

Emission and dry deposition of accumulation mode particles in the Amazon Basin

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Abstract. Size-resolved vertical aerosol number fluxes of particles in the diameter range 0.25–2.5 μm were measured with the eddy covariance method from a 53 m high tower over the Amazon rain forest, 60 km NNW of Manaus, Brazil. This study focuses on data measured during the relatively clean wet season, but a shorter measurement period from the more polluted dry season is used as a comparison.

Size-resolved net particle fluxes of the five lowest size bins, representing 0.25–0.45 μm in diameter, were in general dominated by deposition in more or less all wind sectors in the wet season. This is an indication that the source of primary biogenic aerosol particles may be small in this particle size range. Transfer velocities within this particle size range were observed to increase linearly with increasing friction velocity and increasing particle diameter.

In the diameter range 0.5–2.5 μm , vertical particle fluxes were highly dependent on wind direction. In wind sectors where anthropogenic influence was low, net upward fluxes were observed. However, in wind sectors associated with higher anthropogenic influence, deposition fluxes dominated. The net upward fluxes were interpreted as a result of primary biogenic aerosol emission, but deposition of anthropogenic particles seems to have masked this emission in wind sectors with higher anthropogenic influence. The net emission fluxes were at maximum in the afternoon when the mixed layer is well developed, and were best correlated with horizontal wind speed according to the equation

$$\log_{10} F = 0.48 \cdot U + 2.21$$

where F is the net emission number flux of 0.5–2.5 μm particles [$\text{m}^{-2} \text{s}^{-1}$] and U is the horizontal wind speed [ms^{-1}] at the top of the tower.

1 Introduction

In the Amazon Basin, organic components typically constitute 70 to 90% of the aerosol mass in both the fine and coarse mode (Andreae and Crutzen, 1997; Graham et al., 2003a; Fuzzi et al., 2007). Biomass burning is known to be a large source of organic particles (Reid et al., 2005), thereby explaining the high organic fraction in the dry season. In the wet season, however, the high organic fraction must be attributed to biogenic sources (Guyon et al., 2003a, b), but there are still large uncertainties in the relative importance of different sources and production mechanisms of these biogenic aerosol particles. Measurements of mass concentrations of different aerosol compounds and their diurnal variations in the Amazon boundary layer have provided important knowledge of biogenic aerosol sources in the Amazon Basin (e.g. Artaxo and Maenhaut, 1990; Artaxo and Hanson, 1995; Graham et al., 2003a, 2003b, Chen et al., 2009). However, diurnal variations in concentration of particularly those compounds that have their source at the surface are to a large extent driven by diurnal variations in boundary layer dynamics. Therefore it is necessary to measure vertical fluxes to gain more detailed knowledge of diurnal variations in surface emission and also for quantifying the source strength of individual components.

Vertical fluxes of different volatile organic compounds like isoprene and monoterpenes have been measured over the Amazon rain forest in a number of studies (e.g. Zimmerman et al., 1988; Karl et al., 2007; Kuhn et al., 2007) with the primary goal of improving the knowledge of secondary organic aerosol production. However, fewer studies have focused on the source mechanisms behind primary biogenic aerosol emission, and extremely few studies have done so by measuring vertical fluxes of aerosol particles. Ahlm et al. (2009) investigated surface-atmosphere exchange of particles with



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diameter larger than 10 nm over the Amazon rain forest in the wet season. In that study, it was found that deposition dominated the net particle number flux even in the cleanest prevailing conditions in the wet season. That was an indication that the source of primary biogenic aerosol emission, in terms of number concentrations, may be low when considering the total size range of particles. Also in the dry season, deposition fluxes dominated most of the time (Ahlm et al., 2010). However, upward particle fluxes often appeared in the morning. Rizzo et al. (2010) also observed upward particle number fluxes in the morning simultaneously with increasing particle number concentration, during aerosol flux measurements in the dry season at another tower at the same site. These early morning upward fluxes may have been related to primary biogenic aerosol emission of particles that have been stored in the canopy throughout the night, similar to CO₂ (Goulden et al., 2006; Araújo et al., 2008; Tóta et al., 2008), and the fact that these upward fluxes were sometimes correlated with increasing number concentrations suggest that they cannot be explained only as a result of dilution caused by entrainment during mixed layer growth.

