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Supplement of

Importance of aerosol composition and mixing state for cloud droplet activation over the Arctic pack ice in summer

C. Leck and E. Svensson

Correspondence to: C. Leck (lina@misu.su.se)

1 **Supplementary material**

2 We developed a new implementation (written in Matlab) of the adiabatic parcel model
3 described by Pruppacher and Klett (1997). It consists of essentially the same equations as the
4 model presented in Leitch et al. (1986), with the difference that we use the full implicit
5 dynamic diffusional growth equation for particles in each size bin:

6

$$7 \quad r \frac{dr}{dt} = \frac{D' M_w e_{sat,w}}{\rho_s RT} \left(S - \frac{1}{1+\delta} \exp \left[\frac{L_e M_w}{RT} \left(\frac{\delta}{1+\delta} \right) + \frac{2 M_w \sigma_{s/a}}{RT(1+\delta) \rho_w r_i} - \frac{\nu \Phi_s \varepsilon_m M_w \rho_N r_N^3}{M_s^2 (r^3 - r_N^3)} \right] \right) \quad (1),$$

8 where:

9 r - radius of particle

10 t - time

11 M_w - molecular weight of water

12 $e_{sat,w}$ - saturation pressure of water vapor above a water surface

13 ρ_s - density of the droplet solution

14 R - universal gas constant

15 T - temperature of the gas surroundings

16 S - supersaturation

17 L_e - latent heat of evaporation of water

18 $\sigma_{s/a}$ - surface tension at solution/air interface

19 ρ_w - density of pure water

20 ν - ion number of the salt

21 Φ_s - osmotic coefficient of droplet solution

22 ε_m - mass fraction of soluble material

23 ρ_N - density of dry nucleus

24 M_s - molecular mass of nucleus material

1 $r_{N,1}$ – radius of dry nucleus

2

3 $\delta = T_d/T - 1$, where T_d is the droplet temperature, gives a measure of the heating or
4 cooling caused by latent heat released when the droplet grows or shrinks. It is given by:

5

6
$$\delta = \frac{L_e \rho_s}{TK'} r \frac{dr}{dt} \quad (2)$$

7

8 The D' and K' are modified diffusion and thermal conduction coefficients:

9
$$D' = \frac{D}{\left[\frac{r}{r + \lambda} + \frac{D}{r\alpha_c} \sqrt{\frac{2\pi M_w}{RT}} \right]} \quad (3)$$

10

11
$$K' = \frac{k}{\left[\frac{r}{r + \lambda} + \frac{k}{r\alpha_t \rho c_{pa}} \sqrt{\frac{2\pi M_A}{RT}} \right]} \quad (4)$$

12

13 where:

14 D – pressure and temperature dependent diffusivity of water vapor in air

15 λ – constant ($1.2 \cdot 10^{-7}$ m)

16 α_c – condensation accommodation coefficient

17 k – heat conductivity

18 α_t – condensation accommodation coefficient

19 M_A – mean molecular mass of air

20

1 We simplify the Raoult term in the particle growth equation by identifying it as the water
2 activity.

3

$$\hat{M}\Phi = -\ln(a_w)$$

$$4 \quad \frac{\nu\Phi_s \varepsilon_m M_w \rho_N r_N^3}{M_s \rho_w (r^3 - r_N^3)} = \nu M \quad (5)$$

5 Recent research has shown that both of the accommodation coefficients are likely to have
6 values close to unity for pure water surfaces (e.g. Winkler et al., 2004, Winkler et. al.,
7 2006, Morita et al., 2004).

8 Due to the dependence of δ on dr/dt , the system of ordinary differential equations is implicit,
9 and hence we need to provide both starting values of T, p, S, w and r_i and their time
10 derivatives. There are a few numerical caveats in the model, such as particle sizes shrinking
11 below the dry size or even becoming negative. Unphysical developments like these are
12 avoided by artificially setting the particle size derivatives positive when the droplet size
13 approaches the size of the dry particle.

14 In contrast to the model of Leitch et al. (1986) we make no assumptions of log-normal
15 aerosol distributions. Instead, we directly as input use the measured size distributions by the
16 TDMPS. The lower cut-off for the chemical information is 22 nm in diameter, which is well
17 below the size required to act as CCN in the conditions of his study. Chemical mass
18 concentration data from LPI impactor samples are interpolated on to the TDMPS size bins
19 and thereafter converted to number concentration assuming spherical particles.

20

$$21 \quad r_i \frac{dr_i}{dt} = \frac{D' M_w e_{sat,w}}{\rho_s RT} \left(S - \frac{1}{1+\delta} \exp \left[\frac{L_e M_w}{RT} \left(\frac{\delta}{1+\delta} \right) + \frac{2M_w \sigma_{s/a}}{RT(1+\delta)\rho_w r_i} - \frac{\nu\Phi_s \varepsilon_m M_w \rho_N r_{N,i}^3}{M_s \rho_w (r_i^3 - r_{N,j}^3)} \right] \right) \quad (6)$$

22

1 For an isolated air parcel ascending adiabatically with the vertical speed V , the
2 temperature T , pressure p , supersaturation S and liquid water content are described by the
3 equations:

4

5

6

$$7 \quad -\frac{dT}{dt} = \frac{gV}{c_{pa}} + \frac{L_e}{c_{pa}} \frac{dw}{dt} \quad (7)$$

8

$$9 \quad \frac{dp}{dt} = -\frac{gpV}{R_a T} \quad (8)$$

10

$$11 \quad \frac{dS}{dt} = \frac{p}{\varepsilon e_{sat,w}} - (1+S) \left[\frac{L_e}{R_a T^2} \frac{dT}{dt} + \frac{gV}{R_a T} \right] \quad (9)$$

12

$$13 \quad \frac{dw}{dt} = \sum_i 4\pi m_i r_i^2 \frac{dr_i}{dt} \quad (8)$$

14

15 In this study, however, we use $V=0$ and artificially change S in the way the CCN counter
16 does.

17