor: their mass fractions were on average 2.5%, 1.1%, and 0.2%, respectively. The contribution of OA to the submicrometer aerosol mass increased and that of sulfate decreased in the afternoon hours (12:00–20:00 JST). The mean aerosol number concentration ($N_{\rm CN}$) was 1241±1012 cm⁻³. The geometric mean diameter of the aerosols ranged from 45 to 154 nm with a mean ± SD of 88±17 nm. No strong burst of small particles (i.e., $d_{\rm dry} < 30 \,\rm nm$) was identified during the observation, which is different from the two former observations in 2010 and 2014 (Han et al., 2013; Deng et al., 2018).

The diurnal variation in the number-size distributions and the mass concentrations of the chemical components of aerosols are presented in Fig. 2. The $N_{\rm CN}$, OA, and LOOA presented similar diurnal variation patterns. Their daily minima were observed between 06:00 and 08:30 JST. After 08:30 JST, they increased monotonically and reached their maxima during 15:00-18:00 JST. Then they gradually decreased until approximately 06:00 JST of the next day. MOOA also increased slowly (following the trend of LOOA) in the daytime and reached its maximum around 18:00 JST. The pattern of the enhancement of OA in the daytime followed that of the solar radiation (Fig. 2b), indicating that the enhancement of OA was caused by the formation of BSOA through photochemical reactions of BVOC (Han et al., 2014; Deng et al., 2018). This is supported by an analysis indicating that anthropogenic pollution was not the main contributor to the enhancement of OA, at least during the period 12:00–16:00 JST (Sect. S9), and by the report that primary biogenic OA is mainly in the supermicrometer aerosol diameter range in a forest environment, for the Amazon at least (Pöschl et al., 2010). Furthermore, the stronger enhancement of LOOA than of MOOA indicates that the freshly formed BSOA was mainly composed of LOOA and had a low oxygenation state. The O: C ratio of OA increased slowly from around noon to midnight (Fig. 2c), together with the appearance of MOOA, indicating the aging of freshly formed BSOA (Han et al., 2014). Because of high-RH conditions (Fig. S7), aqueous-phase reactions including in-cloud processes could have played an important role in the aging of fresh BSOA (Han et al., 2014), which could have modified the hygroscopicity of ambient aerosols (Jimenez et al., 2009; Farmer et al., 2015). Although no abrupt increase in sub-30 nm particles was observed, the increase in the number concentration of 30-50 nm particles around noon indicates the formation of new particles near the observation site. These particles had probably grown by the condensation of BSOA formed from BVOC at the time they were transported to the observation site. The concentrations of nitrate and chloride stayed low, although they also presented maxima in the afternoon. Sulfate, which may have been strongly influenced by transported anthropogenic aerosol, did not present an obvious diurnal variation. This result supports the view

that the contribution of anthropogenic OA to the observed enhancement of OA was small.

4.1.2 Hygroscopicity of atmospheric aerosols

Similar to a prior observation at the same site in 2010 (Kawana et al., 2017), the hygroscopic growth factor g_f presented unimodal distributions at respective particle diameters (Figs. 1d, e, and S11), and the mean hygroscopic growth factor $g_{f,m}$ of the aerosols increased with an increase in the particle diameters (Fig. S12). The mean \pm SD of $g_{f,m}$ at 30, 50, 70, 100, 200, 300, and 360 nm were 1.13 ± 0.08 , $1.21 \pm 0.09, 1.22 \pm 0.09, 1.26 \pm 0.10, 1.36 \pm 0.10, 1.40 \pm 0.08,$ and 1.42 ± 0.08 , respectively. The unimodal pattern of the gf-PDF indicated the internal mixing state of the observed aerosol at the respective particle diameters. Decreases in $g_{f,m}$ (Fig. S12) were observed for all particles during periods of intensive BSOA formation (i.e., episodes when the mass concentration of OA especially LOOA greatly increased; such episodes were observed on 31 August, and on 2, 5, 7, 9, 14, 15, 17, 18, 19, 20, and 21 September; Fig. 1b).

