



# Potential of polarization lidar to provide profiles of CCN- and INP-relevant aerosol parameters

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**Abstract.** We investigate the potential of polarization lidar to provide vertical profiles of aerosol parameters from which cloud condensation nucleus (CCN) and ice nucleating particle (INP) number concentrations can be estimated. We show that height profiles of particle number concentrations  $n_{50,\text{dry}}$  considering dry aerosol particles with radius  $> 50$  nm (reservoir of CCN in the case of marine and continental non-desert aerosols),  $n_{100,\text{dry}}$  (particles with dry radius  $> 100$  nm, reservoir of desert dust CCN), and of  $n_{250,\text{dry}}$  (particles with dry radius  $> 250$  nm, reservoir of favorable INP), as well as profiles of the particle surface area concentration  $s_{\text{dry}}$  (used in INP parameterizations) can be retrieved from lidar-derived aerosol extinction coefficients  $\sigma$  with relative uncertainties of a factor of 1.5–2 in the case of  $n_{50,\text{dry}}$  and  $n_{100,\text{dry}}$  and of about 25–50% in the case of  $n_{250,\text{dry}}$  and  $s_{\text{dry}}$ . Of key importance is the potential of polarization lidar to distinguish and separate the optical properties of desert aerosols from non-desert aerosol such as continental and marine particles. We investigate the relationship between  $\sigma$ , measured at ambient atmospheric conditions, and  $n_{50,\text{dry}}$  for marine and continental aerosols,  $n_{100,\text{dry}}$  for desert dust particles, and  $n_{250,\text{dry}}$  and  $s_{\text{dry}}$  for three aerosol types (desert, non-desert continental, marine) and for the main lidar wavelengths of 355, 532, and 1064 nm. Our study is based on multiyear Aerosol Robotic Network (AERONET) photometer observations of aerosol optical thickness and column-integrated particle size distribution at Leipzig, Germany, and Limassol, Cyprus, which cover all realistic aerosol mixtures. We further include AERONET data from field campaigns in Morocco, Cabo Verde, and Barbados, which provide pure dust and pure

marine aerosol scenarios. By means of a simple CCN parameterization (with  $n_{50,\text{dry}}$  or  $n_{100,\text{dry}}$  as input) and available INP parameterization schemes (with  $n_{250,\text{dry}}$  and  $s_{\text{dry}}$  as input) we finally compute profiles of the CCN-relevant particle number concentration  $n_{\text{CCN}}$  and the INP number concentration  $n_{\text{INP}}$ . We apply the method to a lidar observation of a heavy dust outbreak crossing Cyprus and a case dominated by continental aerosol pollution.

## 1 Introduction

Field studies of aerosol–cloud–dynamics interaction are presently in the focus of atmospheric research. Large uncertainties in weather and future-climate predictions (IPCC, 2013) arise from gaps in our knowledge of the detailed impact of aerosols on the evolution of liquid-water, mixed-phase and cirrus clouds. This unsatisfactory situation motivates the strong efforts presently undertaken to investigate formation and evolution of cloud layers and associated aerosol–cloud interactions.

Aerosol particles influence cloud evolution, lifetime, and cloud microphysical properties in two ways. Aerosol particles can serve as cloud condensation nuclei (CCN) in liquid droplet nucleation processes and/or as ice-nucleating particles (INP) in ice nucleation processes, which include the conversion of liquid droplets into ice crystals (immersion freezing). Ground-based active remote sensing (lidar and radar observations) can be used to continuously monitor the evolution of clouds in their natural environment, at given meteorological

logical conditions with high vertical and temporal resolution (Illingworth et al., 2007; Shupe, 2007; Ansmann et al., 2009; de Boer et al., 2011; Schmidt et al., 2014).

Lidar is the most prominent tool for aerosol profiling in terms of particle optical properties. However, to improve the study of aerosol–cloud interaction, the potential of lidar to provide vertical profiles of particle number concentrations such as  $n_{50,\text{dry}}$  (considering all dry particles with radius  $> 50$  nm (reservoir of favorable CCN in the case of marine and anthropogenic particles) (Quinn et al., 2008; Rose et al., 2010; Deng et al., 2011),  $n_{100,\text{dry}}$  (dry particles with radius  $> 100$  nm, reservoir of favorable CCN in the case of desert dust) (Koehler et al., 2009; Kumar et al., 2009, 2011), or of the large particle fraction  $n_{250,\text{dry}}$  (all particles with dry radius  $> 250$  nm, reservoir of favorable INP) (DeMott et al., 2010, 2015a), needs to be explored in detail. The central question of our study is can we use lidar-derived vertical profiles of aerosol backscatter coefficient  $\beta$  and extinction coefficient  $\sigma$ , measured at ambient relative humidity conditions, to estimate vertical profiles of dry particle number concentration  $n_{\text{dry}}$  and surface area concentration  $s_{\text{dry}}$  from which the cloud-relevant particle number concentrations  $n_{\text{CCN}}$  (indicating the CCN particle reservoir) and  $n_{\text{INP}}$  (INP number concentration) can be estimated?

A first promising feasibility study regarding the retrieval of  $n_{\text{INP}}$  profiles from lidar observations was undertaken by Mamouri and Ansmann (2015). Former studies indicate also that measured aerosol optical properties (at wavelengths around 500 nm) can be used to estimate CCN number concentrations, CCNC (Ghan and Collins, 2004; Ghan et al., 2006; Andreae, 2009; Jefferson, 2010; Liu and Li, 2014; Shinzuka et al., 2015).

A crucial point regarding  $n_{\text{CCN}}$  and  $n_{\text{INP}}$  profiling is that the efficacy of aerosol particles to act as CCN or INP depends on aerosol type. In the case of heterogeneous ice nucleation it is found that mineral dust particles are favorable INP at temperatures below about  $-20^\circ\text{C}$  (Ansmann et al., 2009; Murray et al., 2012; Augustin-Bauditz et al., 2014) that marine particles seem to be comparably inefficient INPs (Kanitz et al., 2011) at temperatures  $> -25^\circ\text{C}$ , whereas continental aerosols (mixtures of anthropogenic haze, biomass burning smoke, soil and road dust, and organic and biogenic particles from soils and plants) seem to contain always a significant amount of efficient INPs, already leading to ice nucleation at temperatures as high as  $-5$  to  $-15^\circ\text{C}$  (Seifert et al., 2010; Zhang et al., 2010; Kamphus et al., 2010; Ebert et al., 2011; Augustin et al., 2013; Hartmann et al., 2013; Bühl et al., 2013; Pummer et al., 2015; Umo et al., 2015).

In the case of cloud droplet formation, we have to distinguish at least desert dust and non-desert particles (continental and marine aerosol components) (Koehler et al., 2009; Kumar et al., 2009, 2011; Karydis et al., 2011; Bangert et al., 2012). Marine and hygroscopic continental particles with dry radius  $> 50$  nm get activated even at low supersaturation of 0.1–0.2 % (i.e., at relative humidities over liquid wa-

ter of 100.1 to 100.2 %), whereas the critical activation radius of hydrophobic insoluble desert particles with a negligible amount of soluble material (coating) on the surface is  $> 100$  nm (Koehler et al., 2009; Kumar et al., 2011). Thus, lidar must be able to separate these basic aerosol types and to provide  $n_{\text{CCN}}$  and  $n_{\text{INP}}$  profiles separately for marine, non-desert continental, and desert dust aerosols.

In principle, multiwavelength Raman/polarization or high-spectral-resolution (HSR)/polarization lidars can provide the desired microphysical particle properties (Müller et al., 2005; Veselovskii et al., 2010; Müller et al., 2013, 2014). However comparably complex lidars and comprehensive data analysis methods as well as a good knowledge in the use of ill-posed inversion techniques are required to make these efforts successful. For this reason, we investigate an alternative approach. The overall goal is to develop a robust and easy-to-apply method that allows fast computation and implementation of an automated code in the lidar aerosol and cloud data analysis software. Thus, the method should be simple and applicable to single-wavelength lidar observations at 355, 532, or 1064 nm wavelength to estimate profiles of particle number concentrations  $n_{50,\text{dry}}$ ,  $n_{100,\text{dry}}$ ,  $n_{250,\text{dry}}$ , and surface area concentration  $s_{\text{dry}}$  for the three basic aerosol types. Many lidars are single-wavelength lidars (e.g., 355 or 532 nm backscatter lidars) including the upcoming space lidars of the European Space Agency operating at 355 nm (Ansmann et al., 2007; Illingworth et al., 2015a) which are planned to be launched within the next 1–3 years. Furthermore, a dense European single-wavelength ceilometer network is developing, organized by European weather services (Wiegner and Geiß, 2012; Wiegner et al., 2014; Illingworth et al., 2015b).

To make full use of the retrieval schemes presented in this article, polarization lidars (Freudenthaler et al., 2009) are of advantage. This is a key point of the entire study. By means of the polarization lidar technique, the desert dust aerosol component can be easily separated from other continental aerosol components as well as from marine aerosol. Desert dust causes high depolarization of backscattered linearly polarized laser light, whereas typical non-desert aerosol mixtures lead to very low depolarization. After the separation of the basic aerosol types, in the next step the particle number and surface area concentrations, required as input in the CCN and INP parameterization schemes, are separately determined from the lidar-derived particle extinction coefficients for the basic aerosol types (desert, marine, continental), as outlined in Sects. 3 and 4.

The study presented here is based on our long experience in detection, separation, and quantification of optical and microphysical properties of different aerosol types by using polarization lidars in combination with sun photometers (Tesche et al., 2009, 2011; Ansmann et al., 2011b, 2012; Mamouri et al., 2013; Mamouri and Ansmann, 2014; Nisantzi et al., 2014, 2015). This study can be regarded as a follow-up effort of Mamouri and Ansmann (2015). However, in a much broader and more general sense, we now

illuminate the potential of lidar to provide cloud-formation-relevant aerosol parameters for both liquid-water droplet and ice crystal nucleation. New aspects deal with the estimation of  $n_{50,\text{dry}}$  and  $n_{100,\text{dry}}$ , the CCN parameterization, the retrieval of the particle surface area concentration  $s_{\text{dry}}$  from measured particle extinction coefficients, and the consideration of further dust INP parameterizations developed by Niemand et al. (2012) and Steinke et al. (2015), in which the dust values of  $s_{\text{dry}}$  are input instead of  $n_{250,\text{dry}}$  (DeMott et al., 2010, 2015a). In addition, new findings regarding the efficacy of marine particles to serve as INP are taken into account (DeMott et al., 2015b). In the present study, the wavelength range is extended from 532 nm to all three relevant laser wavelengths so that the CCN and INP-relevant aerosol conversion parameters are available for 355 and 1064 nm as well.

The study makes use of multiyear photometer observations of the Aerosol Robotic Network (AERONET) (Holben et al., 1998) at Leipzig, Germany, Limassol, Cyprus, and at Ragged Point, Barbados. We further include AERONET data from desert dust field campaigns in Morocco, Cabo Verde, and Barbados. The main goal is to investigate the link between the microphysical particle properties such as  $n_{50,\text{dry}}$ ,  $n_{100,\text{dry}}$ ,  $n_{250,\text{dry}}$ , and  $s_{\text{dry}}$  and the ambient particle extinction coefficient, measurable with lidar, for “real-world” aerosol conditions. Long-term AERONET observations reflect best the full range of occurring aerosol mixture and layering scenarios. An alternative approach would be an extended simulation study of the correlation between the cloud-relevant microphysical and measurable optical properties of the aerosol, similar to the study presented by Barnaba and Gobbi (2001) for marine and dust aerosols.

The paper is organized as follows: The AERONET stations and measurement products as well as the lidar sites and lidar products are presented in Sect. 2. Section 3 presents our methodology to obtain profiles of  $n_{50,\text{dry}}$ ,  $n_{100,\text{dry}}$ ,  $n_{250,\text{dry}}$ ,  $s_{\text{dry}}$ ,  $n_{\text{CCN}}$ , and  $n_{\text{INP}}$  from lidar profiles of ambient particle extinction coefficients  $\sigma$  for the three basic aerosol types (desert, marine, continental). The conversion of measured optical properties into particle number and surface area concentrations requires good knowledge of the correlation between optical and microphysical particle properties. This knowledge is gained from the mentioned long-term AERONET measurements and the specific dust field campaigns. The main findings of the AERONET-based correlation studies are presented and discussed in Sect. 4. Section 5 finally deals with the application of the developed methods to two lidar observations conducted during a strong desert dust outbreak towards Cyprus and during conditions with (non-desert) continental aerosol pollution over Cyprus. A summary and concluding remarks are in Sect. 6

## 2 Instrumentation

In Sect. 2.1, we provide an overview of the AERONET stations, the basic AERONET products, and the retrieval of the column-integrated particle number and surface area concentrations. In Sect. 2.2, we briefly describe our lidar instruments.

### 2.1 AERONET sun/sky photometers

The study is based on the analysis of three long-term and four field-campaign AERONET data sets. We investigated 14 years of AERONET observations at Leipzig, Germany, performed by the Leibniz Institute for Tropospheric Research (TROPOS) from 2001–2015. Leipzig is a highly polluted central European city which is affected by Saharan dust outbreaks about 2–10 times per year (Mattis et al., 2004, 2008). We analyzed 4 years of AERONET observations at Limassol, Cyprus, performed by the Cyprus University of Technology (CUT) from 2011 to 2015 (Nisantzi et al., 2014, 2015). This site in the Eastern Mediterranean is a unique station for aerosol studies. Aerosol mixtures of anthropogenic haze, biomass burning smoke, soil and road dust, and marine particles, and strong dust outbreaks from Middle East deserts and the Sahara frequently occur (Nisantzi et al., 2015). Our studies are complemented by AERONET observations conducted during the Saharan Mineral Dust Experiments SAMUM-1 (Ouarzazate, Morocco) (Toledano et al., 2009) and SAMUM-2 (Praia, Cabo Verde) (Toledano et al., 2011; Ansmann et al., 2011a), the Saharan Aerosol Long-range Transport and Aerosol-Cloud interaction Experiments SALTRACE-1 (at the Caribbean Institute for Meteorology and Hydrology (CIMH), Barbados, summer 2013) (Groß et al., 2015) and during SALTRACE-3 (Barbados, summer 2014) (Haarig et al., 2015). The field campaigns offer the unique opportunity to study the correlation between the particle optical properties (extinction coefficient  $\sigma$ , aerosol optical depth AOT) and the microphysical properties (column or layer mean values of, e.g.,  $n_{100}$ ,  $n_{250}$ ,  $s$ ) at pure dust conditions. During SALTRACE in 2013, even aircraft observations of CCNC in lofted dust layers in the Barbados area are available and the link between the in situ-measured CCNC and the lidar-derived particle extinction coefficients will be discussed in a follow-up paper. Furthermore, we used 7.5 years of data from the AERONET station at Ragged Point, Barbados (level 2.0, 2007–2015) (Prospero and Mayol-Bracero, 2013) to study the correlation between the optical and microphysical aerosol properties for pure marine conditions. An overview of the observational periods and amount of available data for the analyzed different aerosol conditions with focus on the three defined aerosol types is given in Table 1. More details of these AERONET stations can be found on the AERONET web page (<http://aeronet.gsfc.nasa.gov>).

**Table 1.** Available AERONET data sets (individual observations) of particle optical and microphysical properties for the three defined basic aerosol types. A total number of 1745 level-2.0 data sets with AOT and column size distributions for Limassol and of 2157 data sets for Leipzig are available for our correlation study. 125 data sets of desert dust optical properties and inverted particle size distributions are selected from the SAMUM and SALTRACE field campaign observations. 123 respective Ragged Point observations (Barbados) for pure marine conditions could be analyzed for our study. CIMH stands for Caribbean Institute for Meteorology and Hydrology.

