

1 **Supplementary Information for**

2 **CCN activity of organic aerosols observed downwind of**  
3 **urban emissions during CARES**

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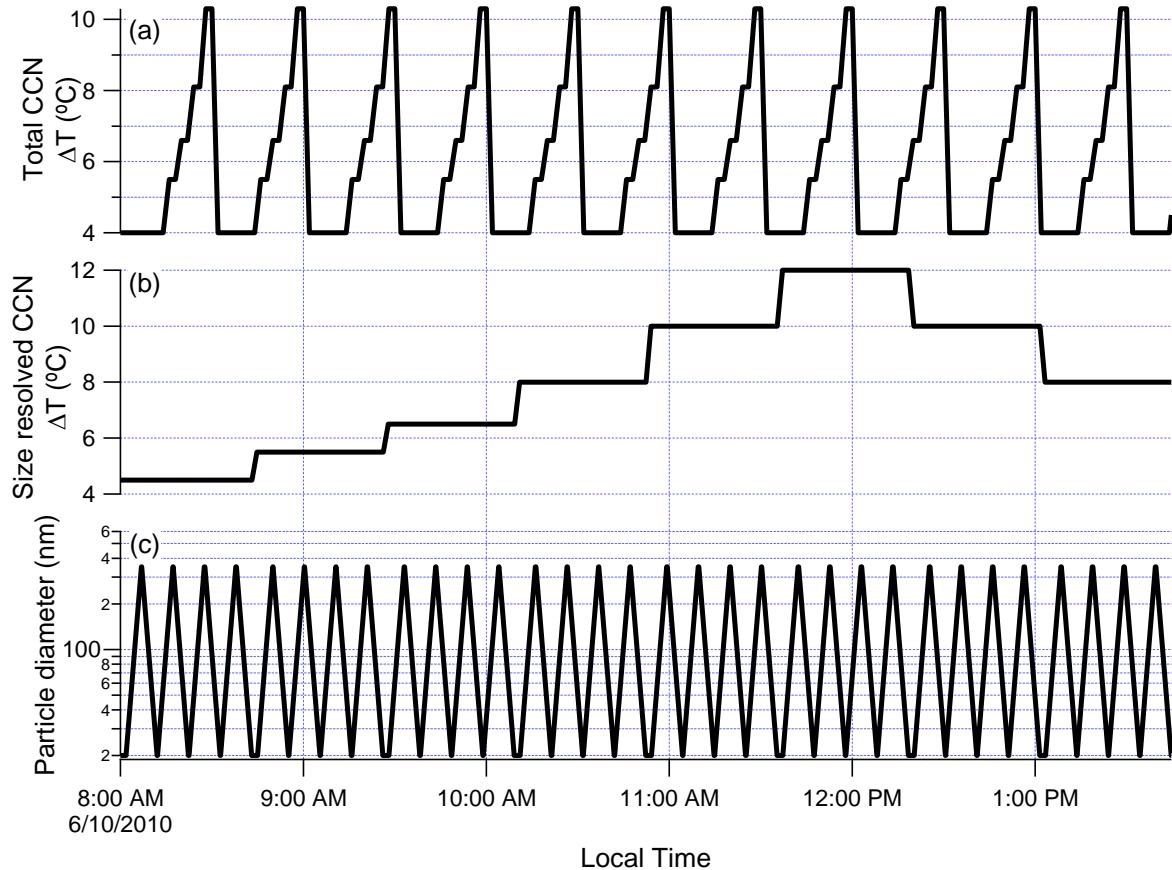
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1    **1. Measurement sequence**



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3    Figure S1. An example of measurement sequences for (a) the CCN counter temperature gradient  
4    for total CCN concentration measurements, (b) the CCN counter temperature gradient and (c) the  
5    particle size classified by DMA for the size resolved CCN measurements.