The hygroscopicity parameter of ambient aerosol particles that corresponds to $g_{f,mw}(\kappa_t)$ also increased with the increase in aerosol particle diameters (Fig. 3a). Similar diurnal variation patterns were observed for all the diameters studied. The κ_t started to decrease around 08:00 JST, then reached daily minima between 13:00 and 19:00 JST. Then it increased continually until around 02:00 JST of the next day, and remained high until 08:00 JST the next morning. For particles with $d_{dry} \ge 100$ nm, the diurnal variation pattern and size dependence of κ_t were opposite to those of the volume fraction of OA (Fig. 3b) and were similar to those of the volume fraction of total inorganic salts (Fig. S13). The results suggest that, at least for ambient aerosol particles with $d_{\rm dry} \ge 100$ nm, OA and inorganic salts had low and high hygroscopicity, respectively, and this resulted in the variations in κ_t . Although $\kappa_{\text{inorgsalt}}$ is much greater than κ_{org} (Petters and Kreidenweis, 2007), the high ε_{org} makes the influence of OA on κ_t significant. Thus, the variation in κ_{org} (Sect. 4.2) may also contribute to the variation in κ_t . For particles with $d_{\rm dry} \leq 70$ nm, the decrease in the particle hygroscopicity with the decrease in the particle diameter is also explained by the accompanying increase in ε_{org} and the decrease in $\varepsilon_{\text{inorgsalt}}$ (Levin et al., 2014). This is indicated by the substantially lower mean mass fraction of inorganic salts than of organics in the corresponding d_{va} range of less than 150 nm (Fig. S2). In particular, for particles with d_{dry} of 30 nm, κ_t remained constant in a low range (0.079–0.089) from 13:00 to 17:00 JST with a mean value of 0.082. BSOA formed during this period probably dominated the particle mass (Kawana et al., 2017; Han et al., 2014).



Figure 1. Time series of (a) sub-micrometer mass concentrations of non-refractory aerosol chemical components (OA, SO₄, NO₃, NH₄, and Chl) from the AMS measurement and BC from the PSAP measurement; (b) mass concentrations of LOOA and MOOA, and the residuals from the PMF analysis; (c) aerosol number–size distributions; and $g_{\rm f}$ -PDF of aerosol particles with $d_{\rm dry}$ of (d) 50 and (e) 200 nm. The two pie charts in panel (a) present the mass fractions of chemical components for the entire study period and for the afternoon hours (12:00–20:00 JST) during the study period. The dashed line in panel (c) represents a diameter of 30 nm.

4.2 Hygroscopicity of organic aerosol components

4.2.1 Variation in κ_{org} and its relation to the chemical structure of OA

The diurnal variation in κ_{org} and the volume fraction of LOOA in OA, $v_{\text{LOOA}}/(v_{\text{LOOA}} + v_{\text{MOOA}})$, where v_{LOOA} and v_{MOOA} refer to the volume concentrations of LOOA and MOOA, respectively, for the entire study period are presented in Fig. 4. Data with $\varepsilon_{\rm org} < 0.40$ were excluded from the $\kappa_{\rm org}$ values presented because the uncertainty that originated from subtraction of the contribution of inorganic components was considered large in the low ε_{org} range (Mei et al., 2013a). The $\kappa_{\rm org}$ decreased rapidly from approximately 08:00 JST in the morning when the mass concentrations of OA and LOOA started to increase (Fig. 2). The κ_{org} reached daily minima during 10:00-18:00 JST and increased after the minima (Fig. 4a). The $v_{\text{LOOA}}/(v_{\text{LOOA}} + v_{\text{MOOA}})$ in Fig. 4b presents the opposite diurnal variation pattern. The characteristics of the size dependence of κ_{org} and $v_{\text{LOOA}}/(v_{\text{LOOA}}+v_{\text{MOOA}})$ were dependent on time periods. To characterize the size dependence of $\kappa_{\rm org}$ and $v_{\rm LOOA}/(v_{\rm LOOA}+v_{\rm MOOA})$, the mean values of κ_{org} and $v_{\text{LOOA}}/(v_{\text{LOOA}}+v_{\text{MOOA}})$ during 12:00–20:00 and 20:00–12:00 JST were plotted separately in the $\kappa_{\rm org}$ – $v_{\text{LOOA}}/(v_{\text{LOOA}}+v_{\text{MOOA}})$ space for different d_{dry} (Fig. 5), and the difference in the diurnal variation data between particles with different diameters were evaluated using a 10 % two-sided t test (Table S7). During 12:00–20:00 JST, opposite size dependences were observed between the mean κ_{org} and $v_{\text{LOOA}}/(v_{\text{LOOA}} + v_{\text{MOOA}})$. Although the differences in $\kappa_{\rm org}$ between 200 and 300 nm particles (p value: 0.71) and 300 and 360 nm particles (p value: 0.15) were not significant, the differences in κ_{org} between 100 and 200 nm particles (p value: 0.01) and 200 and 360 nm particles (p value: 0.07), and the differences in $v_{\text{LOOA}}/(v_{\text{LOOA}}+v_{\text{MOOA}})$ between particles with all different diameters (p value: <0.02), were significant during that period. During 20:00-12:00 JST, the size dependences of $\kappa_{\rm org}$ and $v_{\rm LOOA}/(v_{\rm LOOA} + v_{\rm MOOA})$ were not clear. The clearer size dependence of both κ_{org} and $v_{\text{LOOA}}/(v_{\text{LOOA}} + v_{\text{MOOA}})$ during 12:00–20:00 JST than during 20:00-12:00 JST was explained by the formation of BSOA during the afternoon hours. The patterns of diurnal variation and of size dependence between κ_{org} and $v_{\text{LOOA}}/(v_{\text{LOOA}} + v_{\text{MOOA}})$ during 12:00–20:00 JST indicate that the variation in κ_{org} could be explained at least in part by the relative contributions of LOOA and MOOA to OA. That 5896