Site	Observational period	Observations	Dominating aerosol type
Limassol, Cyprus (CUT-TEPAK, 34.7° N, 33.0° E, 25 m a.s.l.)	Jul 2011–Jun 2015	421	continental aerosol, AE > 1.6
		134	desert dust, AE < 0.5
Leipzig, Germany (TROPOS, 51.4° N, 12.4° E, 125 m a.s.l.)	May 2001–Jun 2015	974	continental aerosol, AE > 1.6
		33	desert dust, AE < 0.5
Ouarzazate, Morocco (SAMUM-1, 30.9° N, 6.9° W, 1150 m a.s.l.)	May–Jun 2006	32	desert dust
Praia, Cabo Verde (SAMUM-2, 14.9° N, 23.4° W, 70 m a.s.l.)	Jan 2008	23	desert dust
Barbados (SALTRACE-1, CIMH, 13.1° N, 59.6° W, 110 m a.s.l.)	Jun–Jul 2013	20	desert dust
Barbados (SALTRACE-3, Ragged Point, 13.2° N, 59.4° W, 40 m a.s.l.)	Jun–Jul 2014	50	desert dust
Barbados (Ragged Point)	Aug 2007–Feb 2015	123	marine aerosol

AERONET provides quality-assured products in terms of AOTs at up to 8 wavelengths (340 to 1640 nm) and column-integrated values for the Ångström exponent (AE, spectral dependence of AOT). The column-integrated particle size distribution is retrieved from the AOT measurements and sky radiance observations at four wavelengths (Dubovik and King, 2000; Dubovik et al., 2006); this allows us to compute particle volume concentration, surface area (column  $s$ ), and column-integrated particle number concentrations (column  $n$ ). All observations are performed at ambient temperature and relative humidity conditions. In Sect. 3.2, we will explain how we corrected for the particle water-uptake effect to obtain the required dry particle values, i.e., of  $n_{50,\text{dry}}$ ,  $n_{100,\text{dry}}$ ,  $n_{250,\text{dry}}$ , and  $s_{\text{dry}}$ .

As explained in the methodology Sect. 3, in the retrieval of  $n_{\text{CCN}}$  and  $n_{\text{INP}}$  we need to know the relationship between the observed (ambient) microphysical particle properties  $n_{50}$ ,  $n_{60}$ ,  $n_{100}$ ,  $n_{250}$ ,  $n_{290}$ ,  $n_{500}$ , and  $s$  and the ambient particle extinction coefficient  $\sigma$  for desert dust, marine, and non-desert continental aerosol conditions. These relationships are quantified by means of the AERONET correlation studies for the particle extinction coefficients at 355, 532, and 1064 nm (Sect. 4). Because AERONET photometers do not directly measure AOTs at the laser wavelengths, we use the measured AOT at 380 nm and the Ångström exponent AE (340–380 nm) to obtain the AOT at 355 nm by interpolation. Similarly, in the case of 532 nm we use the measured AOT at 500 nm and AE (440–870 nm) to derive the 532 nm AOT. The AOT at 1064 nm is obtained by extrapolation based on the measured AOT at 1020 nm and AE (870–1020 nm).

The way to obtain the column-integrated particle number concentrations, e.g., the column values of  $n_{50}$  or  $n_{250}$ ,

from the basic AERONET information (column-integrated particle volume size distribution) is described in detail by Mamouri and Ansmann (2015) (see Sect. 3.2 and Fig. 3 in that article). The particle volume size distribution is retrieved for 22 logarithmically equidistant discrete radius points  $r_j$  with index  $j$  from 1 to 22 (Dubovik and King, 2000; Dubovik et al., 2006). The particle radius spectrum from  $r_1 = 0.05$  to  $r_{22} = 15 \mu\text{m}$  is covered. Each radius  $r_j$  represents a radius interval of logarithmically equal width. To obtain the particle number concentration for each individual radius interval, we divide the determined volume concentration of a given radius interval (or for the discrete radius point  $r_j$ ) by the volume of a single particle with radius  $r_j$  and multiply this ratio with the spectral integral width of 0.2716. Unfortunately, we left out this multiplication with the dimensionless spectral width in the foregoing paper (Mamouri and Ansmann, 2015) so that the presented number concentrations in Mamouri and Ansmann (2015) are a factor of  $1/0.2716$  ( $= 3.68$ ) too high and also the respective conversion factor in Fig. 4 of that paper.

As outlined in Sect. 4, we need the column values of  $n_{50}$ ,  $n_{60}$ ,  $n_{100}$ ,  $n_{250}$ ,  $n_{290}$ ,  $n_{500}$ , and  $s$ . The column value of the particle number concentration  $n_{50}$  is the sum of the number concentrations of all radius classes from 1–22 and thus covers the full size range of optically active particles. Similarly, the column  $n_{60}$  value is obtained by adding all particle number concentrations of the radius classes from 2–22 (particles with radius > 57 nm). The column value of  $n_{100}$  considers the radius classes 4–22 (particles with radius > 98 nm).

The INP-relevant column value of  $n_{250}$  is calculated as follows: This number concentration is the sum of the number concentrations of the radius intervals 8–22 plus an additional contribution by radius interval 7 (centered at  $r_7 = 255$  nm).

This additional contribution is obtained by calculating the mean number concentration of the two intervals 7 and 8 (centered at  $r_8 = 335$  nm), assuming that this mean value represents the number concentration for the radius interval from 255 to 335 nm (centered at about 290 nm), and then taking 50 % of the computed mean value to consider only one half of this size interval. This latter value is interpreted as the number concentration of particles with radius from about 250 to about 290 nm. Furthermore, we make use in Sect. 4 of  $n_{290}$  (radius classes 8–22, particles with radius  $>$  about 290 nm), and  $n_{500}$  (radius classes 10–22, all particles with radius  $>$  about 500 nm).

The total particle surface area concentration  $s$  is obtained by (a) computing the surface area of a sphere with radius  $r_j$  for all 22 radius intervals, (b) multiplying the obtained surface areas for the particles with radius  $r_j$  with the number concentrations of radius interval  $j$  (obtained from the foregoing calculations of  $n$ ), and (c) calculating the total surface area concentration by adding all contributions of the 22 size classes up. According to airborne in situ observations of the particle size distribution during the SAMUM campaigns (Weinzierl et al., 2009), the AERONET-derived values of  $s$  for desert environments explain about 95 % of the total surface area concentration (which includes particles with radius  $<$  50 nm). By inspection of all ground-based in situ measured size distribution at the urban site of Leipzig, taken during the full year of 2008, we found that  $s$  (from AERONET) is about 0.85 ( $\pm 0.1$ ) of the total  $s$ .

Dubovik et al. (2000) carried out a detailed analysis of uncertainties in the AERONET products. Caused by statistical signal noise, the uncertainties in the AERONET  $n_{50}$ ,  $n_{60}$ , and  $n_{100}$  values can be as high as 20 %. For the column values of  $n_{250}$ ,  $n_{290}$ ,  $n_{500}$ , and  $s$ , the uncertainties are around 10 %. Offset errors (caused, e.g., by bad photometer pointing stability, by the use of wrong surface reflectance in the data analysis, and wrong AOT retrievals) can lead to extreme errors of the order of  $>$  50 % for the column  $n$  and  $s$  values in individual observations. On average, uncertainties of 25–35 % are expected. However, in the case of our multiyear AERONET observations with many calibration sessions (photometer calibrations in France or USA) and field campaign measurements with calibration session before and after the campaigns, strong biases and extreme uncertainties in our AERONET data sets can be ruled out. We assume in the following that mean uncertainties in the used long-term and field-campaign mean values of the column values of  $n_{50}$ ,  $n_{60}$ ,  $n_{100}$ ,  $n_{250}$ ,  $n_{290}$ ,  $n_{500}$ , and  $s$  values are about 10–20 %.

## 2.2 Aerosol lidars

The AERONET station of CUT at Limassol is equipped with a polarization/Raman lidar and belongs to the European Aerosol Research Lidar Network EARLINET (Pappalardo et al., 2014). The CUT lidar is described by Mamouri et al.

(2013) and Nisantzi et al. (2015). The case study in Sect. 5.2 is based on the lidar observations at Limassol.

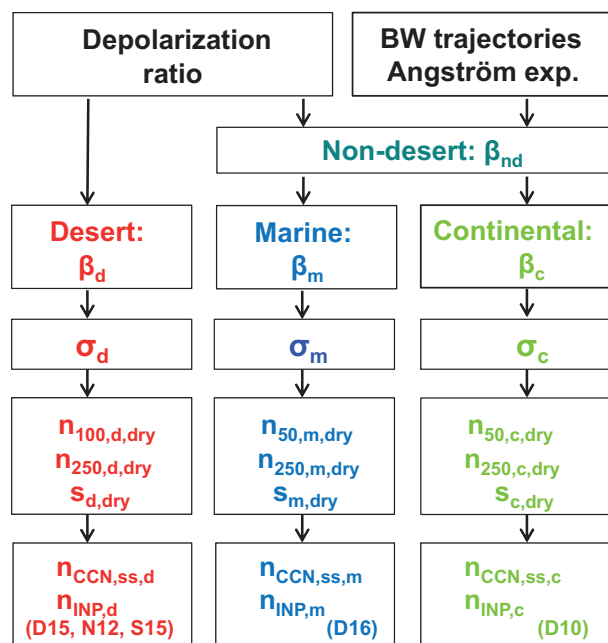
In Sect. 5.1, we discuss a lidar measurement obtained with a mobile system of the Polly<sup>XT</sup> (Portable Lidar System, XT: extended version) series (Engelmann et al., 2016; Baars et al., 2016). This new Polly<sup>XT</sup> was built by TROPOS for the National Observatory Athens (NOA) in 2014. The multiwavelength Raman/polarization lidar was extensively tested and characterized at Nicosia, Cyprus, during a 6-week field campaign in March–April 2015. The field campaign was performed in the framework of the BACCHUS (impact of Biogenic vs. Anthropogenic emissions on Clouds and Climate: towards a Holistic UnderStAnding, [www.bacchus-env.eu](http://www.bacchus-env.eu)) project. BACCHUS is a European Union's Seventh Framework Programme for Research (FP7) collaborative project of 20 institutes (including CUT and TROPOS), coordinated by ETH Zurich, Switzerland. The BACCHUS Cyprus 2015 field campaign focussed on ground-based and airborne in situ observations of  $n_{\text{INP}}$  and comparison of these observations with lidar-derived  $n_{\text{INP}}$  profiles.

The retrieval of the basic lidar products (height profiles of particle backscatter and extinction coefficients) is explained in the next section. In the analysis of lidar data, we need to compute and correct for the contributions of clear air backscattering and extinction (Rayleigh scattering) to the measured total (particle plus Rayleigh) backscatter and extinction coefficients. We downloaded GDAS (Global Data Assimilation System) height profiles of temperature and pressure of the National Weather Service's National Centers for Environmental Prediction (NCEP) for our computations of Rayleigh scattering contributions (NOAA's Air Resources Laboratory ARL, <https://www.ready.noaa.gov/gdas1.php>). The temperature profiles are also used in the INP parameterizations in Sect. 3.4.

## 3 Methodology

In this section, the equations for the conversion of the measured optical aerosol properties into the microphysical properties are presented. Figure 1 illustrates the general idea of our approach. Table 2 provides an overview of the different steps of the entire data analysis. All steps 1–6 are explained in detail in the following Sects. 3.1–3.4.

Section 3.1 starts with a brief explanation how we derive and estimate the required height profiles of particle extinction coefficient  $\sigma_i$  for the three aerosol components (index  $i = \text{d, c, and m}$ ), i.e., for desert dust (d), non-desert continental aerosol contributions (c), and marine particles (m). In Sect. 3.2, we present the conversion method applied to obtain the height profiles of the required particle number and surface area concentrations for dry particles of each defined aerosol type (d, c, m) from the lidar-derived profiles of  $\sigma_{\text{d}}$ ,  $\sigma_{\text{c}}$ , and  $\sigma_{\text{m}}$ . In Sect. 3.3, we provide a simple parameterization scheme which uses the particle number concentra-

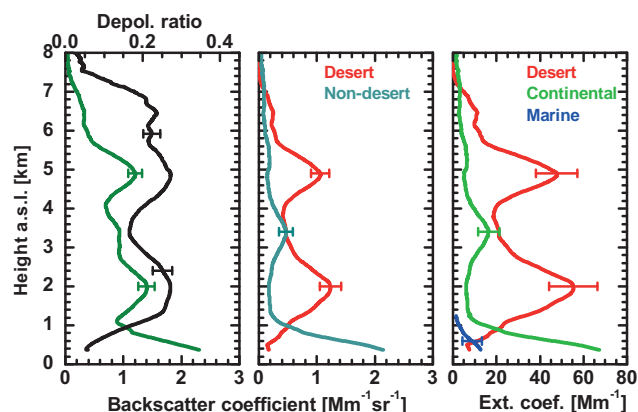


**Figure 1.** Overview of the entire data analysis scheme. The polarization lidar enables us to separate desert and non-desert backscatter coefficients  $\beta_d$  and  $\beta_{nd}$  by means of the particle linear depolarization ratio. In the next step, the backscatter coefficients for marine particles  $\beta_m$  and non-desert continental aerosol mixtures  $\beta_c$  are separated by means of, e.g., backward (BW) trajectory analysis, Ångström exponent information, and by using marine backscatter estimates (see text). The three backscatter coefficients are then converted to aerosol-type-dependent particle extinction coefficients  $\sigma_i$ , which in turn are converted to profiles of particle number concentrations  $n_{100,d,dry}$ ,  $n_{50,m,dry}$ ,  $n_{50,c,dry}$ , and  $n_{250,i,dry}$ , and particle surface area concentration  $s_{i,dry}$ . Finally, CCN-relevant particle number concentrations  $n_{CCN,ss,i}$  are estimated for a given supersaturation (ss) as well as ice-nucleating particle number concentrations  $n_{INP,i}$  by applying INP parameterizations from the literature indicated by D10, D15, D16, N12, and S15 for DeMott et al. (2010), DeMott et al. (2015a), DeMott et al. (2015b), Niemand et al. (2012) and Steinke et al. (2015), respectively.

tions  $n_{50,i,dry}$  for  $i = m$  and  $c$  and  $n_{100,d,dry}$  for desert dust to estimate the CCN-relevant particle number concentration  $n_{CCN,ss,i}$ . In Sect. 3.4, we present the available INP parameterization schemes (DeMott et al., 2010, 2015a; Niemand et al., 2012; Steinke et al., 2015) in which  $n_{250,dry}$  and  $s_{dry}$  profiles are input data. Mamouri and Ansmann (2015) already outlined the principle way to obtain dust-related  $n_{INP}$  from  $n_{250,d,dry}$  profiles by applying the parameterization of DeMott et al. (2015a).

### 3.1 Aerosol-type-dependent $\sigma$ profiles from lidar

Steps 1–3 in Table 2 lead to the required height profiles of the particle extinction coefficients  $\sigma_d$ ,  $\sigma_c$ , and  $\sigma_m$ . The different retrieval steps are explained in Fig. 2. A lidar obser-



**Figure 2.** (Left) 532 nm particle backscatter coefficient (green) and particle linear depolarization ratio (black) as function of height above sea level (a.s.l.), (center) derived particle backscatter coefficients separately for non-desert (marine and continental aerosol, blue-green) and desert particles (red), and (right) particle extinction coefficients separately for marine (blue), continental (green) and desert particles (red). Error bars indicate typical uncertainties. The observation was taken at Nicosia, Cyprus, during a desert dust outbreak from the Sahara on 7 April 2015. Mean profiles for the time period from 20:30–21:30 UTC are shown. The sum of the three extinction profiles match the Raman-lidar-derived total particle extinction profile. Lidar ratios used in the backscatter-to-extinction conversion are 45 sr for desert dust, 35 sr for non-desert continental particles, 20 sr for marine particles in the boundary layer.

vation of a strong Saharan dust outbreak crossing Nicosia during the BACCHUS campaign is presented. This case will be further discussed in Sect. 5.1. As can be seen, the Saharan dust plumes contain (non-desert) aerosol in addition, probably originating from industrial activities in northern Africa (Rodríguez et al., 2011).