6

7    **2. Derivation of particle hygroscopicity and mixing state**

8       The activated fractions measured at the six supersaturations were fitted using following  
9    two different functions (Mei et al., 2013):

10      
$$R_a(S) = \frac{E}{2} \cdot \left( 1 + \operatorname{erf} \left( \frac{\ln S - \ln S^*}{\sqrt{2\sigma_s^2}} \right) \right) \quad (\text{S1})$$

1 and (Lance, 2007; Cerully et al., 2011; Bougiatioti et al., 2011; Padro et al., 2012; Lance et al.,  
2 2013):

$$3 R_a(S) = \frac{E}{1 + \left(\frac{S}{S^*}\right)^C} \quad (S2)$$

4 The fitting parameters are  $E$ ,  $S^*$ , and  $\sigma_S$  for Eqn. (S1) and  $E$ ,  $S^*$ , and  $C$  for Eqn. (S2), where  $\sigma_S$   
5 and  $C$  are related to the slope of the increasing  $R_a$  with  $S$  near  $S^*$ . For each set of measurements,  
6 the function form that yielded the best fit (i.e. smaller least squares residue) was used for  
7 subsequent analysis.

8 For particles with the same size and composition (i.e., hygroscopicity), we would expect  
9 a step function for  $R_a$  as all particles would have the identical  $S_c$ . Ambient aerosols show much  
10 more gradual increase in  $R_a$  (i.e., instead of a step change), suggesting heterogeneity in particle  
11  $S_c$ . The probability density function (PDF) of the critical supersaturation for size selected  
12 particles,  $p(S_c)$  is given by differentiating  $R_a(S_c)$  with respect to  $S_c$ :

$$13 p(S_c) = \frac{1}{E} \cdot \frac{dR_a(S_c)}{dS_c} \quad (S3)$$

14 The dispersion in  $S_c$  is defined as  $\sigma(S_c)/\bar{S}_c$ , where  $\bar{S}_c$  is the average particle critical  
15 supersaturation:

$$16 \bar{S}_c = \int_0^{\infty} p(S_c) \cdot S_c \cdot dS_c \quad (S4)$$

17 and

$$18 \sigma^2(S_c) = \int_0^{\infty} (S_c - \bar{S}_c)^2 p(S_c) dS_c \quad (S5)$$

1 When  $R_a(S)$  is fitted using Eqn. (S1), the hygroscopicity dispersion is:

2 
$$\frac{\sigma(S_c)}{S_c} = \left[ e^{\sigma_s^2} - 1 \right]^{1/2} \quad (S6)$$

3 and for Eqn (S2), the dispersion is given by:

4 
$$\frac{\sigma(S_c)}{S_c} = \left[ \frac{\Gamma\left(\frac{2}{C}+1\right) \cdot \Gamma(1-2/C)}{\Gamma(2)} \right] \left/ \left( \frac{\Gamma\left(\frac{1}{C}+1\right) \cdot \Gamma(1-1/C)}{\Gamma(2)} \right)^2 - 1 \right. \right]^{1/2} \quad (S7)$$

5 The dispersion in  $S_c$  is due to the combination of the width of DMA transfer function  
6 (particles classified by DMA do not have exactly the identical size) and the heterogeneity in  
7 particle composition (i.e., hygroscopicity), and can be expressed as (Lance et al., 2013):

8 
$$\left( \frac{\sigma(S_c)}{S_c} \right)^2 = \frac{9}{4} \left( \frac{\sigma(D_p)}{D_p} \right)^2 + \frac{1}{4} \left( \frac{\sigma(\kappa)}{\kappa} \right)^2 \quad (S8)$$

9 Where the first term on the RHS of the equation represents the contribution due to the width of  
10 DMA transfer function, which was estimated from the dispersion in  $S_c$  measured during  
11 calibration using  $(\text{NH}_4)_2\text{SO}_4$  particles (i.e., the contribution of the second term was essentially  
12 zero during calibrations). The dispersion in hygroscopicity for classified ambient particles was  
13 then derived by subtracting the contribution of DMA transfer function from the total dispersion  
14 in measured critical supersaturation.

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1    **References**  
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