Figure 2. Diurnal variations in (a) number-size distribution (image plot) and number concentration (NCN, right axis) of ambient aerosols; (b) the mass concentrations of OA, SO₄, and NH₄ (left axis), and NO₃, BC, and Chl (right axis in black); and (c) the mass concentrations of LOOA, MOOA, and the residual, and the O : C ratio of bulk OA (only data with $m_{\rm org} > 0.3 \,{\rm mg m}^{-3}$ are included). (A box-and-whisker plot of the diurnal variation in O : C ratio is presented in Fig. S19.) The diurnal variation in solar radiation is superimposed in panel (b). All data presented are averages for the entire study period.

is, the presence of LOOA with low oxygenation state (O : C ratio of 0.47) lowered the observed κ_{org} , while the presence of MOOA with high oxygenation state (O : C ratio of 0.95) increased the observed κ_{org} . A similar relationship was observed in a former study at the observation site: κ_{org} was positively correlated with the O : C ratio of the organics (Deng et al., 2018).

The hygroscopicity parameters of LOOA (κ_{LOOA}) and MOOA (κ_{MOOA}) were determined to evaluate the variations in $\kappa_{\rm org}$ that can be explained by the relative contributions of LOOA and MOOA to OA. For the determination of κ_{LOOA} and κ_{MOOA} , κ_{org} was plotted against $v_{\rm LOOA}/(v_{\rm LOOA} + v_{\rm MOOA})$ and their correlation was analyzed based on linear regression analysis (Figs. 5 and S14). For particles with d_{dry} of 100, 200, 300, and 360 nm, the correlation coefficients between κ_{org} and $v_{\text{LOOA}}/(v_{\text{LOOA}}+v_{\text{MOOA}})$ were -0.50, -0.58, -0.27, and -0.099, respectively. Relatively high correlations were observed for particles with $d_{\rm drv}$ of 100 and 200 nm, probably because higher particle number concentrations (Fig. 1c) and higher OA volume fractions (Fig. 3) led to smaller uncertainties in the derived κ_{org} in those diameter ranges than in those of 300 and 360 nm. The regression line for particles with $d_{\rm drv}$ of both 100 and 200 nm (Fig. 5) were used to derive κ_{LOOA} and κ_{MOOA} by applying $v_{\text{LOOA}}/(v_{\text{LOOA}}+v_{\text{MOOA}})$ of zero and unity to the obtained regression equation, respectively. The derived κ_{LOOA}

and κ_{MOOA} were 0.083 and 0.28, respectively. This result is in between the results if particles with d_{dry} of only 100 nm (derived κ_{LOOA} and κ_{MOOA} were 0.060 and 0.25, respectively) and only 200 nm (derived κ_{LOOA} and κ_{MOOA} were 0.095 and 0.34, respectively) were used. Compared with the κ of PMF factors reported by Jimenez et al. (2009), the derived κ_{LOOA} and κ_{MOOA} are within the ranges of κ for semivolatile oxygenated OA (0.04-0.18) and low-volatility oxygenated OA (0.18-0.35), respectively (Fig. 5). The sizeresolved $v_{\text{LOOA}}/(v_{\text{LOOA}}+v_{\text{MOOA}})$ and the above derived κ_{LOOA} and κ_{MOOA} were used to reconstruct the diurnal variation data and mean values of size-resolved κ_{org} during 12:00– 20:00 JST based on the volume additivity assumption. The reconstructed κ_{org} and the κ_{org} described above (measured $\kappa_{\rm org}$) were compared and their correlations were evaluated through linear regression analysis (Fig. S15). The slope and r^2 of the regression line are used to assess the ability of LOOA and MOOA to explain the variations in κ_{org} . For 100 and 200 nm particles, the relative contribution of LOOA and MOOA to OA can explain the majority of the variations in $\kappa_{\rm org}$ (the slope and r^2 for 100 nm particles were 0.72 and 0.79, respectively, and those for 200 nm particles were 0.63 and 0.68, respectively). The variations in $\kappa_{\rm org}$ of 300 and 360 nm particles are explained less by the relative contributions of LOOA and MOOA to OA (the slope and r^2 for 300 nm particles were 0.31 and 0.48, respectively; for 360 nm particles were 0.16 and 0.07, respectively). The slope and r^2 of the regression line between reconstructed and measured $\kappa_{\rm org}$ for particles with all four $d_{\rm dry}$ were 0.39 and 0.44, respectively. This result indicates that the relative contribution of LOOA and MOOA can explain around 40 % of the observed diurnal variations in $\kappa_{\rm org}$. The slope and r^2 of the regression line over the mean reconstructed and observed κ_{org} of the four different sizes during afternoon hours (12:00-20:00 JST) were 0.39 and 0.84, respectively, which indicates that the size dependence of κ_{org} is explained by the relative contribution of LOOA and MOOA to OA by at least $\sim 40\%$ during the time period.