In the first step, we determine the height profiles of particle backscatter coefficient  $\beta_p$  and particle linear depolarization ratio  $\delta_p$ , here for the transmitted laser wavelength of 532 nm (Fig. 2, left panel). These profiles of  $\beta_p$  and  $\delta_p$  allow us to separate the desert dust backscatter coefficient  $\beta_d$  and the non-desert backscatter contribution  $\beta_{nd}$  (Fig. 2, center panel). This part of the data analysis is explained in detail by Tesche et al. (2009), Groß et al. (2011), Mamouri et al. (2013), Mamouri and Ansmann (2014), and Nisantzi et al. (2015), and will therefore not be outlined here.

To keep the following steps of the complex data analysis as simple as possible, we concentrate on the aerosol conditions over the polluted European continent and the Eastern Mediterranean. We assume that the optical properties over continental sites are related to desert dust and non-desert continental aerosol (urban haze, smoke, soil and road dust, biological particles), only. The impact of marine particles on the overall aerosol optical properties is ignored. Only over the Mediterranean Sea, the North Atlantic, over islands, and in coastal regions we assume that marine particles significantly

**Table 2.** Overview of the data analysis from the basic lidar-derived aerosol optical properties (particle backscatter and extinction coefficients, linear depolarization ratio) to the height profiles of CCN-relevant particle and INP number concentrations. Indices p, m, c, d, and nd stand for particle, and marine, non-desert continental, desert, and non-desert particles, respectively. ss indicates the supersaturation level.

Step	Computed parameters	Equation terms
1	Profiles of particle backscatter coefficient and particle linear depolarization ratio	$\beta_p, \delta_p$
2	Separation of desert dust and non-desert-dust backscatter coefficients	$\beta_d, \beta_{nd}$
3	Conversion to desert dust, marine and continental particle extinction profiles	$\sigma_d, \sigma_m, \sigma_c$
4	Conversion to particle number and surface area concentrations (aerosol type $i = d, m, \text{ and } c$ )	$n_{100,d,dry}, n_{50,m,dry}, n_{50,c,dry}$ $n_{250,i,dry}, s_{i,dry}$
5	Estimation of $n_{CCN,ss}$ from $n_{50,c,dry}, n_{50,m,dry}, \text{ and } n_{100,d,dry}$	$n_{CCN,ss,i}$
6	Estimation of $n_{INP}$ from $n_{250,dry}$ and $s_{dry}$ for each aerosol type $i$	$n_{INP,i}$

contribute to the observed optical properties. To keep again the CCN and INP retrievals simple, we ignore a potential marine contribution to aerosol extinction in the free troposphere. This is justified as our numerous lidar observations in remote oceanic areas indicate, as will be discussed in more detail in Sect. 3.4.1. Backward trajectories, AE values from photometer observations, and the usually available retrievals of the particle extinction-to-backscatter ratio (Nisantzi et al., 2015) will support us to estimate the contribution of marine particles in the planetary boundary layer (PBL) to the determined non-desert backscatter coefficient  $\beta_{nd}$ . In Fig. 2, we assume a small marine contribution of the order of 20 % to the non-desert backscatter coefficient.

After the separation of the backscatter contribution, we multiply the three backscatter profiles of  $\beta_d, \beta_c$ , and  $\beta_m$  with appropriate lidar ratios of 35–40 sr for Middle East dust, 45–55 sr for Saharan dust, 35–75 sr for continental non-desert aerosol, and 15–20 sr for marine particles to obtain the  $\sigma_i$  profiles for the three aerosol components (see Fig. 1, step from  $\beta$  to  $\sigma$ , and Fig. 2, right panel). The overall uncertainty in the  $\sigma$  retrieval is estimated to be of the order of 15–25 % for desert dust extinction coefficients and 20–40 % for the non-desert continental extinction contribution (Tesche et al., 2009; Mamouri et al., 2013; Mamouri and Ansmann, 2014). A significant source of uncertainty is the lidar ratio for continental aerosol which can vary between about 35 sr for almost non-absorbing anthropogenic haze and 75 sr for strongly absorbing biomass burning smoke (Müller et al., 2007; Groß et al., 2013). If a combined Raman/polarization lidar is used, the Raman-lidar derived  $\sigma_p$  must be in agreement with the sum of the three  $\sigma_i$  profiles (for desert, marine, and non-desert continental aerosol particles) in Fig. 2 (right panel). Strong deviations then usually indicate a wrong estimate of the lidar ratio for continental aerosol pollution, as our experience shows.

### 3.2 Profiles of $n_{50,dry}, n_{100,dry}, n_{250,dry}, \text{ and } s_{dry}$ from lidar-derived $\sigma$ profiles

In the next step (step 4 in Table 2, and the step from  $\sigma$  to  $n$  and  $s$  in Fig. 1), we derive profiles of the required parti-

cle number and surface area concentrations for dry particles from the particle extinction coefficients  $\sigma_d, \sigma_c$ , and  $\sigma_m$ . Because the observed ambient particle extinction coefficients are related to microphysical properties such as  $n_{50}, n_{100}, n_{250}$ , and  $s$  at ambient relative humidity conditions, we need to consider water-uptake by hygroscopic particles.

The respective conversion parameters, required to estimate  $n_{dry}$  and  $s_{dry}$  from the ambient  $\sigma$  values, are obtained from the AERONET correlation study presented in Sect. 4. Regarding water uptake by desert dust, we assume in the correlation studies that desert particles are hydrophobic so that a correction is not necessary. Therefore, we directly used the measured column values of  $n_{100,d}, n_{250,d}$ , and  $s_d$  in the AERONET correlation study as proxies for  $n_{100,d,dry}, n_{250,d,dry}$ , and  $s_{d,dry}$ , respectively. As already mentioned in the introduction (Sect. 1) and explained in more detail in Sect. 3.3,  $n_{100,d,dry}$  is the appropriate number concentration in the CCN parameterization for desert dust.

For hygroscopic continental aerosol particles, we assume a typical relative humidity of 60 % ( $\pm 20$  %) for boundary layer aerosols (main reservoir of continental aerosol) as well as for lofted aerosol plumes in the free troposphere. According to 20 years (1995–2015) of radiosonde ascents in Germany (Essen, Munich, and Lindenberg) the mean relative humidity in the boundary layer is 70 % (March to October, most AERONET observations are performed during these months) and 75 % (January to December, M. Pattantyús-Ábrahám, Meteorological Observatory Hohenpeissenberg, personal communication, 2016). Keeping in mind that lidar observations (and AERONET observations in Sect. 4) are predominantly performed at comparably dry conditions, the assumption of an average relative humidity of around 60 % is justified. We assume similar relative humidity conditions in the aerosol layers over Cyprus during times with dominating continental aerosol pollution. The particle radius of continental aerosol particles at 60 % relative humidity is about a factor of 1.15 ( $\pm 0.05$ ) larger than the respective dry particle radius (Skupin et al., 2016). Therefore we use  $n_{60,c}, n_{290,c}$ , and  $s_c/1.33$  in the following as proxies for  $n_{50,c,dry}, n_{250,c,dry}$ , and  $s_{c,dry}$ , respectively. As explained in Sect. 2.1,

$n_{60,c}$  and  $n_{290,c}$  consider all particles with radius  $> 57$  nm and  $>$  about 290 nm, respectively.

For marine particles we assume a relative humidity of 80 % in the water-uptake correction (in the AERONET correlation study). A relative humidity of around 80 % is typical for marine boundary layers. At these moist conditions, marine particles are about a factor of 1.6–2 larger than dry marine particles (O'Dowd and de Leeuw, 2007; Zieger et al., 2010, 2013; Zhang et al., 2014). For our study, we use  $n_{100,m}$ ,  $n_{500,m}$ , and  $s_m/4$  in the following as proxies for  $n_{50,m,dry}$ ,  $n_{250,m,dry}$ , and  $s_{m,dry}$ , assuming that at sea-salt-controlled conditions (sea salt is the most important aerosol type with respect to CCN and INP studies) the particle growth can be as large as a factor of 2 in radius increase. The comparison of the results obtained with our CCN retrieval for marine particles with in situ observed marine CCNC and particle extinction coefficients (Shinozuka et al., 2015) in Sect. 4 will demonstrate that our selection of  $n_{100,m}$  as a basis for the estimation of marine  $n_{CCN}$  is appropriate.

In accordance with Shinozuka et al. (2015), we now can make use of the following approach to estimate  $n_{50,c,dry}$ ,  $n_{50,m,dry}$ , and  $n_{100,d,dry}$  from ambient  $\sigma_i$  for the aerosol types  $i = d, c$ , and  $m$ :

$$n_{100,d,dry}(z) = c_{100,d} \times \sigma_d^{x_d}(z), \quad (1)$$

$$n_{50,c,dry}(z) = c_{60,c} \times \sigma_c^{x_c}(z), \quad (2)$$

$$n_{50,m,dry}(z) = c_{100,m} \times \sigma_m^{x_m}(z), \quad (3)$$

with  $n_{100,d,dry}$ ,  $n_{50,c,dry}$ , and  $n_{50,m,dry}$  in  $\text{cm}^{-3}$ , the conversion factor  $c_{100,d}$ ,  $c_{60,c}$ , and  $c_{100,m}$  in  $\text{cm}^{-3}$  for the ambient particle extinction coefficient  $\sigma_i = 1 \text{ Mm}^{-1}$ , the ambient particle extinction coefficient  $\sigma_i$  in  $\text{Mm}^{-1}$ , and the aerosol extinction exponent  $x_i$ . Equations (1)–(3) assume a linear correlation of  $\log n_{100,d}$  with  $\log \sigma_d$ ,  $\log n_{60,c}$  with  $\log \sigma_c$ , and  $\log n_{100,m}$  with  $\log \sigma_m$ . Values for  $c_{100,d}$ ,  $c_{60,c}$ ,  $c_{100,m}$ , and  $x_i$  are given in Table 3 for all three laser wavelengths. Determination of the specific parameters  $c_{100,d}$ ,  $c_{60,c}$ , and  $c_{100,m}$  and  $x_i$  is explained in Sect. 4 (AERONET correlation study).

$n_{250,i,dry}$  for aerosol type  $i$  is related to the corresponding particle extinction coefficient  $\sigma_i$  as follows:

$$n_{250,d,dry}(z) = c_{250,d} \times \sigma_d(z), \quad (4)$$

$$n_{250,c,dry}(z) = c_{290,c} \times \sigma_c(z), \quad (5)$$

$$n_{250,m,dry}(z) = c_{500,m} \times \sigma_m(z), \quad (6)$$

with  $n_{250,i,dry}$  in  $\text{cm}^{-3}$ , the conversion factors  $c_{250,d}$ ,  $c_{290,c}$ , and  $c_{500,m}$  in  $\text{cm}^{-3} \text{ Mm}$ , and the particle extinction coefficient  $\sigma_i$  in  $\text{Mm}^{-1}$ . Equations (4)–(6) assume a linear relationship between the large particle fraction  $n_{250}$  and  $\sigma_d$ ,  $n_{290}$  and  $\sigma_c$ , and  $n_{500}$  and  $\sigma_m$ . Again, the conversion factors  $c_{250,d}$ ,  $c_{290,c}$ , and  $c_{500,m}$  are listed in Table 3. They are obtained from the correlation analysis in Sect. 4.

Finally, we obtain the particle surface area concentration  $s_{dry}$  for aerosol type  $i$  from

$$s_{d,dry}(z) = c_{s,d} \times \sigma_d(z), \quad (7)$$

$$s_{c,dry}(z) = c_{s,c}/1.33 \times \sigma_c(z), \quad (8)$$

$$s_{m,dry}(z) = c_{s,m}/4 \times \sigma_m(z), \quad (9)$$

with  $s_{i,dry}$  in  $\text{m}^2 \text{ cm}^{-3}$  and the conversion factor  $c_{s,i}$  in  $\text{m}^2 \text{ cm}^{-3} \text{ Mm}$ . Again, a linear relationship between particle surface area  $s_i$  and particle extinction coefficient  $\sigma_i$  is assumed. The  $c_{s,i}$  values are listed in Table 3. The overall uncertainties in all retrievals will be discussed in Sect. 4.4. Standard deviations of all conversion parameters in Table 3 are the basic information in the uncertainty analysis.

### 3.3 Profiles of $n_{CCN,ss}$ from $n_{50,dry}$ and $n_{100,dry}$ profiles

In the next step (step 5 in Table 2 and in Fig. 1, the step from  $n_{50,dry}$  and  $n_{100,dry}$  to  $n_{CCN}$ ), we estimate the profiles of CCN-relevant particle number concentrations. The CCN parameterization is a crucial task. Therefore only the basic approach is presented here. The ability of aerosol particles to serve as CCN is a function of their size, chemical composition, and the level of supersaturation in the ambient cloud layer. The supersaturation (ss) depends on the updraft velocities and typically is in the range of  $ss = 0.1$  to 1 %. The higher the supersaturation, the smaller the particles that can be activated, and thus the higher the number concentration of potential CCN. We will restrict our CCN parameterization here to the low supersaturation of 0.1–0.2 %. The CCN number concentration can be easily a factor of 2–3 higher when the updraft speeds causes conditions with supersaturation of 0.4–1 %.

In the case of desert dust, the situation is even more complex (Kumar et al., 2009, 2011; Koehler et al., 2009). During emission, desert dust particles may contain negligible amounts of soluble material. They are typically hydrophobic. During long-range transport, dust particles undergo atmospheric processing and soluble species may form on the particle surfaces. In this way, the ability of desert dust particles to serve as CCN may be significantly improved. A factor 2 or even more particles may be activated. Observations by Shinozuka et al. (2015) and our own SALTRACE observations (CCN number concentrations from airborne in situ observations and particle extinction coefficients from ground-based lidar) however suggest that the soluble fraction must be small, at least for Saharan dust after the long-range transport over 5000–8000 km, so that  $n_{100,d,dry}$  is a good proxy for the particle number concentration of the desert-dust-related CCN reservoir here. This aspect will be further discussed in Sect. 4.



**Table 3.** Conversion parameters required in the conversion of particle extinction coefficients into particle number and surface area concentrations with Eqs. (1)–(9) in Sect. 3.2. The values are derived from the extended AERONET data analysis (Sect. 4) and are given for the laser wavelengths of 355 (380 nm), 532, and 1064 nm.  $c_{50,d}$ ,  $c_{60,c}$ ,  $c_{100,m}$  (in  $\text{cm}^{-3}$  for  $\sigma_i = 1 \text{ Mm}^{-1}$ ), and  $x_i$  and respective standard deviations (SD) are obtained from the log-log regression analysis presented in Sect. 4. The maximum (positive) SD is given in the table. The mean values and SD of  $c_{250,d}$ ,  $c_{290,c}$ ,  $c_{500,m}$  (in  $\text{Mm cm}^{-3}$ ) and  $c_{s,i}$  (in  $10^{-12} \text{ Mm m}^2 \text{ cm}^{-3}$ ) are computed from averaging of all individual observations of these conversion factors of a given data set (listed in Table 1). In the case of the Limassol (Cyprus) and Leipzig (Germany) data, all observations with AE (440–870 nm) > 1.6 are interpreted as continental-aerosol-dominated cases, and the observations with AE (440–870 nm) < 0.5 are assumed to be desert-dust-dominated. During SALTRACE-3 the 340 nm channel of the AERONET photometer was not working properly, so that we provide the respective values for 380 nm (instead of 355 nm).