In this study, results from measurements of size-resolved fluxes of accumulation mode particles over the Amazon rain forest are presented. The measured particles have diameters between 0.25 and 2.5 μm . Particles larger than 0.25 μm only makes up a minor fraction of the total aerosol number population in the Amazon boundary layer, but their size make them efficient as cloud condensation nuclei (CCN). The primary goal of this study is to investigate the number source of primary biogenic aerosol particles in different particle diameter intervals. In addition, triggering mechanisms of biogenic aerosol emission and diurnal variations in both flux and concentration are investigated. Finally, dry deposition is analyzed for those particle size ranges that are dominated by deposition fluxes.

Most data to be presented have been measured during the wet season when biogenic sources are expected to dominate the aerosol population within the atmospheric boundary layer. Additionally, data from measurements during a shorter period in the dry season is used as a comparison to the wet season data. This Brazilian-Swedish project AMAFLUX (Amazonian Biosphere-Atmosphere Aerosol Fluxes in view of their potential control of cloud properties and climate) was carried out as a part of the larger international project LBA (The Large Scale Biosphere Atmosphere Experiment in Amazonia).

2 Method

2.1 Site description

The measurements were carried out at the top of the 53 meter high tower K34 at the rain forest site Reserva Biológica do Cuieiras (2° 35.37' S, 60° 06.92' W), approximately 60 km

NNW of Manaus, Brazil. The tower K34 is a research facility operated by INPA (The Brazilian National Institute for Research in Amazonia). The canopy height in the Cuieiras Reserve is between 30 and 35 m (Kruijt et al., 2000). Southeasterly winds are associated with air transport from Manaus and thereby some anthropogenic influence. The centre of the research station is located approximately 2 km NNW of K34. A diesel generator at the centre of the station provides the research station with electrical power. A more detailed description of the site can be found in Ahlm et al. (2009).

2.2 Eddy covariance measurements

The eddy covariance method was used to measure the mean vertical turbulent aerosol number flux $\overline{N'w'}$, where N' and w' represent fluctuations in aerosol number concentration and vertical wind speed from the temporal means of these parameters, and the cross bar represents a temporal mean of the product of the two fluctuations. $\overline{N'w'}$ was calculated over periods of 30 minutes. The fluctuations N' and w' were separated from the mean by linear de-trending, which also removes the influence of low frequency trends. The aerosol data were shifted in relation to the wind data to correct for the time lag in the sampling line (calculated from the maximum correlation).

The eddy covariance method requires stationary conditions. In this study, the instationarity test by Foken and Wichura (1996) was applied to the particle fluxes in order to filter out fluxes measured in non-stationary conditions. The averaging period 30 min was divided into sub-periods of 5 min. If the difference between the flux calculated over 30 min and the mean of the covariances calculated over the 5 min intervals was larger than 60% (Järvi et al., 2009), the flux was rejected. In total, 22% of the particle fluxes were rejected.

Data collected when the friction velocity (u_*) is lower than 0.1 m s^{-1} have been rejected in figures that focus on the sign of the flux, in order to reduce the uncertainty of the flux. However, u_* -filtering may result in an overestimation of the average flux. Therefore, u_* -filtering has not been applied in figures with focus on the magnitude of the flux.

All aerosol data collected during rainfall have been excluded from the results in order to simplify interpretation of the flux. Technical problems due to instrumental failure, computer software problems and electrical power breaks resulted in some loss of data. The removal of data due to rainfall together with data losses due to technical problems resulted in a total data removal of 43% in the wet season period. In the short dry season period, the corresponding removal was only 14%.

2.3 Instrumentation

2.3.1 Flux measurements

The 3-D wind components and temperature were measured with a Gill Windmaster Pro 1561 ultrasonic anemometer, and logged at 20 Hz. An Optical Particle Counter (OPC), model Grimm 1.109, was used to measure the aerosol number concentration in 15 size classes from 0.25 μm to 2.5 μm at 1 Hz, using a separate 1/4-inch stainless steel 2.5 m long sampling line with a flow rate of 1.2 Lmin⁻¹ (laminar flow). The particle measurements were performed inside a weather housing that contained one OPC and a drying system. The sampled air was dried by 1:1 diffusion of particle free air with zero humidity. This method minimizes the risk of unwanted losses of semi-volatile compounds, compared to when the air is heated to dry it. The drying procedure before sampling removed 50% of the water vapor in the atmosphere, thereby lowering the relative humidity to 50% of its atmospheric value (since the partial pressure of water vapor is directly proportional to the number of moles of water). Rissler et al. (2006) has investigated hygroscopic properties of the aerosol in the Amazon and found no step-like deliquescent behavior of the aerosol in the dry or wet period. The particle growth factor could in that study be described by following relation:

$$Gf = \left[1 + A \cdot \left(\frac{\text{RH}/100}{1 - \text{RH}/100} \right) \right]^{1/3} \quad (1)$$

where RH represents relative humidity and A is a factor with a value of ~ 0.1 for the most hygroscopic particles ($D_p > 250$ nm). At a relative humidity of 90%, Gf is ~ 1.24 , but at a relative humidity of 50%, Gf is only ~ 1.03 . Since the drying procedure in this study will keep the relative humidity below 50%, the maximum hygroscopic growth for the most hygroscopic particles in the most humid conditions will not exceed 3%. In daytime, when the average atmospheric relative humidity was 70–80% in this study, the hygroscopic growth of the most hygroscopic particles will be only 1–2%. The error due to hygroscopic growth (Fairall, 1984; Kowalski, 2001; Vong et al., 2004) can be estimated by calculating the saturation ratio flux (Fairall, 1984) and by assuming a particle size distribution in the Junge form (Junge, 1963). This error can be large when the slope of the number distribution is steep. However, due to the drying procedure in this study, the error due to hygroscopic growth was estimated to be below 10%. In addition to water vapor, semi-volatile organic and inorganic vapors may result in particle growth and shrinkage, leading to spurious number fluxes for fixed size-bins, as pointed out by Nemitz and Sutton (2004) and Nemitz et al. (2009). Due to the very large uncertainty in particle growth rates in the Amazon Basin (Martin et al., 2010), these potential errors are harder to quantify, though.

Aerosol number concentrations decreased rapidly with increasing particle size within the OPC size range. In the channels representing the larger diameters, average number concentrations were lower than 0.1 cm⁻³, which is the resolution of the Grimm 1.109. Therefore, size resolved fluxes have only been calculated for the first five OPC channels, where number concentrations are at least one order of magnitude higher than the resolution of the OPC. The first five channels represent a size interval of $D_p = 0.25\text{--}0.45$ μm , where D_p represents particle diameter. However, by summing up several of the size bins that contain larger particle in order to increase the number of particles, also fluxes of the larger particles could be calculated. These fluxes were calculated over the interval $D_p = 0.5\text{--}2.5$ μm , representing the nine highest OPC size bins. Additionally, particle fluxes covering the whole OPC size range have been calculated.

2.3.2 Additional measurements used during data analysis

Mass concentration of equivalent black carbon (BC_e) was provided by São Paulo University using a Thermo Scientific Multi-Angle Absorption Photometer (MAAP), model 5012. This measurement derives the concentration of BC_e from the determination of light absorption at a wavelength of 637 nm. The BC_e measurements were performed close to the centre of the research station, about 2 km from K34.

Additionally, mass concentrations of trace elements were provided by São Paulo University using particle-induced X-ray emission analysis (PIXE). The sampling for the PIXE trace element analysis was also performed at the centre of the research station.

Meteorological parameters (rain amounts and photosynthetic active radiation) were measured at the K34 tower and provided by INPA. These were logged on a Campbell CR-10 (Campbell Scientific UK) data logger with a sampling interval of 30 s and stored as either 10 or 30 min averages.

2.4 Flux corrections

2.4.1 Corrections

Particle fluxes measured with the eddy covariance method are underestimated due to the limited time response of the OPC and attenuation of turbulent fluctuations in the sampling line. The time response of a first-order instrument for a step change in concentration from c_1 to c_2 can be described as (Doebelin, 1990)

$$c_m(t) = c_2 + (c_1 - c_2)e^{-t/\tau} \quad (2)$$

where t represents time, τ is the first-order time constant of the instrument, and $c_m(t)$ is the measured particle concentration by the instrument. The first-order time constant of the Grimm 1.109 was determined using a zero-filter and by estimating the time response for a step decrease in concentration