4.2.2 Hygroscopicity of biogenic secondary organic aerosols

The hygroscopicity parameter of freshly formed BSOA (κ_{BSOA}) was calculated as the volume-weighted mean of κ_{LOOA} and κ_{MOOA} . As discussed in Sect. 4.1.1, the enhanced OA mass in the daytime can be regarded as fresh BSOA. To simplify the analysis, the remaining part of the observed OA can be regarded as regionally transported OA (ROA), which may contain some aged, locally formed BSOA (Deng et al., 2018). To estimate the size-resolved contributions of LOOA and MOOA to BSOA and ROA, the size-resolved diurnal variations in sulfate were assumed to be a tracer of regionally transported aerosol and were scaled to represent the diurnal variations in LOOA and MOOA that constitute ROA (LOOA-ROA and MOOA-ROA;



Figure 3. Diurnal variations in (a) size-resolved hygroscopicity of aerosols (κ_t) and (b) size-resolved volume fractions of OA (ε_{org}) for the entire study period. Note that for particles with d_{dry} of 30 nm, κ_t is presented in 2 h time resolution because of the low data coverage (Fig. S12).



Figure 4. Diurnal variations in (**a**) size-resolved hygroscopicity of OA (κ_{org}) and (**b**) size-resolved volume fractions of LOOA in OA ($v_{LOOA}/(v_{LOOA} + v_{MOOA})$) for the entire study period. For κ_{org} in panel (**a**), only data with $\varepsilon_{org} > 0.40$ were considered. The values in panel (**b**) were calculated from the diurnal variations in the average volume concentrations of LOOA and MOOA (Fig. S16), not from the averages of $v_{LOOA}/(v_{LOOA} + v_{MOOA})$.

Fig. S16). For the scaling, the period of 06:00–08:00 JST, when OA and its subcomponents reached their daily minima (Figs. 2 and S16), was regarded as the background period, and all the LOOA and MOOA during the period were considered constituents of ROA. The remaining fractions of LOOA and MOOA were regarded as constituents of BSOA, referred to as LOOA-BSOA and MOOA-BSOA, respectively. The $v_{\text{LOOA}}/(v_{\text{LOOA}} + v_{\text{MOOA}})$ of ROA was estimated to be 0.52, 0.53, 0.50, and 0.45 for particles with d_{dry} of 100, 200, 300, and 360 nm, respectively (the range is presented by the area filled with a left-slash pattern in Fig. 5). The period when the diurnal variations in the sizeresolved concentration of LOOA reached their maxima (i.e., 14:00–16:00 JST; Fig. S16) was chosen to estimate the sizeresolved $v_{\text{LOOA}}/(v_{\text{LOOA}} + v_{\text{MOOA}})$ of BSOA. The estimated $v_{\text{LOOA}}/(v_{\text{LOOA}} + v_{\text{MOOA}})$ values of the BSOA were 0.97, 0.88, 0.80, and 0.79 for particles with d_{dry} of 100, 200, 300, and 360 nm, respectively (the range is presented by the area filled with a right-slash pattern in Fig. 5). Although the estimated BSOA could have aged to some extent, it was defined as fresh BSOA. The κ_{BSOA} (and κ_{ROA}) values were calculated using the derived $v_{\text{LOOA}}/(v_{\text{LOOA}} + v_{\text{MOOA}})$ of BSOA (ROA) for particles with different d_{dry} and κ_{LOOA} and κ_{MOOA} , and found to be 0.089 (0.18), 0.11 (0.18), 0.12 (0.18), and 0.12 (0.19) for particles with d_{dry} of 100, 200, 300, and 360 nm, respectively. The result indicates that κ_{BSOA} may increase with an increase in the particle diameter as a result of the size-dependent contribution of LOOA