Desert dust	$c_{100,d}$	$x_d$	$c_{250,d}$	$c_{s,d}$
Cabo Verde, Barbados, 380 nm	$5.8 \pm 1.7$	$0.72 \pm 0.05$	$0.19 \pm 0.02$	$1.90 \pm 0.25$
Cabo Verde, Barbados, 532 nm	$6.5 \pm 1.8$	$0.70 \pm 0.05$	$0.20 \pm 0.02$	$1.94 \pm 0.26$
Cabo Verde, Barbados, 1064 nm	$7.5 \pm 2.1$	$0.69 \pm 0.05$	$0.22 \pm 0.03$	$2.21 \pm 0.29$
Cyprus, dust, 355 nm	$8.5 \pm 2.0$	$0.80 \pm 0.04$	$0.16 \pm 0.03$	$2.60 \pm 0.55$
Cyprus, dust, 532 nm	$11.8 \pm 2.7$	$0.76 \pm 0.04$	$0.18 \pm 0.03$	$2.90 \pm 0.61$
Cyprus, dust, 1064 nm	$20.2 \pm 4.9$	$0.69 \pm 0.04$	$0.23 \pm 0.05$	$3.65 \pm 0.85$
Germany, dust, 355 nm	$9.1 \pm 5.7$	$0.79 \pm 0.09$	$0.17 \pm 0.03$	$2.32 \pm 0.52$
Germany, dust, 532 nm	$13.9 \pm 8.6$	$0.73 \pm 0.09$	$0.20 \pm 0.03$	$2.66 \pm 0.68$
Germany, dust, 1064 nm	$20.3 \pm 14.0$	$0.68 \pm 0.10$	$0.23 \pm 0.03$	$3.14 \pm 1.02$
Continental aerosol	$c_{60,c}$	$x_c$	$c_{290,c}$	$c_{s,c}/1.33$
Cyprus, 355 nm	$105 \pm 28$	$0.67 \pm 0.04$	$0.05 \pm 0.02$	$2.19 \pm 0.73$
Cyprus, 532 nm	$102 \pm 26$	$0.75 \pm 0.05$	$0.09 \pm 0.02$	$3.87 \pm 1.23$
Cyprus, 1064 nm	$460 \pm 79$	$0.59 \pm 0.04$	$0.31 \pm 0.10$	$13.51 \pm 5.17$
Germany, 355 nm	$12.1 \pm 1.7$	$0.97 \pm 0.02$	$0.06 \pm 0.03$	$1.55 \pm 0.46$
Germany, 532 nm	$25.3 \pm 3.3$	$0.94 \pm 0.03$	$0.10 \pm 0.04$	$2.80 \pm 0.89$
Germany, 1064 nm	$108 \pm 14$	$0.85 \pm 0.03$	$0.33 \pm 0.16$	$8.98 \pm 3.69$
Marine aerosol	$c_{100,m}$	$x_m$	$c_{500,m}$	$c_{s,m}/4$
Barbados, 355 nm	$2.7 \pm 1.6$	$1.06 \pm 0.11$	$0.05 \pm 0.01$	$0.52 \pm 0.09$
Barbados, 532 nm	$7.2 \pm 3.7$	$0.85 \pm 0.11$	$0.06 \pm 0.01$	$0.63 \pm 0.11$
Barbados, 1064 nm	$35.4 \pm 12.3$	$0.50 \pm 0.08$	$0.09 \pm 0.02$	$0.95 \pm 0.22$

We estimate  $n_{\text{CCN},ss,i}$  now in the following way:

$$n_{\text{CCN},ss,d}(z) = f_{ss,d} \times n_{100,d,\text{dry}}(z), \quad (10)$$

$$n_{\text{CCN},ss,c}(z) = f_{ss,c} \times n_{50,c,\text{dry}}(z), \quad (11)$$

$$n_{\text{CCN},ss,m}(z) = f_{ss,m} \times n_{50,m,\text{dry}}(z), \quad (12)$$

with  $f_{ss,i} = 1.0$  for  $ss = 0.15\%$ . According to the literature, non-desert aerosol particles with dry particle radius of about > 40 nm (at  $ss = 0.25\%$ ) and > 30 nm (at  $ss = 0.4\%$ ) form the reservoir of potential CCN (Quinn et al., 2008; Rose et al., 2010; Deng et al., 2011; Ditas et al., 2012; Siebert et al., 2013; Henning et al., 2014). This was found from a variety of studies conducted in very different regions of the world and for very different aerosol mixtures. Only for supersaturation values of about 0.2% and lower does  $n_{50,\text{dry}}$  seem to represent the particle number concentration of the CCN reservoir. By inspection of the size distributions for pure marine aerosols (Bates et al., 2000), continental pollution aerosol (Beddows et al., 2014) and our own Leipzig city size distributions (measured at TROPOS throughout the year 2008), we found  $n_{30,\text{dry}}/n_{50,\text{dry}}$  ratios on the order of 1.7 ( $\pm 0.8$ ) and  $n_{40,\text{dry}}/n_{50,\text{dry}}$  of about 1.35 ( $\pm 0.7$ ). These values

may be used as the enhancement factor  $f_{ss,i}$  in Eqs. (10)–(12), i.e.,  $f_{ss=0.25\%,i} = 1.35$  and  $f_{ss=0.4\%,i} = 1.70$ . Ji and Shaw (1998) found for pure ammonium sulfate in laboratory studies enhancement factors of 1.26 ( $ss = 0.25\%$ ) and 1.46 ( $ss = 0.4\%$ ). Shinozuka et al. (2015) assumes an increase of  $n_{\text{CCN}}$  by a factor of 2 when the supersaturation increases from 0.2 to 0.4%. Hiranuma et al. (2011), however, also mentioned that natural aerosols show a much more complex behavior regarding these enhancement factors than discussed here.

In the case of desert dust, cloud droplet activation may include particles with dry radius as low as 50 nm at supersaturation of 0.15–0.2%, when the particles are coated with soluble material. According to the AERONET size distributions, the number concentration  $n_{50,\text{dry}}$  is roughly a factor of 4 higher than  $n_{100,\text{dry}}$ . All these uncertainties lead to the conclusion of Shinozuka et al. (2015) that the uncertainty range for  $n_{\text{CCN},ss,i}$  is best described by a factor of 3 around the derived solutions. By using  $n_{60,c}$ ,  $n_{100,m}$ , and  $n_{100,d}$  as proxies for  $n_{50,c,\text{dry}}$ ,  $n_{50,m,\text{dry}}$ , and  $n_{100,d,\text{dry}}$  in Eqs. (10)–(12), the  $n_{\text{CCN},ss=0.15\%}$  values presented in Sects. 4 and 5 may be

therefore interpreted as the minimum values of the possible solution space for  $n_{\text{CCN,ss}}$ .

### 3.4 Profiles of $n_{\text{INP}}$ from $n_{250,\text{dry}}$ and $s_{\text{dry}}$ profiles

The final step of the retrieval (step 6 in Table 2, and in Fig. 1, the step from  $n_{250,i,\text{dry}}$  and  $s_{\text{dry}}$  to  $n_{\text{INP},i}$ ) leads to the estimation of the INP number concentration profiles. Different parameterizations can be used based on  $n_{250,\text{dry}}$  (DeMott et al., 2010, 2015a) or  $s_{\text{dry}}$  profiles (Niemand et al., 2012; Steinke et al., 2015).

#### 3.4.1 Estimation of $n_{\text{INP}}$ from $n_{250,\text{dry}}$

The INP parameterizations introduced by DeMott et al. (2010, 2015a) hold for  $n_{250,\text{dry}}(p_0, T_0)$  and thus standard (std) pressure ( $p_0 = 1013$  hPa) and temperature ( $T_0 = 273.16$  K) conditions (see Eqs. 13 and 14). Therefore, we have to convert each profile value  $n_{250,\text{dry}}(p_z, T_z)$  from ambient pressure  $p_z$  and temperature  $T_z$  at height  $z$  to  $n_{250,\text{dry}}(p_0, T_0)$  by using the factor  $(T_z p_0)/(T_0 p_z)$ .

DeMott et al. (2010) introduced a so-called global INP parameterization which is based on nine field campaigns conducted in Colorado (4 campaigns), eastern Canada (2 campaigns), Amazonia, Alaska, and in the Pacific Basin. This INP characterization scheme is, in our opinion, suitable for an INP parameterization of non-desert continental aerosols (for mixtures of anthropogenic haze, biomass burning smoke, biological particles, soil and road dust):

$$n_{\text{INP,c}}(p_0, T_0, T_z) = a_1(273.16 - T_z)^{b_1} \times n_{250,c,\text{dry}}(p_0, T_0)^{[c_1(273.16 - T_z) + d_1]}, \quad (13)$$

with  $n_{250,c,\text{dry}}$  in  $\text{std cm}^{-3}$ ,  $n_{\text{INP,c}}$  in  $\text{std L}^{-1}$ ,  $a_1 = 0.0000594$ ,  $b_1 = 3.33$ ,  $c_1 = 0.0265$ ,  $d_1 = 0.0033$ , and temperature  $T(z)$  in K (and  $< 273.16$  K). Note that the values of  $a_1$ ,  $b_1$ ,  $c_1$  and  $d_1$  given in Mamouri and Ansmann (2015) are erroneous. However, all computations presented in that article were performed with the correct values listed here.

Finally, we transfer the obtained values of  $n_{\text{INP,c}}(p_0, T_0, T_z)$  to the ones for ambient pressure and temperature conditions,  $n_{\text{INP,c}}(p_z, T_z)$ , by multiplying  $n_{\text{INP,c}}(p_0, T_0, T_z)$  with the factor  $(T_0 p_z)/(T_z p_0)$ .

It should be emphasized that this INP parameterization shows an uncertainty in the range of a factor of 5–10 as recent observation indicate (McCluskey et al., 2014; Mason et al., 2015; Taylor et al., 2016a, b). The most obvious reason for the remaining uncertainty is that the specific aerosol composition, i.e., the mixture of aerosol types (the exact amount of pollen, dust, soot, organic material, sulfates, etc.), during an actual measurement case is not known. Strong differences in the INP efficacy of different aerosol types is found in laboratory studies (see review of Murray et al., 2012). Furthermore, observations also indicate that particles with radii  $< 250$  nm may be activated as well (Mason et al., 2016). The size effect

was found to increase with decreasing temperature. Nevertheless, we use this schemes for continental aerosol mixtures (by excluding explicitly desert dust) because it explains many of the details of the found relationship between the observed fractions of ice-containing clouds and cloud top temperature of altocumulus layers which formed over the European continent in aged aerosol mixtures. This INP parameterization especially predicts significant heterogeneous ice nucleation already at high temperatures of  $-5$  to  $-15$  °C as observed (Seifert et al., 2010; Kanitz et al., 2013).

The INP parameterization scheme for mineral dust of DeMott et al. (2015a) is used here explicitly for desert dust:

$$n_{\text{INP,d}}(p_0, T_0, T_z) = f_{\text{d}} n_{250,d,\text{dry}}(p_0, T_0)^{[a_2(273.16 - T_z) + b_2]} \times \exp[c_2(273.16 - T_z) + d_2], \quad (14)$$

with the so-called atmospheric correction factor  $f_{\text{d}} = 3$ , and the coefficients  $a_2 = 0.0$ ,  $b_2 = 1.25$ ,  $c_2 = 0.46$ , and  $d_2 = -11.6$ . Again, to obtain the  $n_{\text{INP}}$  profile for ambient temperature and pressure conditions, we have to transfer the obtained values of  $n_{\text{INP,d}}(p_0, T_0, T_z)$  to the ones for ambient pressure and temperature conditions in the same way as described above for  $n_{\text{INP,c}}(p_z, T_z)$ , namely by multiplying  $n_{\text{INP,d}}(p_0, T_0, T_z)$  with the factor  $(T_0 p_z)/(T_z p_0)$ .

According to DeMott et al. (2015a), Eqs. (13) and (14) can be used to estimate  $n_{\text{INP}}$  for immersion freezing processes. The formulas are applicable to the temperature range from  $-9$  to  $-35$  °C (Eq. 13) and  $-21$  to  $-35$  °C (Eq. 14). In Sect. 5 (lidar case studies), we use these immersion-freezing-based parameterizations for higher temperatures as well. According to Wex et al. (2014) ice nucleation for anthropogenic particles (with an insoluble part) and coated mineral dust particles (coated with natural and/or anthropogenic soluble material) can be described as immersion freezing even at temperatures  $< -35$  °C. Above the deliquescence relative humidity, additional water is added to the coating and a solution shell forms around the insoluble part of the particles, causing them to nucleate ice from concentrated solutions via the immersion freezing pathway, taking a freezing point depression into account.

Regarding the uncertainties in the INP computation, we assume that Eq. (14) allows a prediction of dust  $n_{\text{INP}}$  within an uncertainty range of a factor of 2–5 (DeMott et al., 2015a; Schrod et al., 2016). An overview of all uncertainties in the basic lidar-derived particle optical properties, the retrieved microphysical aerosol properties, and the finally estimated  $n_{\text{INP},i}$  values is given in Sect. 4.4.

Recently, DeMott et al. (2015b) compared the potency of marine and continental INPs. By comparing laboratory studies and field observations it was found that for typical marine (sea spray) and continental aerosol conditions characterized by ambient particle extinction coefficients of  $50$ – $100$   $\text{Mm}^{-1}$  at  $500$  nm wavelength, the marine INP number concentration was lower by about 3 orders of magnitude than the continental INP number concentration. Compared to terrestrial parti-

cles, sea spray particles are obviously bad INPs (efficacy is a factor of 300–500 lower) which is in agreement with mixed-phase cloud observations in the northern midlatitudes (high amount of terrestrial particles) and in the southern Ocean (Punta Arenas, Chile, very low amount of terrestrial particles) (Seifert et al., 2010; Kanitz et al., 2011). However, the temperature dependence of heterogeneous ice formation caused by marine and terrestrial particles (as given by Eq. 13) was found to be similar (DeMott et al., 2015b). Therefore, in order to roughly estimate marine  $n_{\text{INP,m}}$  we simply use Eq. (13) with  $n_{250,\text{m,dry}}$  (after Eq. 6) as input and divide the resulting  $n_{\text{INP}}$  value by 350 (DeMott et al., 2015b). Alternatively, one may use Eqs. (15)–(16), introduced in the next Sect. 3.4.2 for mineral dust, but here for marine particles with  $s_{\text{m,dry}}$  as input and divide the result by 350 to obtain  $n_{\text{INP,m}}$ .

In the estimation of actual marine INP number concentrations at given environmental conditions (mixture of marine and terrestrial aerosols) one should mention again that the polarization lidar technique allows us to separate dust from non-dust aerosol components, but not a further separation of marine from continental aerosol pollution. We must therefore estimate the impact of marine particles on the non-desert aerosol extinction coefficient. As stated in Sect. 3.1, over the oceans, we can assume that the extinction coefficient in the PBL is widely determined by marine particles. In continental outflow regimes and at coastal sides (because of sea breeze effects) we must estimate the contribution by continental particle scattering and absorption to the overall aerosol extinction coefficient in the PBL. We may use the AE information from AERONET observations or multiwavelength lidar observations to estimate the contributions by marine and continental particles to the observed overall non-desert extinction coefficient.

As stated in Sect. 3.1, we ignore a marine contribution of  $\sigma_{\text{m}}$  to the particle extinction coefficient in the free troposphere, and therefore a marine contribution to the CCN and INP reservoirs ( $n_{\text{CCN}}$ ,  $n_{\text{INP}}$ ) in the free troposphere. This is corroborated by our lidar observation at Punta Arenas, Chile, Cape Town, South Africa, aboard the R/V Polarstern, and many Polly lidar sites around the globe (Kanitz et al., 2013; Seifert et al., 2015; Baars et al., 2012, 2016). We conclude from these lidar observations that the marine extinction coefficient  $\sigma_{\text{m}}$  is  $< 1\text{--}2 \text{ Mm}^{-1}$  for free-tropospheric heights  $< 3\text{--}5 \text{ km}$ , and of the order of  $0.01\text{--}0.2 \text{ Mm}^{-1}$  for heights  $> 5 \text{ km}$ . Only by strong updrafts below cumulus towers with cloud base in the marine PBL can a large amount of marine particles over oceanic sites be injected into the free troposphere, potentially triggering strong heterogeneous ice formation when the air parcels ascend to heights with temperatures below  $-25^\circ\text{C}$ . For typical marine ambient particle extinction coefficients of  $50\text{--}100 \text{ Mm}^{-1}$  in the marine PBL, we obtain an estimate of roughly  $5\text{--}10 \text{ INP per m}^3$  at  $-25^\circ\text{C}$ . At free tropospheric aerosol background conditions with marine particle extinction coefficients of the order of  $0.1\text{--}1 \text{ Mm}^{-1}$ ,  $n_{\text{INP,m}}$  is in the range from  $0.01\text{--}0.1 \text{ m}^{-3}$  at temperatures of

$-25^\circ\text{C}$  according to the study of DeMott et al. (2015b), and thus 4–5 orders of magnitude lower than  $n_{\text{INP,d}}$  in our dust outbreak case study at  $-25^\circ\text{C}$  (at 6.5 km height) presented in Sect. 5.1.