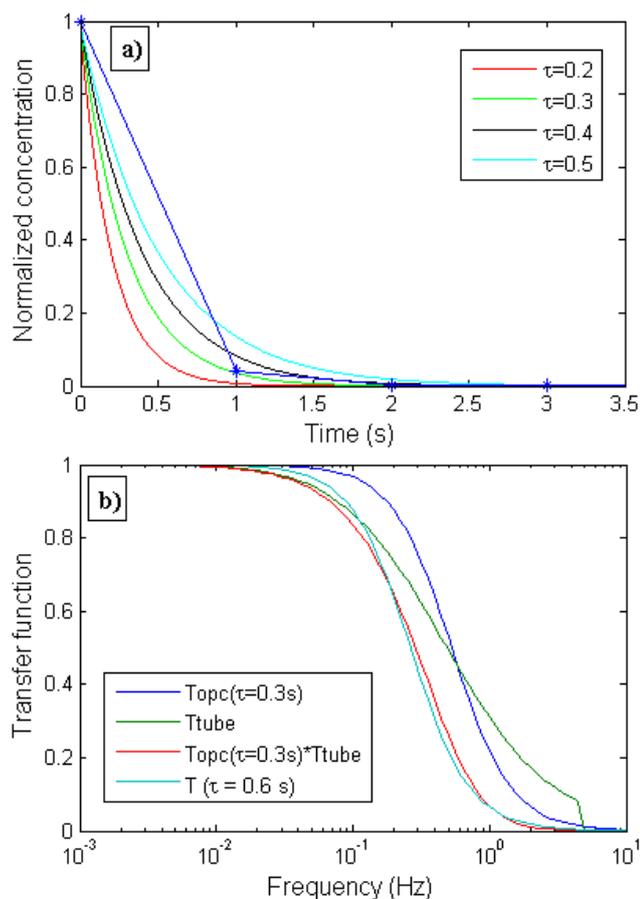


Fig. 1. (a) Time response of the Grimm 1.109 for a step decrease in particle concentration (blue) and time response curves for different values on τ (Doebelin, 1990). —bf(b) Cospectral transfer functions for the OPC (Horst et al., 1997) (blue), tube (Lenschow and Raupach, 1991) (green) and in total (red) compared with a cospectral transfer function (Horst et al. 1997) for $\tau = 0.6$ (cyan).

to zero. The result can be seen in Fig. 1a, where $\tau = 0.3$ gives the best fit to the data points. The experiment was repeated several times, with similar results, so the estimation of τ seems reliable. The total time constant for both OPC and sampling line was estimated using cospectral transfer functions (Fig. 1b) for damping of particle fluctuations in laminar flow (Lenschow and Raupach, 1991) and in a sensor (Horst et al., 1997). Multiplying these two transfer functions gave a total time constant of $\tau = 0.6$.

The aerosol fluxes in this study have been corrected for these fluctuation attenuations according to Horst et al. (1997). The median increase in flux due to this correction was 5% in the wet season. The 25 and 75 percentiles were 3 and 14%, respectively. The increase in flux in the dry season was also 5%, and the 25 and 75 percentiles were 3 and 22%.

In addition, aerosol fluxes and concentrations were corrected for losses in the sampling line. The sampling line was

very close to vertical so losses due to settling were negligible. However, the inlet of the tubing was bent downwards to prevent rainwater entering the sampling line. The sampling line was also slightly bent at two other points. For laminar flow, the inertial deposition of particles due to tube bends is given by the equation (Crane and Evans, 1977)

$$\text{Bend loss} = Stk \cdot \frac{\Phi}{2} \quad (3)$$

where Stk is Stokes number and Φ is the bend angle in radians. All size resolved concentrations and fluxes in this study have been corrected according to Eq. (3) for the three bends. The relative losses due to these bends were less than 0.05% for the first OPC channel ($0.25 < D_p < 0.28 \mu\text{m}$), gradually increasing with increasing D_p within the OPC size range, to reach 1.6% for the last OPC channel ($2.0 < D_p < 2.5 \mu\text{m}$). The reason for the small inertial loss is the laminar flow.

Losses due to diffusion were estimated using relations for diffusion in laminar flow in Hinds et al. (1999). The average diffusional loss was 0.4% for the first OPC channel and 0.07% for the last OPC channel. Since these values are small, the impact of diffusional losses within the sampling line has been neglected.

No Webb correction (Webb et al., 1980) has been applied to the fluxes. Temperature fluctuations are dampened in a tubing of several meters (Rannik et al., 1997) and therefore the particle fluxes were not corrected with regard to sensible heat fluxes. Nor were the fluxes corrected with regard to latent heat fluxes, since the sampled air was dried before entering the OPC.

2.4.2 Error treatment

The uncertainty in number concentration due to discrete counting was calculated from the equation (Hinds et al., 1999)

$$\delta(N) = \frac{1}{\sqrt{n}} \quad (4)$$

where $\delta(N)$ is the relative error in number concentration for half-hourly averaged data and n is the number of counts per half hour. The median relative uncertainty in particle concentration of the first OPC channel was 0.14% and 0.07% in the wet and dry season, respectively. The corresponding relative uncertainties for the last OPC channel, where the number of particles is much lower, were 2.3% and 1.8%.