Figure 5. κ_{org} versus $v_{\text{LOOA}}/(v_{\text{LOOA}} + v_{\text{MOOA}})$ for particles with $d_{\rm dry}$ of 100 nm (gray open squares) and 200 nm (gray cross markers) over the entire study period. The time resolution of individual data is 2 h. Only data with $\varepsilon_{org} > 0.40$ are considered. The shortdashed, long-dashed, and solid lines are the regression lines for particles with $d_{\rm drv}$ of 100, 200 nm, and the sum of the particles with the two sizes, respectively. The κ values of LOOA and MOOA derived from the regression lines (Sect. 4.2.1) are indicated by the light and dark blue pentagons, respectively. The size-resolved mean κ_{org} values during 12:00-20:00 and 20:00-12:00 JST are indicated as filled circles and diamond markers, respectively. The standard deviations of the mean κ_{org} are indicated by the whiskers. The standard deviations of the mean $v_{\text{LOOA}}/(v_{\text{LOOA}} + v_{\text{MOOA}})$ are presented in Table S6. The size-resolved κ values of BSOA and ROA are indicated by the three-pointed stars and triangles, respectively. The diameters of κ_{org} , κ_{BSOA} , and κ_{ROA} are differentiated by colors. The ranges of κ for low-volatility oxygenated OA (LVOOA) and semivolatile oxygenated OA (SVOOA) from Jimenez et al. (2009) are superimposed on the left and right axes, respectively. The shaded areas represent the estimated ranges of $v_{LOOA}/(v_{LOOA} + v_{MOOA})$ for ROA (left slash pattern) and BSOA (right slash pattern) of 100-360 nm particles.

and MOOA to BSOA (Fig. 5; colored three-pointed stars), which, however, needs to be confirmed by further studies. The size dependence of the estimated κ_{ROA} (Fig. 5; colored triangles) was less obvious than that of κ_{BSOA} . The κ_{BSOA} derived at 85% RH for particles with d_{dry} of 100 nm in this study (0.089) was slightly smaller than that in a previous study for particles with similar diameters under SUPS (0.10 at 94 ± 11 nm) at the same site (Deng et al., 2018). The derived κ_{ROA} is similar to the average κ_{org} during the nighttime (Fig. 4). The size-resolved volume concentrations of BSOA (v_{BSOA}) and ROA (v_{ROA}) were also estimated using those size-resolved $v_{\text{LOOA}}/(v_{\text{LOOA}} + v_{\text{MOOA}})$ values of BSOA and ROA (Sect. S10). The obtained volume fraction of BSOA in aerosol particles ($\varepsilon_{\text{BSOA}}$; Fig. S18) presented diurnal variation patterns that were similar to those of

 $v_{\text{LOOA}}/(v_{\text{LOOA}} + v_{\text{MOOA}})$. Furthermore, the size dependence of $\varepsilon_{\text{BSOA}}$ during afternoon hours was also similar to that of $v_{\text{LOOA}}/(v_{\text{LOOA}} + v_{\text{MOOA}})$. Because both the volume concentrations and hygroscopicity of BSOA and ROA were derived from that of LOOA and MOOA, the variation in the relative contributions of the estimated BSOA and ROA to OA can explain 40 % of the diurnal variation and size dependence of the measured κ_{org} .

4.3 Ranges of the variations in κ_t , κ_{org} , and κ_{BSOA}

The ranges of the diurnal variations in κ_t and κ_{org} , the difference between their maxima and minima, were obtained from their diurnal variation data with 2 h resolution (Table S8). The variation ranges of κ_t were 0.14, 0.091, 0.084, 0.10, 0.11, 0.11, and 0.070 for particles with d_{dry} of 30, 50, 70, 100, 200, 300, and 360 nm, respectively. The variation ranges of κ_{org} for particles with d_{dry} of 100, 200, 300, and 360 nm, respectively. The variation ranges of κ_{org} for particles with d_{dry} of 100, 200, 300, and 360 nm were 0.091, 0.079, 0.096, and 0.11, respectively. The size dependence of κ_t and κ_{org} were quantified by the mean κ_t and κ_{org} values for the entire study period (Table S8). The difference in κ_t between particles with d_{dry} of 100 and 300 nm was 0.13, and that of κ_{org} was 0.056.