### 3.4.2 Estimation of $n_{\text{INP}}$ from $s_{\text{dry}}$

$n_{\text{INP,d}}$  profiles can also be estimated from the  $s_{\text{d}}$  profiles. An immersion-freezing INP parameterization is provided by Niemand et al. (2012):

$$n_{\text{INP,d}}(T_z) = 1000 \times s_{\text{d,dry}}(z) \times \eta_{\text{im}}(T_z), \quad (15)$$

$$\eta_{\text{im}}(T_z) = \exp[-0.517(T_z - 273.16) + 8.934] \quad (16)$$

with  $n_{\text{INP,d}}$  in  $\text{L}^{-1}$ ,  $s_{\text{d,dry}}$  in  $\text{m}^2 \text{ cm}^{-3}$  (so that a multiplication by 1000 is needed to obtain  $s$  in  $\text{m}^2 \text{ L}^{-1}$ ), and  $\eta_{\text{im}}$  in  $\text{m}^{-2}$ . The  $n_{\text{INP,d}}$  profile holds for temperatures from 237–261 K ( $-12$  to  $-37^\circ\text{C}$ ).

Steinke et al. (2015) provides a deposition-freezing parameterization:

$$n_{\text{INP,d}}(T_z) = 1000 \times s_{\text{d,dry}}(z) \times \eta_{\text{dep}}(T_z), \quad (17)$$

$$\eta_{\text{dep}}(T_z) = 1.88 \times 10^5 \times \exp(0.2659\chi(T_z)), \quad (18)$$

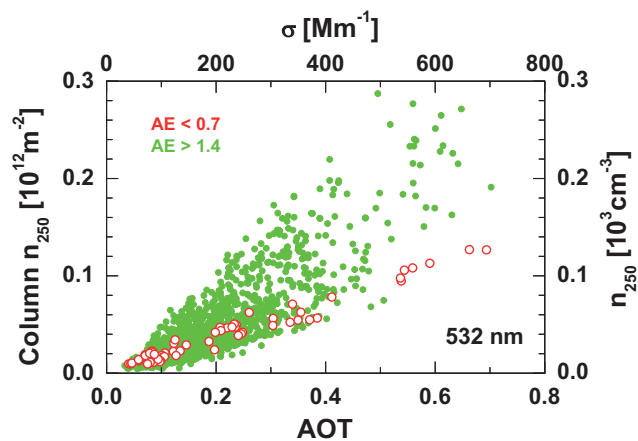
$$\chi(T_z) = -(T_z - 273.16) + (s_{\text{ICE}} - 1) \times 100 \quad (19)$$

with ice supersaturation of  $ss_{\text{ICE}}$ . We assume a constant, but reasonable value of 1.15 for  $ss_{\text{ICE}}$  indicating frequently occurring moderate supersaturation conditions in ice clouds (Comstock et al., 2008). The  $n_{\text{INP,d}}$  profile holds for temperatures from 220–253 K ( $-20$  to  $-53^\circ\text{C}$ ). This deposition freezing parameterization, however, is based on laboratory studies of heterogeneous ice nucleation on artificially produced mineral dust particles (Arizona test dust) which usually show an enhanced freezing efficacy compared to natural desert dust aerosols.

## 4 AERONET correlation study

Of key importance for the entire retrieval of cloud-relevant microphysical aerosol parameters from lidar-derived particle extinction coefficient profiles at ambient conditions are trustworthy conversion parameters  $c_{60,\text{c}}$ ,  $c_{100,\text{d}}$ ,  $c_{100,\text{m}}$ ,  $c_{250,\text{d}}$ ,  $c_{290,\text{c}}$ ,  $c_{500,\text{m}}$ ,  $c_{s,i}$ , and exponents  $x_i$  as required to solve Eqs. (1)–(9). These conversion parameters are derived from the long-term AERONET observations at Leipzig and Limassol (for northern and southern European continental aerosol mixtures), at Ragged Point (for pure marine conditions), and the short-term dust-related field campaigns in Morocco, Cabo Verde, and Barbados (for pure desert dust scenarios, see Table 1). The main results of the AERONET data analysis are presented and discussed in this section.

We performed the AERONET correlation study separately for all three laser wavelengths, but show the results for the mostly used lidar wavelength of 532 nm, only. To facili-



**Figure 3.** Replacement of the column integrals (column  $n_{250}$ , aerosol optical thickness AOT) obtained from the AERONET observations by volume-related quantities ( $n_{250}$ , particle extinction coefficient  $\sigma$ ) by assuming an arbitrarily chosen vertical height of the column of 1000 m. In this example, green symbols show all 1523 pollution-dominated Leipzig AERONET observations (2001–2015, level 2.0) with high Ångström exponent ( $AE > 1.4$ ), and red symbols all 59 dust-dominated cases ( $AE < 0.7$ ).

tate our studies and to be consistent with the work of Shinzuka et al. (2015), who investigated the correlation between CCNC and  $\sigma$  at 500 nm, we replaced all column integrals, i.e., AOTs and the column values of  $n$  and  $s$  by respective volume-related values. For this, we introduced a normalizing, arbitrarily selected vertical column height of 1000 m and divided all basic AERONET observational data points by 1000 m. An example of the transformation is illustrated in Fig. 3 for the Leipzig observation of the column-integrated  $n_{250}$  and AOT at 532 nm. The volume-related values can be interpreted as the vertical averages of  $n$ ,  $s$ , and  $\sigma$  in the assumed 1000 m deep column. It should be mentioned that the selected column height has no impact on the data analysis, but is set to a realistic value so that the range of  $\sigma$ , typically measured with lidar for a given site, is covered.

#### 4.1 Leipzig and Limassol long-term observations of mixed aerosols

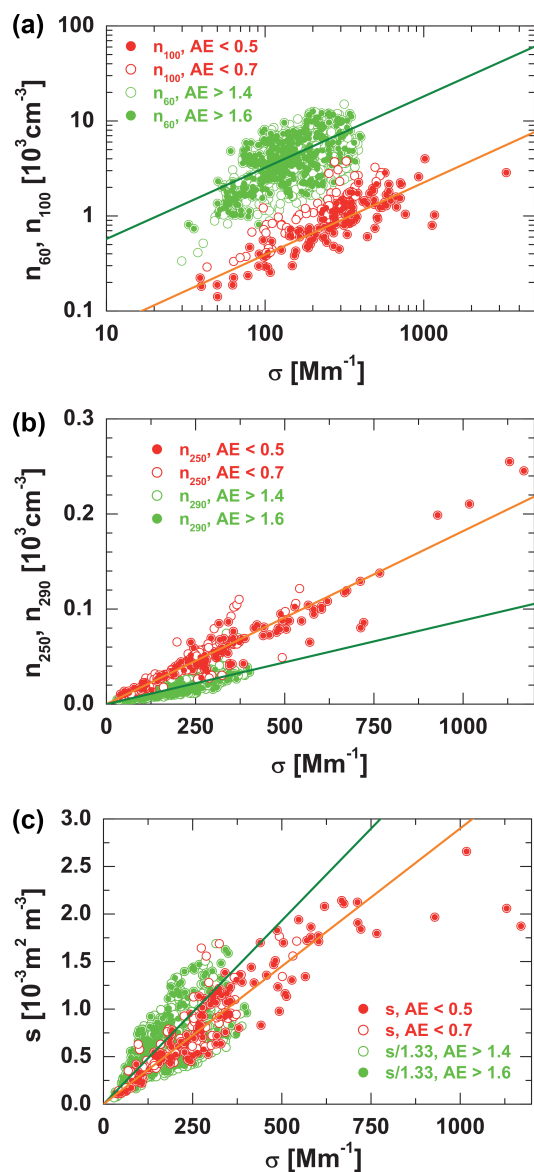
A total number of 48 474 and 34 982 sun/sky photometer observation (level 1.5) were taken at Limassol and Leipzig, respectively, during the 2011–2015 (Limassol) and 2001–2015 (Leipzig) time periods. 4190 and 4651 of these measurements at Limassol and Leipzig could be used to derive particle size distributions and thus column values of  $n$  and  $s$ . 1745 Limassol and 2157 Leipzig quality-assured level-2.0 data sets were finally available for our study. The observed correlations of  $n_{60}$ ,  $n_{290}$ , and  $s/1.33$  vs. 532 nm  $\sigma$  for pollution-dominated scenarios (Ångström exponents  $AE > 1.4$  and  $> 1.6$ ) and of  $n_{100}$ ,  $n_{250}$ , and  $s$  for the desert-dust-dominated cases ( $AE < 0.7$  and  $< 0.5$ ) are shown in Figs. 4 and 5. The

conversion parameters derived from the correlation analysis are given in Table 3 and used in Eqs. (1)–(9).

As shown in Figs. 4 and 5, at both sites a large variability in the aerosol conditions is observed. Limassol in the Eastern Mediterranean experiences complex aerosol conditions almost every day. This Middle East (Eastern Mediterranean) station is influenced by frequent dust outbreaks from the Sahara and the deserts of the Middle East (Nisantzi et al., 2015), biomass burning smoke and fire-induced soil dust injections (Nisantzi et al., 2014) from Turkey, the Black Sea area, and European regions further to the north, and anthropogenic haze from eastern, southeastern and southern Europe, northern Africa, and western Asia. Marine particles form the background aerosol at Limassol at the south coast of Cyprus. In contrast, the central European AERONET station of Leipzig is heavily influenced by fresh and aged anthropogenic pollution, which dominates the boundary layer aerosol (Mattis et al., 2004; Wandinger et al., 2004). A few Saharan dust outbreaks towards central Europe (Ansmann et al., 2003; Papayannis et al., 2008) and long-range transport of biomass burning aerosol and anthropogenic haze from southern Europe and North America determine the aerosol conditions in the free troposphere (Mattis et al., 2008). On average, the free-tropospheric AOT contributes 20 % to the overall AOT (Mattis et al., 2004). The impact of marine aerosol on the Leipzig observations is negligible.

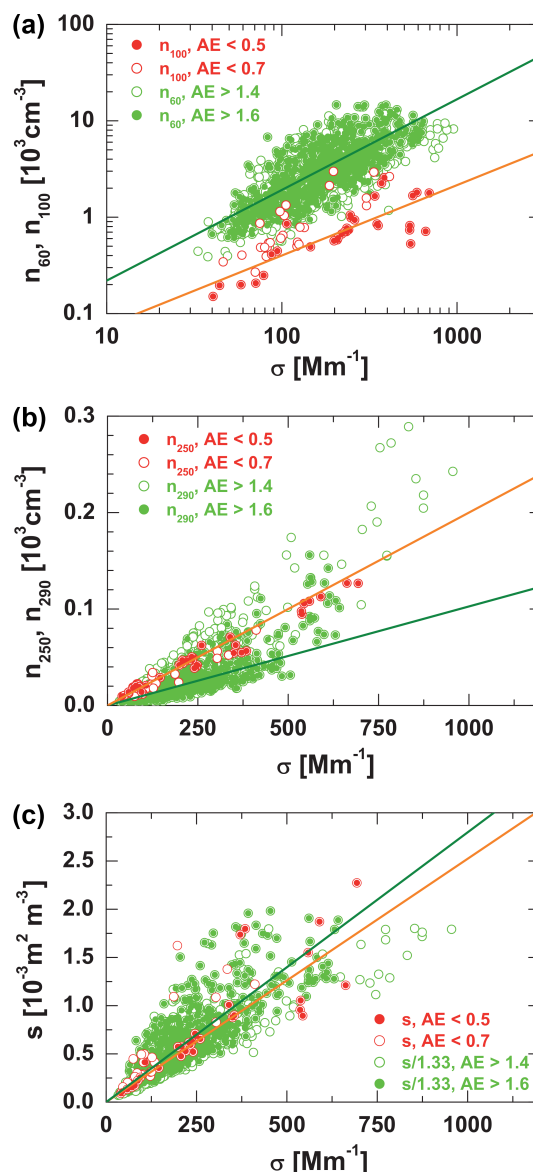
The top panels in Figs. 4 and 5 nicely show that the Limassol and Leipzig AERONET observations are complementary from the statistics point of view. Many more cases with a strong desert dust impact are measured at Limassol (133 cases with  $AE < 0.5$  within 4 years) than at Leipzig (only 33 dust cases in 14 years). The opposite is true for well-mixed anthropogenic haze (with  $AE > 1.6$ ). About 1000 observations are available for Leipzig covering a broad range of particle extinction values from 40–700  $Mm^{-1}$ , whereas at Limassol homogeneous haze/smoke situations are less frequent (421 observations with  $AE > 1.6$ ) and the ambient extinction values cover a range from 30–400  $Mm^{-1}$  only. We used  $AE$  calculated from the AOT values from 440 to 870 nm here to filter out dust-dominated and haze-dominated aerosol observations.

The found scatter in the correlations of  $n_{60}$ ,  $n_{100}$ ,  $n_{250}$ ,  $n_{290}$ , and  $s$  with  $\sigma$  in Figs. 4 and 5 is caused by many reasons. First of all, different particle size distributions (leading to different  $n$  and  $s$  values) can produce the same  $\sigma$  value. The optical efficiency (optical cross section divided by the geometrical cross section  $s/4$ ) of a given log-normal aerosol size distribution can easily vary between 0.3 and 3 as a function of a shifting mode radius of the fine-mode particle spectrum towards larger or smaller sizes without leading to significant changes in the  $n$  and  $s$  values. The particle optical effects depend on ambient relative humidity (significant water up-take by particles occurs when the relative humidity in the vertical column exceeds 75–80 %) so that large changes in  $\sigma$  (within a factor of 1.5–2.5) may be correlated with compar-



**Figure 4.** (a) Particle number concentrations  $n_{100}$  (red),  $n_{60}$  (green), (b)  $n_{250}$  (red),  $n_{290}$  (green), and (c) particle surface area concentration  $s$  (red) and  $s/1.33$  (green) vs. 532 nm particle extinction coefficient  $\sigma$ . AERONET observations performed at Limassol, Cyprus, from 1 July 2011 to 30 June 2015 are shown. 839, 421, 213, and 134 observations are available at Ångström exponents of  $AE > 1.4$  (open green circles),  $> 1.6$  (solid green circles),  $< 0.7$  (open red circles), and  $< 0.5$  (solid red circles), respectively. The olive lines (for  $AE > 1.6$ ) and orange lines (for  $AE < 0.5$ ) indicate the mean increase of  $\log n_{100}$  and  $\log n_{60}$  with  $\log \sigma$  (532 nm, top panel), and the mean increase of  $n_{250}$ ,  $n_{290}$  and  $s$  with 532 nm  $\sigma$ .

bly small changes in  $n_{60}$ ,  $n_{100}$ ,  $n_{250}$ ,  $n_{290}$ , and  $s$ . The aerosol mixtures (or the overall chemical composition of the particles including the type-dependent water uptake and growth effects) may be different for relatively clean aerosol conditions (low  $\sigma$  values) and scenarios with heavy pollution or



**Figure 5.** Same as Fig. 4, except for the AERONET observations at Leipzig from 1 May 2001 and 30 June 2015. 1523, 974, 59, and 33 observations are available at Ångström exponents of  $AE > 1.4$  (open green circles),  $> 1.6$  (solid green circles),  $< 0.7$  (open red circles), and  $< 0.5$  (solid red circles), respectively.

dust outbreaks (high  $\sigma$  values). All this systematically influence the correlation features. The discussed uncertainties in the retrieval of the particle size distribution,  $n$ , and  $s$  (case by case, Sect. 2.1) as quantified by Dubovik et al. (2000) also contribute to the observed variability in the correlations.

As recommended by Shinozuka et al. (2015), we applied the regression analysis to the  $\log n_{60}$ - $\log \sigma$  and  $\log n_{100}$ - $\log \sigma$  data fields (top panels in Figs. 4 and 5). The regression lines in the figures go through the geometric averages of  $n_{100}$  and  $n_{60}$  for the average  $\sigma$  value. The slope of the regression line is the extinction exponent  $x$  in Eqs. (1) and (2). The obtained

numbers for  $x_d$ ,  $x_c$ ,  $c_{100,d}$ , and  $c_{60,c}$  of the log-log regression analysis are given in Table 3 (Cyprus and Germany observations). The standard deviations (root mean square values) of the regressions analysis are mostly 0.15–0.25 in the log scale and thus indicate overall uncertainties within a factor of 1.4–1.8 for  $n_{60,c}$  and  $n_{100,d}$  when estimated from  $\sigma$ . Taking an additional uncertainty in the water-uptake correction into account, we estimate that  $n_{100,d,dry}$  (Eq. 1) and  $n_{50,c,dry}$  (Eq. 2) can be estimated with an uncertainty of a factor of 1.5–2.

We compared our results with respective ones presented by Shinozuka et al. (2015) for likewise rural and background sites (Southern Great Plains, Oklahoma, Cape Cod, Massachusetts, Black Forest, Germany). In these measurements, the dry extinction coefficients for the 500 nm wavelength mainly ranged from 5 to 100  $Mm^{-1}$ . The comparison reveals that the Limassol and Leipzig AERONET data sets clearly represent highly polluted urban conditions. Our observations considered in Figs. 4 and 5 cover an AE range from 1.6–2.2 and thus indicate the strong impact of fine-mode aerosol in these measurements. By using the Leipzig conversion parameters in Table 3 ( $c_{60,c} = 25.3 \text{ cm}^{-3}$  at  $\sigma_c = 1 \text{ Mm}^{-1}$ ,  $x_c = 0.94$ ) we obtain  $n_{CCN} \approx 1000 \text{ cm}^{-3}$  for an ambient extinction value of  $\sigma_c = 50 \text{ Mm}^{-1}$  in Eq. (2) and when inserting the resulting  $n_{50,c,dry}$  in Eq. (11). For Limassol we get even higher CCN-relevant values ( $n_{CCN} \approx 2000 \text{ cm}^{-3}$  for  $\sigma_c = 50 \text{ Mm}^{-1}$ ). Similar values are obtained from horizontal long-path particle extinction measurements at ambient conditions at TROPOS, Leipzig, (Skupin et al., 2014, 2016) and accompanying continuous dry-particle size distribution observations (A. Skupin, personal communication, 2016).

The measurements of Shinozuka et al. (2015) at more rural and background sites indicate  $n_{CCN}$  of 400–500  $\text{cm}^{-3}$  (Southern Great Plains), 350–400  $\text{cm}^{-3}$  (Black Forest), and around 700  $\text{cm}^{-3}$  in the case of Cape Cod at the Atlantic Ocean in the northeastern United States for ambient extinction coefficients of 50  $Mm^{-1}$  at 500 nm and for the AE class from 1.5–1.7 (indicating less fine-mode dominated aerosols compared to the Limassol and Leipzig aerosol conditions). To compare our data (for supersaturations of 0.15 % and ambient instead of dry particle extinction coefficients), we divided the CCNC numbers of Shinozuka et al. (2015), measured at supersaturations around 0.4 %, by a factor of 2, assuming that the resulting numbers then represent CCNC values for  $ss = 0.15 \%$ , and we multiplied the dry extinction coefficients with a factor of 1.4 to obtain ambient extinction coefficients, assuming relative humidities of 60–70 % prevail also in the aerosol layers over Oklahoma, Massachusetts, and the Black Forest in southern Germany.

In this context, it is also noteworthy to mention that Liu and Li (2014) showed that the product of  $\sigma \times AE$  (denoted as Aerosol Index AI, introduced by Nakajima et al., 2001) provides a better correlation with  $n_{CCN}$  than  $n_{CCN}$  with  $\sigma$ . By using AI instead of  $\sigma$  in the correlation, Liu and Li (2014) consider information on the aerosol type and the related size distribution (high AI for fine-mode-dominated aerosol,

low AI for coarse-mode-dominated aerosol conditions). Similarly, Shinozuka et al. (2015) separated the observations in classes with AE from 1.5–1.7 and from 0.3–0.5, and derived AE-dependent parameterizations to obtain estimates of  $n_{CCN}$  from  $\sigma$  observations. In contrast to these approaches, the advantage of our lidar technique is that we separate the different aerosol types by means of the polarization lidar technique first, i.e., before we apply our parameterization and conversion procedures to estimate the microphysical and cloud-relevant aerosol parameters for each aerosol type separately. Nevertheless, we conclude that different parameterizations are required for more rural, background, or free-tropospheric aerosol conditions and for urban aerosols.

A complex regression data analysis as in the top panels of Figs. 4 and 5 is not needed in the study of the  $n_{250}-\sigma$ ,  $n_{290}-\sigma$ , and  $s-\sigma$  relationships. We can assume simple linear relationships because the optical effects of the aerosol mixtures depend approximately linearly on  $s$ ,  $n_{250}$ , and  $n_{290}$ . For all individual, single AERONET observations (belonging to the separate data sets for  $AE > 1.6$  and  $AE < 0.5$ ) we calculated the  $n_{250}/\sigma$ ,  $n_{290}/\sigma$ , and  $s/\sigma$  ratios for all three laser wavelengths. In Figs. 4 and 5, the geometrical averages of these ratios (for 532 nm  $\sigma$ ) define the slopes of the shown straight lines. Shown are both slopes for the fine-mode ( $AE > 1.6$ ) and coarse-mode ( $AE < 0.5$ ) classes. The mean values of  $n_{250}/\sigma$ ,  $n_{290}/\sigma$ , and  $s/\sigma$  for each aerosol subdata set ( $AE < 0.5$ ,  $AE > 1.6$ ) are used as  $c_{250,d}$ ,  $c_{290,c}$ ,  $c_{s,d}$  and  $c_{s,c}$ , respectively, in Eqs. (4), (5), (7), and (8). All Leipzig and Limassol values of  $c_{250,d}$ ,  $c_{290,c}$ ,  $c_{s,d}$  and  $c_{s,c}/1.33$  together with SD (obtained from the averaging procedure) are given in Table 3.

Our results are in good agreement with combined airborne in situ observations of particle number concentration  $n_{150,dry}$  (particles with dry radius  $> 150 \text{ nm}$ ) and lidar-derived particle backscatter coefficients at 532 nm in southern Japan at marine, moderately polluted summer conditions (Sakai et al., 2013). Measurements were performed between 500 m and 5 km height and were influenced by long-range transport of pollution and dust from eastern Asia. By assuming a particle extinction-to-backscatter ratio of 50 sr (typical for a mixture of aged pollution and dust), the conversion factor for the measurements in southern Japan is  $c_{150} \approx 1.0 \text{ Mm cm}^{-3}$  (AE ranged from about 0.3 to 1.0). Our AERONET study indicates for dusty environments that  $n_{150,dry}$  is a factor of roughly 5 higher than  $n_{250,dry}$ . Consequently, our conversion factor  $c_{250} \approx 0.2 \text{ Mm cm}^{-3}$  is 5 times lower than the  $c_{150}$  value.

Our results are also in good agreement with respective model studies of  $s$  for various aerosol types conducted by Barnaba and Gobbi (2001, 2002). According to Barnaba and Gobbi (2001), the  $s-\sigma$  ratio for 500–550 nm for example should be in the range of 2–4 for particle size distributions with strong coarse mode. Kolgotin et al. (2015) found a value of 1.6 ( $\pm 20 \%$ ) for the  $s/\sigma$  ratio at the 355 nm wavelength. In their simulations, they considered mono-modal log-normal

size distributions with mean radius from 20 to 300 nm. We conclude from their study that the  $s/\sigma$  values are in the range from 2.4–3 for 532 nm for haze and dust conditions.

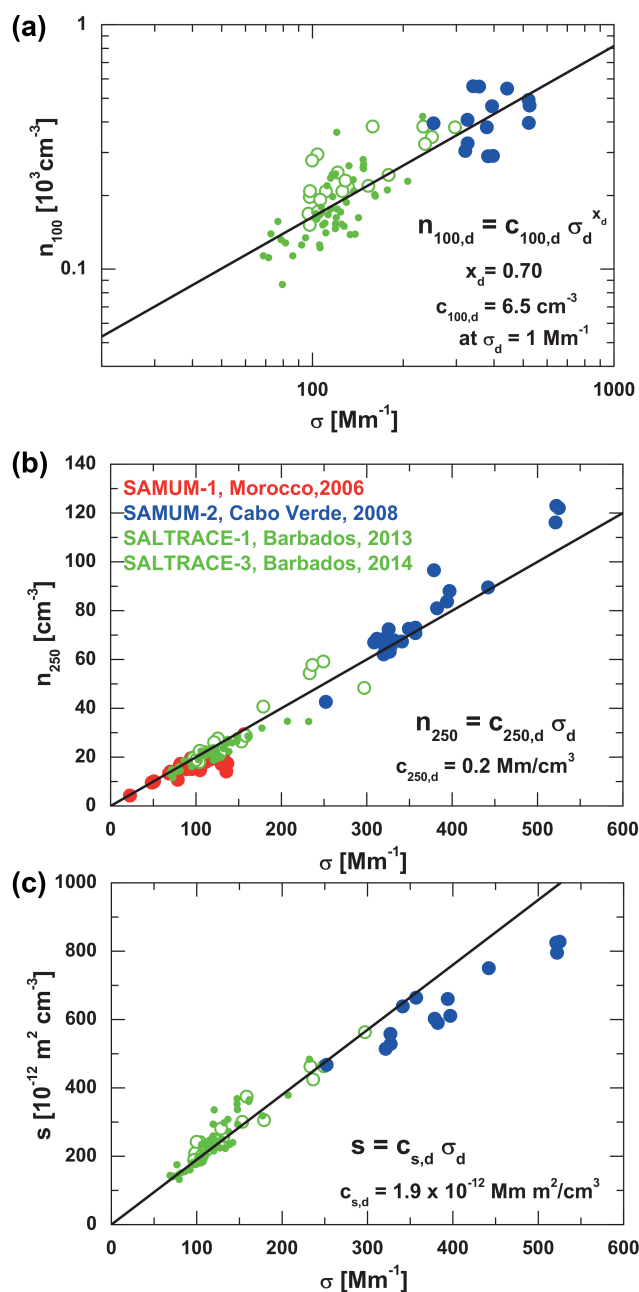
The scatter of the individual observations for the typical range of  $\sigma$  from 50–400  $\text{Mm}^{-1}$  provides insight into the uncertainty in the retrieval of the particle number concentrations and surface area concentrations from the measured particle extinction coefficients. The respective standard deviations of  $c_{250,d}$ ,  $c_{290,c}$ ,  $c_{s,d}$ , and  $c_{s,c}$  in Table 3 are used in the error analysis in Sects. 4.4. The standard deviations roughly indicate that conversions of  $\sigma$  into  $n_{250,d}$ ,  $n_{290,c}$ , and  $s$  are possible with a relative error of 20–30 %.

#### 4.2 Field campaign data sets for pure dust conditions

Unique combined AERONET photometer and multiwavelength lidar observations are available for pure Saharan dust conditions, sampled during several field campaigns in southeastern Morocco (SAMUM-1), close to the dust source at a minimum influence by marine particles and anthropogenic pollution, at Cape Verde (SAMUM-2) during a heavy dust outbreak from 28–30 January 2008, and at Barbados (SALTRACE-1 and 3, lofted dust plumes during the summer months) in the long-range dust transport regime, 5000–8000 km west of the Saharan dust sources (see Table 1).

From all the SAMUM and SALTRACE observations we were able to carefully select 125 cases with dominant dust conditions (indicated by AE values  $< 0.2$ ). For all these data sets, detailed lidar observations of dust layering (layer structures, base and top heights of main dust layers) (Tesche et al., 2009, 2011; Haarig et al., 2015; Groß et al., 2015) are available so that mean dust extinction coefficients and mean values of particle number and surface area concentrations could be calculated for the observed dust layers by combining the AERONET column observations and the layer depth information from the lidar. The results shown in Fig. 6 are based on these dust layer mean values. We also checked all AERONET measurements carefully regarding cloud contamination (subvisible and thin cirrus) by means of the lidar observations. Furthermore, we launched 1–3 radiosondes per day. The relative humidity in the dust layers over Cabo Verde and Barbados was always  $< 50\%$ .

Unfortunately, problems with the AERONET 340 nm channel in Morocco and Barbados (in 2014, SALTRACE-3) prohibit the retrieval of conversion parameters at 355 nm. So, we present the conversion parameters at 380 nm in Table 3 which fairly well represent the parameters for 355 nm in the case mineral dust. Furthermore, the Morocco size distributions are not trustworthy for small particles (sometimes rather high peaks in number concentrations occurred for size bins from 50 to 112 nm radius). The reasons may be related to the missing 340 nm channel and to the fact that the occurrence of very large particles with radius  $> 15\ \mu\text{m}$  at a site close to the desert dust sources can never be excluded (Müller



**Figure 6.** Relationship between dust layer mean 532 nm extinction coefficient  $\sigma$  and (a) particle number concentrations  $n_{100}$ , (b)  $n_{250}$ , and (c) surface area concentration  $s$  for observations taken during the desert dust field campaigns in Morocco (red, SAMUM-1, 2006), Cape Verde (blue, SAMUM-2, 2008), and Barbados (open green circles, SALTRACE-1, 2013, solid green circles, SALTRACE-3, 2014). The slope of the black lines are obtained in the same way as in Figs. 4 and 5. Note again that the  $n_{250}/\sigma$  conversion factor is 0.2 and not  $0.67 \text{ Mm cm}^{-3}$  as erroneously given in Mamouri and Ansmann (2015).

et al., 2010). The AERONET size distribution retrieval, however, considers particles with radius up to 15  $\mu\text{m}$ , only. We therefore did not consider the Morocco AERONET observations in the correlation analysis for  $n_{100,d}$  and  $s_d$  with  $\sigma_d$  in Fig. 6. The desert dust conversion factors in Table 3 are exclusively derived from the Cabo Verde and Barbados observational data.

The results of the correlation analysis in Fig. 6 (for 532 nm) are given in the same way as for the multi-year Leipzig and Limassol data in Figs. 4 and 5. As can be seen, there is much less scatter in the SAMUM and SALTRACE dust observations compared to the observations for the aerosol mixtures over the urban sites of Leipzig and Limassol.

The CCN-relevant correlation study ( $n_{100,d}$  vs  $\sigma_d$ ) is in good agreement with field observations of Shinozuka et al. (2015) at the dusty site of Niamey (Niger, western Africa, south of the Sahara). The simultaneous observation of CCNC and dry extinction coefficients observations (for the class with AE from 0.3–0.5) yield  $n_{CCN} \approx 110 - 120 \text{ cm}^{-3}$  for a supersaturation level of 0.2 % and an (ambient) dust extinction coefficient of  $\sigma_d = 50 \text{ Mm}^{-1}$ . Our parameterization yields  $n_{CCN} \approx 100 \text{ cm}^{-3}$  for  $\sigma_d = 50 \text{ Mm}^{-1}$  with the dust parameters in Table 3 inserted in Eqs. (1) and (10).

A clear linear relationship between  $n_{250,d}$  and  $\sigma_d$  is given which corroborates the usefulness of lidar for dust INP profiling after DeMott et al. (2010, 2015a). The almost linear behavior of  $s_d$  with  $\sigma_d$  also suggests that surface-area-based INP parameterization (Niemand et al., 2012; Steinke et al., 2015) for desert dust is possible with good accuracy. The relationship between  $n_{250,d}$  and  $\sigma_d$  values was already discussed by Mamouri and Ansmann (2015). However, as mentioned in Sect. 2.1, a wrong conversion factor was used in our foregoing article, the true one is  $c_{250,d} = 0.20 \text{ Mm cm}^{-3}$ .