The ranges of both the diurnal variations and the size dependence of κ_t are similar to those reported from a previous study at the same site in 2010 (the mean of the differences between 09:00-21:00 and 21:00-09:00 JST on NPF event days and on nonevent days for the d_{dry} range of 28.9–359 nm was in the range of 0.09-0.13, and the difference between Aitken mode and accumulation mode particles was 0.12) (Kawana et al., 2017), which pointed out the importance of the variation in particle hygroscopicity with time and size to the CCN number concentration. The ranges of both the diurnal variations and the size dependence of κ_{org} are comparable to the range of 0.05 (from 0.08 to 0.13) that could lead to 30%or more bias in the predicted CCN number concentration if not considered (Mei et al., 2013b). Here, only the ranges of the variation in κ_t and κ_{org} are discussed; other factors such as the absolute values of κ_t or κ_{org} should also be important to the prediction of CCN number concentrations. If SS is 0.1 %–1%, typical maximum values in cloud systems (Farmer et al., 2015), the d_{dry} of 100 nm is close to the mode diameters of the CCN number-size distributions in previous studies at the same site (Kawana et al., 2017; Deng et al., 2018), whereas the $d_{\rm dry}$ of 300 nm is close to the mode mobility diameters of the mass-size distributions of OA and other aerosol components in this study (Fig. S2). The difference in the two types of mode diameters indicates that significant bias could be introduced if the bulk aerosol composition and/or OA composition is used for the prediction of CCN number concentrations.

The difference in the estimated κ_{BSOA} between particles with d_{dry} of 100 and 300 nm was estimated to be 0.031. The difference implies the importance of the size dependence of

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 κ_{BSOA} in the prediction of the contribution of BSOA to the CCN number concentration.

4.4 Contributions of OA and BSOA to CCN concentrations

The contributions of OA and BSOA to CCN number concentrations were assessed from the viewpoint of their contributions to the aerosol water uptake, which are size-dependent. For the estimate, the observed aerosols were assumed to be internally mixed. This is supported by a result from a previous study at the observation site: there was almost no difference in the prediction of the number fractions of CCN between the use of time- and size-resolved $g_{\rm f}$ distributions and time- and size-resolved $g_{f,m}$ (Kawana et al., 2017). The CCN activation diameters (d_{act}) for respective 2 h time-of-day bins at 0.4 % SS were calculated from reconstructed hygroscopicity parameter values of aerosol particles based on κ -Köhler theory (Sect. S11). The estimated total CCN number concentration is referred to as $N_{\text{CCN},t}$. For each size range, the contribution of OA (BSOA) to the aerosol water uptake was represented as the product of the volume fraction of OA (BSOA) and κ_{org} (κ_{BSOA}) divided by κ_t (i.e., $\varepsilon_{\text{org}}\kappa_{\text{org}}/\kappa_t$ $(\varepsilon_{\text{BSOA}}\kappa_{\text{BSOA}}/\kappa_t))$ and was used to represent the fractional contribution of OA (and BSOA) to $dN_{CN}/d\log d_{dry}$ in the size range. The fractional contribution of OA (and BSOA) to $N_{\text{CCN},t}$, hereafter referred to as $F_{\text{CCN},\text{OA}}$ ($F_{\text{CCN},\text{BSOA}}$), was derived by integrating the product of $dN_{CN}/d\log d_{dry}$ and $\varepsilon_{\rm org} \kappa_{\rm org} / \kappa_t$ ($\varepsilon_{\rm BSOA} \kappa_{\rm BSOA} / \kappa_t$) above the $d_{\rm act}$ and by dividing the obtained value by $N_{\text{CCN},t}$. Details for the estimation of $F_{\text{CCN,OA}}$ and $F_{\text{CCN,BSOA}}$ are presented in Sect. S11.

The diurnal variation in $F_{\rm CCN,OA}$ estimated using timeand size-resolved $\kappa_{\rm org}$ and that of $F_{\rm CCN,BSOA}$ using sizeresolved $k_{\rm BSOA}$ are presented in Fig. 6. Both $F_{\rm CCN,OA}$ and $F_{\rm CCN,BSOA}$ reached their maxima during 12:00–16:00 JST, when intensive BSOA formation was observed. The magnitude of the variation in $F_{\rm CCN,BSOA}$ (from 0.00 to 0.28) was larger than that of $F_{\rm CCN,OA}$ (from 0.40 to 0.52). This is explained by the larger magnitude of the diurnal variation range of $\varepsilon_{\rm BSOA}$ (Fig. S18), compared to that of $\varepsilon_{\rm org}$ (Fig. 3). The $F_{\rm CCN,BSOA}$ of 0.28 during 14:00–16:00 JST indicates a significant contribution of BSOA to the CCN number concentration.