### 4.3 Pure marine aerosol conditions

Barbados offers also the unique opportunity to analyze the relationship between the microphysical and optical properties for pure marine conditions. As mentioned, Barbados is located more than 4000 km west of Africa. No anthropogenic aerosol sources exist upwind Barbados over the tropical Atlantic (except ships). We selected 123 AERONET Barbados Ragged Point observations (level 2.0, 2007–2015) for our correlation study. To identify these pure marine conditions we used the criteria of AOT < 0.07 at 500 nm and AE between 0.25 and 0.6. The AE value for pure marine conditions is clearly higher than for Saharan dust and smaller for cases with local pollution (mainly biomass burning). The marine AE values accumulate at 0.45–0.55. The conversion parameters for pristine marine conditions are given in Table 3.

As mentioned, because of the high relative humidity around 80 % in the marine PBL, we use  $n_{100,m}$  as a proxy for  $n_{50,m,dry}$ . Similarly  $s_m/4$  is assumed to represent  $s_{m,dry}$ .

As for desert dust, the comparison with the CCN- $\sigma$  correlation studies of Shinozuka et al. (2015) show good agreement. On Graciosa Island (Azores), for marine conditions (in summer, AE between 0.3 and 0.5), on average,  $n_{CCN}$  between 400 and 500  $\text{cm}^{-3}$  was observed for a supersaturation of 0.3–0.5 % and a mean dry extinction coefficient of 20  $\text{Mm}^{-1}$ . For relative humidities of 80 % the ambient extinction coefficients are roughly a factor of 3 larger than the dry extinction coefficients, and thus around 60  $\text{Mm}^{-1}$  (Zieger et al., 2010; Zhang et al., 2014). By further assuming that the average  $n_{CCN}$  is approximately between 200 and 250  $\text{cm}^{-3}$  when changing the supersaturation level from 0.3–0.5 to 0.2 %, these transformed values are close to the ones obtained with our parameterization drawn from the AERONET observations. By using the parameters in Table 3 and Eqs. (3) and (12), we get  $n_{CCN} \approx 200 \text{ cm}^{-3}$  for an ambient  $\sigma_m = 50 \text{ Mm}^{-1}$ .

The good agreement between our parameterization and the CCNC- $\sigma$  correlation study of Shinozuka et al. (2015) suggests that our way to handle the water uptake effect by using  $n_{100,m}$  as a proxy for  $n_{50,m,dry}$  is reasonable. Similarly, the good agreement with the results of Shinozuka et al. (2015), discussed in Sect. 4.2, indicates that the use of  $n_{100,d}$  (in the case of hydrophobic dust particles) to estimate dust  $n_{CCN,d}$  is justified, too.

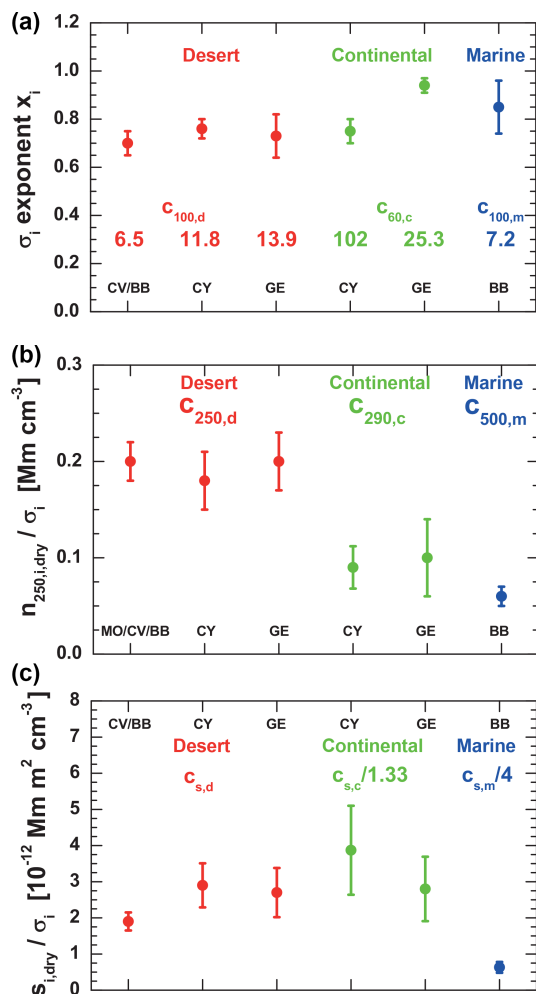
### 4.4 Continental, marine, and desert aerosol conversion parameters and uncertainties

Figure 7 provides finally an overview of all 532 nm mean conversion parameters for the three aerosol types and different AERONET data sets. Clear differences in the conversion parameters for dust, marine, and continental (urban) aerosol conditions are visible. The error bars are related to the atmospheric variability (scatter in the correlations shown in Figs. 4–6). In Table 3, all conversion parameters (with corresponding SD) required to solve the equations in Sect. 3 are given for the three laser wavelengths.

Typical uncertainties in the basic particle optical properties, the derived particle number and surface area concentrations, and the CCN-relevant particle and INP number concentrations are provided in Table 4. They result from uncertainties in the lidar retrievals (uncertainties in the input parameters, statistical noise), and retrieval uncertainties as discussed in Sects. 2, 3, and 4. The uncertainties are similar for all three laser wavelengths.

The parameterizations and corresponding uncertainties given in Table 4 hold for relative humidities up to about 80 % in the case of continental aerosols. The estimated  $n_{CCN}$  values are no longer trustworthy for higher relative humidities, i.e., for example in the subcloud layer (from 500 m below cloud base to cloud base), i.e., in the humid layer right below the base of a convective cloud system (Schmidt et al., 2014). It remains to be investigated to what extent our method can





**Figure 7.** Overview of derived values for (a) the extinction exponent  $x_i$  (required to compute  $n_{100,d,dry}$ ,  $n_{50,c,dry}$ , and  $n_{50,m,dry}$  with Eqs. 1–3), (b)  $c_{250,d}$ ,  $c_{290,c}$ , and  $c_{500,m}$  (required to compute  $n_{250,i,dry}$  with Eqs. 4–6) and (c)  $c_{s,i}$  (required to compute  $s_{i,dry}$  with Eq. (7)–(9) for 532 nm). Values for  $c_{100,d}$ ,  $c_{60,c}$ , and  $c_{100,m}$  (in  $\text{cm}^{-3}$  at  $\sigma = 1 \text{ Mm}^{-1}$  are required to solve Eq. 1–3) are given as numbers at the bottom of the top panel. Error bars (SD) indicate the uncertainties in the derived parameters. MO/CV/BB and CV/BB indicate SAMUM/SALTRACE dust observations, GE Leipzig, CY Limassol measurements, and BB (marine) denotes the Barbados Ragged Point 2007–2015 long-term observations.

be used for relative humidities  $> 80\%$  and also for humidities  $< 40\%$ .

Further uncertainty sources, not considered in Table 4, are the unknown updraft velocities at cloud base so that  $n_{CCN}$  can easily be underestimated by a factor of 2–3 when the supersaturation is between 0.4 and 1 instead of at 0.15 % (as we assume). In the case of mineral dust the amount of soluble material on the dust particle surface sensitively influences the ability to act CCN so that  $n_{CCN}$  of aged dust particles coated with hygroscopic species may be a factor of 2–4 higher than

**Table 4.** Typical uncertainties in the lidar-derived particle optical properties (for 532 nm wavelength), in the retrieved microphysical particle properties, and in the estimated cloud-relevant quantities.

Parameter		Relative uncertainty
Backscatter coefficient	$\beta_p$	5–10 %
Backscatter coefficient	$\beta_d$	10–15 %
Backscatter coefficient	$\beta_c$	10–20 %
Backscatter coefficient	$\beta_m$	20 % (PBL)
Extinction coefficient	$\sigma_d$	15–25 %
Extinction coefficient	$\sigma_c$	20–30 %
Extinction coefficient	$\sigma_m$	25 % (PBL)
Number concentration	$n_{50,i,dry}$	Factor of 1.5–2
Number concentration	$n_{100,i,dry}$	Factor of 1.5–2
Number concentration	$n_{250,i,dry}$	30–50 %
Surface area concentration	$s_{i,dry}$	30–50 %
CCN number concentration	$n_{CCN,ss,i}$	Factor of 2–3
INP number concentration	$n_{INP,i}$	Factor of 3–10

predicted by our parameterization. As mentioned in Sect. 3.3, the derived  $n_{CCN,ss=0.15\%}$  values can be interpreted as the minimum values of the possible solution space of  $n_{CCN,ss}$  for  $ss$  from 0.1 to 1 %.

Nevertheless, the consistency with the direct observations of CCNC and extinction coefficients by (Shinozuka et al., 2015) for marine, desert, and continental aerosol conditions corroborates that our lidar-based parameterizations are trustworthy. However, because the CCN retrievals based on the Leipzig and Limassol AERONET data sets for continental aerosols represent urban conditions, this parameterization may overestimate  $n_{CCN}$  in rural environments (aerosol background conditions) and probably also in the free troposphere (aged, long-range transported particles). It may be therefore advisable to use at least two sets of parameterizations for urban and rural sites and for the PBL (regional aerosol, high amount of freshly produced fine-mode particles) and the free troposphere (aged particles, partly originating from other continents) instead of just one generalized parameterization scheme. This aspect is further discussed in the next section.

## 5 Lidar estimates of $n_{CCN,ss}$ and $n_{INP}$ profiles: case studies

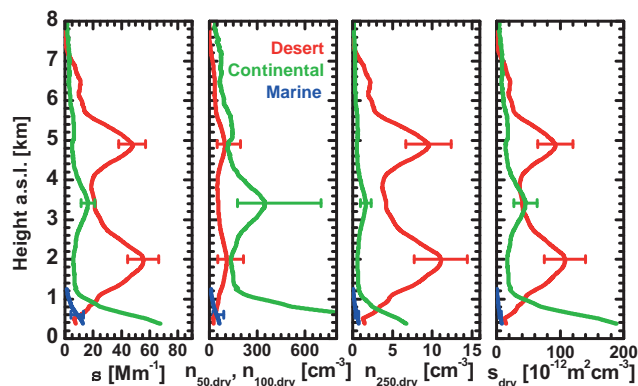
In this section, we apply the developed methodology presented in Sect. 3 to two lidar observations. The first lidar measurement was performed during a strong dust outbreak crossing Cyprus in the spring of 2015. The second case was measured during an episode with continental aerosol pollution advected from the European continent to Cyprus in the summer of 2012.

### 5.1 Lidar profiling during a dust outbreak

During the BACCHUS field campaign in March–April 2015, described in Sect. 2.2, many dust outbreaks from the Middle East deserts and the Sahara were observed over Cyprus. We selected the case from 7 April to apply our methods to a dust-dominated aerosol scenario. The basic lidar observations of height profiles of particle backscatter, linear depolarization ratio, and derived  $\sigma_i$  profiles were already shown in Fig. 2. The  $\sigma_i$  profiles for 532 nm wavelength are the input parameters for the retrieval of the particle number concentrations  $n_{50,\text{dry}}$ ,  $n_{100,\text{dry}}$ , and  $n_{250,\text{dry}}$ , and of the particle surface area concentration  $s_{\text{dry}}$  shown in Fig. 8. Equations (1)–(9) and the conversion parameters in Table 3 are used to obtain the presented profiles. For desert dust we used the SAMUM/SALTRACE conversion parameters, for continental pollution the Leipzig data, and for marine aerosols the Barbados (2007–2015) conversion parameters. The error bars indicate typical uncertainties (as summarized in Table 4). For  $n_{50,\text{dry}}$  and  $n_{100,\text{dry}}$  we simply assume an overall uncertainty factor of 2 in Fig. 8 (and in the following figures also for  $n_{\text{CCN}}$ ). For  $n_{250,\text{dry}}$  and  $s_{\text{dry}}$  the error bars show relative uncertainties of 30 % (dust) to 50 % (continental aerosol).

As can be seen in Fig. 8, the fine-mode-dominated continental aerosol fraction contains more CCN-relevant small particles ( $n_{50,\text{dry}}$  vs.  $n_{100,\text{dry}}$ ) than the desert aerosol, although the ambient extinction coefficients  $\sigma_c$  are smaller than the dust-related  $\sigma_d$  values. Vice versa, the coarse-mode-dominated dust aerosol controls the overall large-particle number concentrations  $n_{250,\text{dry}}$  and surface area concentration  $s_{\text{dry}}$ .

Figure 9 shows the retrieval products in terms of  $n_{\text{CCN}}$  and  $n_{\text{INP}}$ . In addition, the GDAS temperature and relative humidity (RH) profiles are given. Different  $n_{\text{CCN}}$  retrievals are presented. The profiles for GE(c) (thick green profile in Fig. 9) and for CY(c) (thin light green profile) are calculated with Eq. (11) and the conversion parameters for Leipzig (GE for Germany) and Limassol (CY for Cyprus) for continental aerosol (c) in Table 3, respectively. For comparison, also results obtained with the generalized parameterization scheme of Shinozuka et al. (2015) for a supersaturation level of 0.2 % and the AE class of 1.5–1.7 are plotted (SHI(c), thin green profile in Fig. 9). In our notation (according to Eq. 11), the SHI(c) parameters in Fig. 9 are  $c_{\text{CCN}} = 30 \text{ cm}^{-3}$  (for  $\sigma = 1 \text{ Mm}^{-1}$ ) and  $x_{\text{CCN}} = 0.75$ , and represent rural-like rather than urban aerosol conditions. Similarly, the correlation study of Sakai et al. (2013) based on vertical profiles of ambient particle backscatter coefficients measured with lidar over southern Japan and airborne in situ CCN observations, yield  $c_{\text{CCN}} = 30 \text{ cm}^{-3}$  (for  $\sigma = 1 \text{ Mm}^{-1}$ ), but  $x_{\text{CCN}} = 0.5$ . These observations also indicate aerosol background conditions (AE values mostly from 0.3–1.0). As mentioned above, we used an extinction-to-backscatter ratio of 50 sr to transfer the backscatter into extinction coefficients at 532 nm laser wavelength.

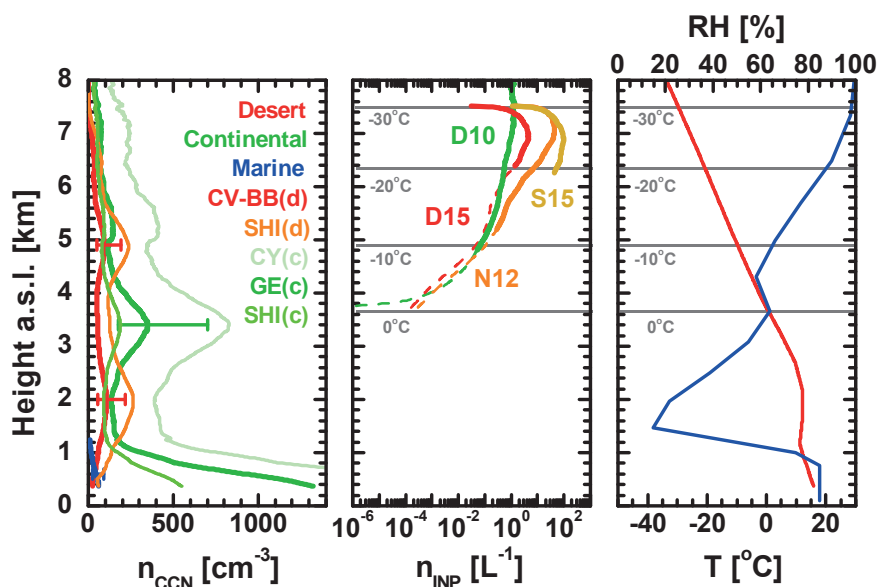


**Figure 8.** Height profiles of the 532 nm aerosol extinction coefficient  $\sigma$  (as shown in Fig. 2), and derived profiles of particle number concentrations  $n_{50,\text{dry}}$  (marine, continental) and  $n_{100,\text{dry}}$  (desert), of the large particle fraction in terms of  $n_{250,\text{dry}}$ , and surface area concentration  $s_{\text{dry}}$ , separately for all three aerosol types. The BACCHUS lidar observation was taken with Polly<sup>XT</sup> at Nicosia on 7 April 2015 during a major dust outbreak from the Sahara. Error bars show typical overall retrieval uncertainties.

The desert-dust-related  $n_{\text{CCN}}$  profile (thick red curve in Fig. 9) is calculated with Eq. (10) and the conversion parameters in Table 3, derived from the Cabo-Verde and Barbados AERONET observations of pure dust (denoted as CV-BB(d)). Again for comparison, the thin orange profile (SHI(d)) shows the  $n_{\text{CCN}}$  profile obtained with the generalized aerosol parameterization of Shinozuka et al. (2015) for the supersaturation level of 0.2 % and AE from 0.3–0.5. The SHI(d) conversion parameters are in this case  $c_{\text{CCN}} = 13 \text{ cm}^{-3}$  (for  $\sigma_d = 1 \text{ Mm}^{-1}$ ) and  $x_{\text{CCN}} = 0.75$ . It is interesting to note that the Limassol dust conversion parameters in Table 3 ( $c_{100,d} = 11.8 \text{ cm}^{-3}$  for  $\sigma_d = 1 \text{ Mm}^{-1}$ ,  $x_d = 0.76$ ) and the Leipzig dust conversion parameters ( $c_{100,d} = 13.9 \text{ cm}^{-3}$  for  $\sigma_d = 1 \text{ Mm}^{-1}$ ,  $x_d = 0.73$ ) are similar to the SHI(d) conversion parameters. All three parameterizations obviously represent slightly polluted desert dust conditions. For pure desert dust scenarios (CV-BB(d), thick red profile in Fig. 9, based on the Cabo-Verde and Barbados AERONET observations) the respective  $n_{\text{CCN}}$  values are lower by almost a factor of 2.

The different profiles for continental aerosols (GE(c), CY(c), and SHI(c)) provide an impression of the uncertainty in the  $n_{\text{CCN}}$  retrieval for this aerosol type. Similarly, the orange and red curves may indicate the overall uncertainty in the retrieval of  $n_{\text{CCN}}$  for desert dust.

The relative humidity profile indicates that RH is  $< 80 \%$  for the range from 1–6 km height. For this region, our parameterization (for RH around  $60 \pm 20 \%$ ) is valid. Care has to be taken in the interpretation of the continental  $n_{\text{CCN}}$  values in the PBL (RH  $> 80 \%$ ) and in the upper troposphere (above 6 km height, RH  $> 80 \%$ ).



**Figure 9.** (Left) Particle number concentration  $n_{\text{CCN}}$  for a supersaturation of  $ss = 0.15\%$  (obtained with Eqs. 10–12). Different parameterizations are used (thick green, GE(c), Germany conversion parameters in Table 3, light green, CY(c), Cyprus parameters, thick red, CV-BB(d), Cabo Verde and Barbados dust conversion parameters, see text for more explanations). In the case of SHI(c) (thin green) and SHI(d) (thin orange) the CCNC parameterization of Shinozuka et al. (2015) is applied. Blue line segment in the PBL shows the estimated marine contribution to CCNC. (Center) Ice-nucleating particle number concentration  $n_{\text{INP}}$ , computed with the parameterization schemes after DeMott et al. (2010) (D10, Eq. 13), DeMott et al. (2015a) (D15, Eq. 14), Niemand et al. (2012) (N12, Eqs. 15–16), and Steinke et al. (2015) (Eqs. 17–19). The respective particle input parameters  $n_{250,d,dry}$  and  $s_{d,dry}$  are shown in Fig. 8. Solid line segments show the temperature range for which the parameterizations were developed. (Right) GDAS temperature and relative-humidity profiles for Limassol, 7 April 2015, 21 UTC. Error bars (left panel) indicate the estimated uncertainties (factor of 2).  $n_{\text{INP}}$  errors are estimated to be within a factor of 3–10.

In the central panel of Fig. 9, the retrieved INP profiles are shown, obtained with the different parameterizations discussed in Sect. 3.4. Mamouri and Ansmann (2015) already discussed the retrieval of  $n_{\text{INP}}$  from  $n_{250,d}$  by using the D10 method (DeMott et al., 2010) and the D15 approach (DeMott et al., 2015a). Figure 9 also contains the  $n_{\text{INP}}$  profiles when the desert particle surface-area concentration  $s_d$  is used as input in the N12 approach (Niemand et al., 2012) and S15 parameterization (Steinke et al., 2015). Because air temperatures are all above  $0^\circ\text{C}$  at heights below 3.6 km as the horizontal temperature lines in Fig. 9 (right panel) indicate,  $n_{\text{INP}}$  values are only given for the upper part of the desert dust plume. According to Table 4, the uncertainty in the INP retrieval is within a factor of 3–10.

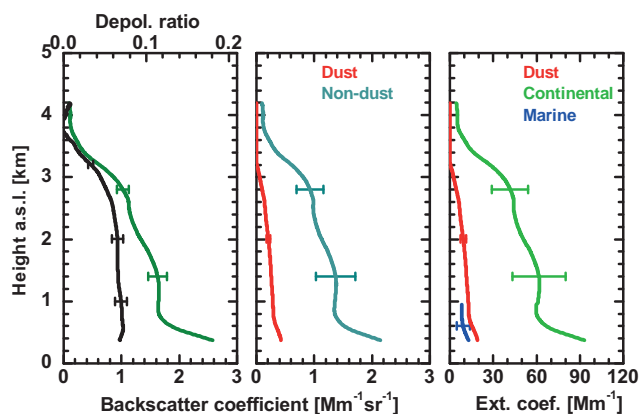
It is not the aim of the paper to discuss in detail the reasons for the differences between the different immersion freezing parameterizations of DeMott et al. (2015a) and Niemand et al. (2012), which partly exceed 1 order of magnitude. The higher  $n_{\text{INP}}$  values obtained with the procedure developed by Niemand et al. (2012) compared to the one presented by DeMott et al. (2015a) may result from the fact that  $s$  covers all particles even particles with radius  $< 250$  nm. The combination of the parameterizations of Niemand et al. (2012) (dust aerosol, immersion freezing) and Steinke et al. (2015) (dust, deposition freezing) provides the opportunity to deliver  $n_{\text{INP}}$

profiles from about  $-10$  to  $-50^\circ\text{C}$  and thus up to cirrus level. The parameterization scheme of Steinke et al. (2015) need to be tested for natural desert dust. As mentioned in Sect 3.4.2, it is based on laboratory studies with Arizona test dust.

At the end of this subsection, it is noteworthy to mention that similar profiles as shown in Figs. 8 and 9 can be obtained with a polarization lidar operated at the laser wavelength of 355 nm. The respective conversion parameters are given in Table 3. This means that ESA's EarthCARE lidar (satellite-borne 355 nm polarization/HSR lidar) (Illingworth et al., 2015a) can also provide these CCN and INP number concentration profiles, however, on a global scale.

## 5.2 Lidar profiling during an episode with European continental pollution

In contrast to the BACCHUS dust case in Fig. 2, the backscatter and depolarization profiles in Fig. 10 show a case with strong advection of aerosol from the European continental. The measurement was taken at Limassol on 16 August 2012. The AOT at 532 nm was about 0.2 and AE close to 1.8. Aerosols up to 3.5 km height were detected and the air masses came from northern to northwestern directions, from Europe and Turkey according to backward trajectory analysis. We used lidar ratios of 50–60 sr for continental pollution and 45 sr for mineral dust in the backscatter-to-extinction



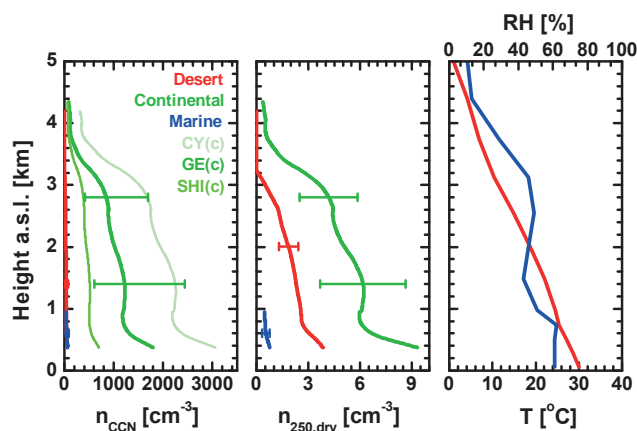
**Figure 10.** Same as Fig. 2, except for a lidar observation at Limassol on 16 August 2012. On this day, continental aerosol pollution from Turkey, the Black Sea area, and from southeastern and central Europe was advected to Cyprus at different heights up to 4 km. Lidar ratios used in the conversion of backscatter into extinction profiles were 50–60 sr for continental pollution and 45 sr for mineral dust.

conversion to obtain the  $\sigma_i$  profiles from the backscatter coefficients. We again assume a small contribution of marine particles in the boundary layer over the coastal city of Limassol. The almost height-independent particle linear depolarization ratio indicates an aged, well-mixed pollution plume. The values of the depolarization ratio are 5–8 %. Such values indicate the presence of a small amount of soil and road dust, or even traces of desert dust.

Figure 11 shows the derived profiles of the CCN-relevant particle number concentration  $n_{CCN}$  and of the INP-relevant  $n_{250,dry}$  values. The respective  $n_{INP}$  values are at all zero for this case with ambient temperatures  $> 0^\circ\text{C}$  up to 5 km height, as shown in the right panel.

Most interesting for such a pollution case in the lower troposphere are the  $n_{CCN}$  profiles. We show again the profiles for different parameterization. The Limassol conversion parameters (CY(c)), the Leipzig parameters (GE(c)), and the parameterization after Shinozuka et al. (2015) for the supersaturation level of 0.2 % and AE around 1.6 (SHI(c)) are applied. The contributions of the dust and marine aerosol components ( $20\text{--}40\text{ cm}^{-3}$ ) to  $n_{CCN}$  are almost negligible in this case.

The overall uncertainty in the retrieval of a factor of 3 is again well covered by the three different parameterizations. Because the relative humidity is mostly between 40 and 60 %, an RH-related error can be regarded to be small. As can be seen, even at moderate pollution levels with  $\sigma_c$  of  $30\text{--}60\text{ Mm}^{-1}$ , the  $n_{CCN}$  values can be of the order of  $1000\text{--}2000\text{ cm}^{-3}$  at urban sites. By using the parameterization of Shinozuka et al. (2015) (SHI(c) profile), which is more appropriate for rural aerosol conditions, we obtain  $n_{CCN} \approx 300\text{--}500\text{ cm}^{-3}$ .



**Figure 11.** (Left) Particle number concentration  $n_{CCN}$  for a supersaturation of  $ss = 0.15\%$ . The thick green (GE(c) conversion parameters), red, and blue profiles are obtained with Eqs. (10)–(12) for continental, desert, and marine aerosol, respectively. The thin light green profile (CY(c)) is obtained with Cyprus conversion parameters, and SHI(c) (thin green) with the parameterization of Shinozuka et al. (2015) (see text for more explanations). (Center) Large-particle number concentration  $n_{250,dry}$ , computed with Eqs. (4)–(6). (Right) GDAS temperature and relative-humidity profiles for Limassol on 16 August 2012.

## 6 Conclusions

For the first time, a comprehensive study on the potential of polarization lidar to provide vertical profiles of CCN-relevant particle and INP number concentrations has been presented. Of key importance is the separation of the basic aerosol types (desert, continental, marine) by means of the polarization lidar technique. Based on an in-depth correlation study applied to long-term and field campaign AERONET observations, it has been demonstrated that a solid path exists from the particle extinction coefficients, as measurable with lidar, to the basic aerosol parameters from which the  $n_{CCN}$  and  $n_{INP}$  profiles can be estimated.

We showed that height profiles of CCN-relevant number concentrations of aerosol particles with dry radius  $> 50\text{ nm}$  (marine and continental particles) and  $> 100\text{ nm}$  (desert dust), and of the INP-relevant number concentration of particles with dry radius  $> 250\text{ nm}$ , as well as profiles of the aerosol particle surface area concentration can be retrieved from lidar-derived ambient aerosol extinction coefficients with relative uncertainties of a factor of around 2 (CCN reservoir) and about 25–50 % (INP reservoir). The overall uncertainties in the  $n_{CCN}$  estimation of a factor 3 and in the  $n_{INP}$  estimation of a factor 3–10 result, to a large extent, from the unknown aerosol types and properties (origin of the aerosol components, chemical composition of the aerosol, aging and coating effects).

The full methodology was applied to two contrasting cases: a heavy desert dust outbreak crossing Cyprus with

mineral dust up to 8 km height in the spring of 2015 and a case with aged anthropogenic haze from the European continent. These case studies clearly demonstrated the attractiveness of lidar to provide simultaneously height profiles of  $n_{\text{CCN}}$  and  $n_{\text{INP}}$  estimates up to the mixed-phase and cirrus cloud level.

There is room for improvements. Our study may be regarded as a starting point for a deeper discussion on the role of lidar (organized in ground-based networks and operated in space) to provide height profiles of cloud-relevant aerosol parameters and to support in this way atmospheric research regarding the aerosol impact on cloud evolution and precipitation formation processes. It is an open question how to handle the water-uptake effect by the particles in the retrieval of the required dry-particle microphysical properties. Should one also move from lidar-derived ambient particle extinction coefficients to dry-particle extinction coefficients as in the study of Shinozuka et al. (2015)? Can we significantly improve the accuracy in the  $n_{\text{CCN}}$  and  $n_{\text{INP}}$  retrievals by combining the polarization lidar technique with the Raman lidar technique for water vapor and temperature profiling so that actual height profiles of relative humidity (Mattis et al., 2002) are available for the necessary water-uptake corrections? Do we need at least two CCN parameterization schemes to cover contrasting environments (rural vs. urban sites, PBL vs. free tropospheric height regions)? This seems to be obvious. In this context we may follow the way of Sakai et al. (2013) to use the AE profile (measured in the case of a multiwavelength lidar) as a guide in the selection of the most appropriate CCN parameterization scheme. Regarding  $n_{\text{INP}}$  profiling, the following question arises: in which way we may better consider the different INP efficacy of different aerosol types in the INP parameterizations, especially in cases with good knowledge on the amount of biological particles, biomass burning smoke, or urban haze in observed complex aerosol mixtures as a function of site, season of the year, and height range in close combination with backward trajectory analysis or more complex aerosol transport modeling? All these questions need to be answered in followup studies.

Our future plans comprise extended comparisons of the lidar observations of  $n_{\text{CCN}}$  and  $n_{\text{INP}}$  profiles with respective surface and airborne in situ observations of these quantities. The efforts should also include comparisons of the basic aerosol parameters such as  $n_{50, \text{dry}}$ ,  $n_{100, \text{dry}}$ , and  $n_{250, \text{dry}}$ , and the surface area concentration  $S_{\text{dry}}$ . Several measurement campaigns and long-term monitoring aerosol data sets will be used in these quality assurance activities. We will, e.g., compare the lidar retrieval products with aircraft measurements of desert-dust and marine-aerosol-related CCNC profiles, collected during the SALTRACE-1 campaign (Barbados, 2013) and with ground-based in situ  $n_{\text{INP}}$  observations during the BACCHUS campaign (Cyprus, 2015).

Furthermore, it is time for well-designed INP campaigns with aircraft measurements around laser beams. Airborne in situ observations (including aboard unmanned aerial vehi-

cles – UAVs) in a desert environment, at pure marine conditions, and at mixed aerosol conditions would be desirable. The aerosol components (origin, chemical composition), the particle size distributions, and INP number concentrations need to be measured in the vicinity or around the laser beams of a polarization lidar. Such field campaigns would provide ideal conditions for in-depth characterization of the potential of lidar-based INP parameterization efforts. This would also provide the unique opportunity to identify the gaps in our knowledge regarding heterogeneous ice formation when combined with cloud observations. The lidar monitors the evolution of cloud layers (altocumulus and cirrus layers) embedded in the aerosol layers from cloud base to cloud top, whereas aircraft can probe the aerosol and cloud layers, height level by height level, in terms of ice crystal and INP number concentrations.

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