Because obvious diurnal variations and size dependence of κ_{org} were found and because κ_{BSOA} was also estimated to be size-dependent (Sect. 4.2), the sensitivities of the estimated $F_{CCN,OA}$ and $F_{CCN,BSOA}$ on the variations in κ_{org} and κ_{BSOA} were assessed. To assess the influence of the variation in κ_{org} on $F_{CCN,OA}$, the diurnal variations in $F_{CCN,OA}$ were estimated using (0a, base case) time- and size-resolved κ_{org} ; (1a) size-resolved, time-averaged (note that the average here refers to the arithmetic mean, and it is the same in other places of this paragraph) κ_{org} ; (2a) time-resolved, size-averaged κ_{org} ; and (3a) time- and size-averaged κ_{org} (Table S10). Cases (1a)–(3a) were compared with the base



Figure 6. Diurnal variation in the fractional contribution of OA to the total CCN number concentration ($F_{\text{CCN,OA}}$) estimated using time- and size-resolved κ_{org} , and diurnal variation in the fractional contribution of BSOA to the total CCN number concentration ($F_{\text{CCN,BSOA}}$) estimated assuming fresh BSOA (using size-resolved κ_{BSOA}) and aged BSOA (using size-resolved κ_{ROA}).

case (0a), as presented in Fig. 7a. Using time-averaged $\kappa_{\rm org}$ (case 1a), the $F_{\text{CCN,OA}}$ was overestimated by 18% during 16:00-18:00 JST and underestimated by 13 % during 06:00–08:00 JST. Using size-averaged $\kappa_{\rm org}$ (case 2a), the F_{CCN,OA} was overestimated by 2%–19% on a diurnal basis. Using time- and size-averaged κ_{org} (i.e., a single mean $\kappa_{\rm org}$, case 3a), $F_{\rm CCN,OA}$ was overestimated by 26 % during 16:00–18:00 JST and underestimated by 5.0 % during 06:00– 08:00 JST. The deviation of case (3a) from the base case (0a) resulted from the factors leading to the deviations of cases (1a) and (2a). The magnitudes of the deviations, defined here as the difference between the lowest and highest values of the ratios in Fig. 7a, for (1a), (2a), and (3a) are 31 %, 17 %, and 31 %, respectively. The substantial differences suggest that the diurnal variations and size dependence of κ_{org} are important for accurate prediction of the contribution of OA to the CCN number concentration in modeling studies. To assess the influence of the size dependence of κ_{BSOA} on $F_{\text{CCN,BSOA}}$, the diurnal variations in $F_{\text{CCN,BSOA}}$ was estimated using (0b, base case) size-resolved κ_{BSOA} and (1b) size-averaged κ_{BSOA} (Table S11). Case (1b) was compared with case (0b), as presented in Fig. 7b. Using size-averaged $\kappa_{\rm BSOA}$ caused overestimation of $F_{\rm CCN,BSOA}$ by 12%–19%, which relates to the decrease in the estimated κ_{BSOA} and the increase in $dN_{CN}/d\log d_{drv}$ (Fig. S17) with decrease in the dry particle diameter.

Furthermore, because fresh BSOA probably become aged after atmospheric transport, the influence of the aging of the estimated fresh BSOA (assuming κ_{BSOA} for the calculation of d_{act} and $F_{CCN,BSOA}$ was as large as that of κ_{ROA} , Table S11) on $F_{CCN,BSOA}$ was also evaluated. Here, the estimation of $F_{CCN,BSOA}$ in the aged condition ignored the possible change in the aerosol size distribution accompanying the aging process. Aged BSOA can contribute more to the aerosol water uptake and thus to the CCN number concentra-



Figure 7. (a) Diurnal variation in the ratios of the $F_{\text{CCN},\text{OA}}$ derived using time- and size-resolved κ_{org} (TimeSizeReso κ_{org}), time-averaged and size-resolved κ_{org} (SizeReso κ_{org}), size-averaged and time-resolved κ_{org} (TimeReso κ_{org}), and time- and size-averaged κ_{org} (Single κ_{org}) to that derived using the time- and size- resolved κ_{org} (TimeSizeReso κ_{org}). (b) Diurnal variation in the ratios of the $F_{\text{CCN},\text{BSOA}}$ derived using size-resolved κ_{BSOA} (SizeReso κ_{BSOA}), size-averaged κ_{BSOA} (SizeReso κ_{BSOA}), and aged size-resolved κ_{BSOA} (AgedSizeReso κ_{BSOA}) to that derived using the size-resolved κ_{BSOA} (SizeReso κ_{BSOA}). In panel (b), the condition of aged size-resolved κ_{BSOA} assumes that the value of κ_{BSOA} equals that of κ_{ROA} , and the data during 04:00–08:00 JST, when the concentration of BSOA was low (volume concentration less than $0.01 \times 10^{-6} \text{ cm}^3 \text{ m}^{-3}$), are not presented (data are presented in Table S12).

tion. Assuming that the BSOA was as aged as ROA, the estimated $F_{\text{CCN,BSOA}}$ was 52 %–84 % larger than that estimated assuming fresh BSOA (Fig. 7b), and it could have been 0.42 if the aerosol observed during 14:00–16:00 JST aged (Fig. 6). The result suggests that, whereas the contribution of BSOA to CCN was substantial at the study site, the magnitude of the contribution might be increased substantially by aging of the BSOA during transport after its formation in the forest.

5 Summary and conclusions

The size-resolved hygroscopicity at 85% RH, chemical composition, and number–size distributions of atmospheric aerosols were observed at a forest site in Wakayama, Japan, in August and September 2015. The diurnal variation and size dependence in the hygroscopicity of the observed aerosol and organic aerosol (OA) components were discussed in view of the formation of BSOA. The fractional contributions of OA and BSOA to the total CCN number concentration were discussed in view of the variations in the hygroscopicity parameter of OA and BSOA.

Similar to two previous observations at the same site (Han et al., 2013, 2014; Kawana et al., 2017; Deng et al., 2018), OA was the dominant sub-micrometer aerosol component, followed by sulfate. While the mass concentration of sulfate, on average, did not vary much in a day, the mass concentration of OA increased substantially in the afternoon hours, which was presumably explained by the condensation of BSOA. The hygroscopicity of ambient aerosol (κ_t) and of OA (κ_{org}) increased with an increase in the dry particle diameter and presented daily minima in the afternoon hours. In this study, the ranges of the diurnal variations in κ_{org} of 100–360 nm particles were 0.079–0.11 and the κ_{org} of 300 nm particles. The

diurnal variations and size dependence of κ_t can be explained by the relative contributions of OA and inorganic salts in the observed aerosol. The relative contributions of the estimated fresh BSOA and regional OA can explain 40 % of the diurnal variation and size dependence of $\kappa_{\rm org}$. The hygroscopicity of fresh BSOA ($\kappa_{\rm BSOA}$) was estimated to increase (0.089–0.12) with an increase in the dry particle diameter (100–300 nm).

The fractional contributions of OA and fresh BSOA to CCN number concentrations, estimated from the viewpoint of their contributions to the water uptake by the aerosol, were in the ranges 0.40–0.52 and 0.00–0.28, respectively. Compared with the use of time- and size-resolved κ_{org} , the use of time- and size-averaged κ_{org} overestimated the contribution of OA to the CCN number concentration by up to 26 % (16:00-18:00 JST) and underestimated the contribution by up to 5.0% (06:00–08:00 JST). These results indicate the importance of the diurnal variations and size dependence of $\kappa_{\rm org}$ in the prediction of the contribution of OA to the CCN number concentration. The use of size-averaged κ_{BSOA} overestimated the contribution of fresh BSOA to the CCN number concentration by 12 %-19 % compared with the use of size-resolved κ_{BSOA} . If aging of BSOA after atmospheric transport occurs, the contribution of fresh BSOA to the CCN number concentration could be increased by 52 %-84 %, and could have reached a high value of 0.42 if the aerosol observed during 14:00–16:00 JST aged.

This study revealed the large magnitude of the diurnal variation and size dependence of κ_{org} at the observation site under the influence of the formation of BSOA. Also revealed was the importance of the variation in κ_{org} to the estimation of the contribution of OA to the CCN number concentration from the viewpoint of the size-resolved contribution of OA to the water uptake of aerosols. Because both the diurnal variation and size dependence of κ_{org} in the studied forest are dif-

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ferent from those in some other forest environments (Cerully et al., 2015; Thalman et al., 2017), further studies on the variation in the hygroscopicity of organics and on the contributions of OA and BSOA to the CCN concentrations should be performed in other forest environments. Furthermore, the size dependence of the hygroscopicity of fresh BSOA estimated here should be confirmed by additional studies.

Data availability. Raw data used to derive the data presented are available on request to the corresponding author. All of the final derived data supporting the findings of this study are available in the article or in the Supplement.

Supplement. The supplement related to this article is available online at: https://doi.org/10.5194/acp-19-5889-2019-supplement.

Author contributions. MM and YD designed the experiments, and YD, HY, MM, HF, and TN performed them. YD analyzed the data with contributions from MM, TN, and KK. YD prepared the manuscript with contributions from MM, KK, and TN.

Competing interests. The authors declare that they have no conflict of interest.

Acknowledgements. We thank the faculty and staff of the Wakayama Forest Research Station, Field Science Education and Research Center of Kyoto University, Japan, for the provision of the study site and the meteorological data. We thank Qingcai Chen, Kouji Adachi, Yuuki Kuruma, Ryuji Fujimori, and Takayuki Yamasaki for their help in the field observation. We acknowledge Kazuma Aoki for the use of the 1λ -PSAP instrument. We acknowledge the NOAA Air Resources Laboratory (ARL) for providing the HYSPLIT transport and dispersion model. This study was supported in part by JSPS KAKENHI grant numbers JP26281007 and JP18K19852.

Review statement. This paper was edited by Veli-Matti Kerminen and reviewed by three anonymous referees.

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