The uncertainty in flux due to discrete counting can be expressed as

$$\overline{\delta(N'w')} = \frac{\sigma_w \bar{N}}{\sqrt{\bar{N} Q \Delta t}} \quad (5)$$

where σ_w is the standard deviation of the vertical wind, \bar{N} is the aerosol number concentration averaged over the sampling period Δt and Q is the sampling volume flow rate through

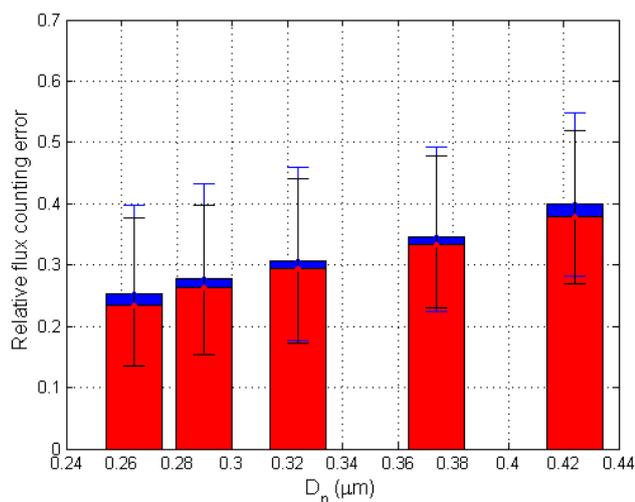


Fig. 2. Median relative counting error (bars) in particle flux for the five first OPC channels during the wet (blue) and dry (red) season. Error bars represent 25 and 75 percentiles of the counting error in the wet (blue) and dry (black) season

the particle counter (Fairall, 1984). Thus, the relative uncertainty ε in flux due to discrete counting becomes

$$\varepsilon(\overline{N'w'}) = \frac{\sigma_w}{\sqrt{Q\Delta t}} \cdot \frac{1}{\sqrt{\overline{N}} \cdot v_t} \quad (6)$$

where

$$v_t = -\frac{\overline{N'w'}}{N} \quad (7)$$

is the particle transfer velocity.

This means that the relative flux uncertainty in many cases becomes larger when the aerosol number concentration decreases. Figure 2 shows the median relative counting uncertainty in flux for the first five OPC channels. The wet season uncertainty (blue bars) are slightly higher than the dry season uncertainty (red bars). The flux uncertainty increases with increasing particle size, and thereby decreasing particle numbers, in both seasons. However, the median uncertainty in flux is lower than 50% for all the five size bins. Median counting errors for fluxes calculated over the whole OPC size range, with interquartile ranges within brackets, were 17% (8%, 32%) in the wet season, and 16% (9%, 29%) in the dry season. The flux of 0.5–2.5 μm particles was only calculated for the wet season data and the flux was there associated with a median uncertainty of 41% (23%, 82%) due to discrete counting.

2.5 Turbulence spectra

Turbulence spectra were calculated in order to investigate the performance of the eddy covariance system. The spectra were obtained using fast Fourier transform (FFT) for data

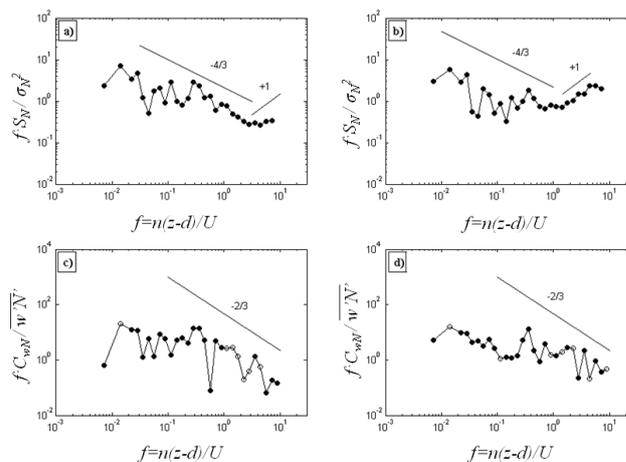


Fig. 3. Averaged frequency weighted variance spectra (S) of particle number concentration within the total OPC size range (a), and within the diameter interval 0.5–2.5 μm (b), and averaged frequency weighted covariance spectra (C) between vertical wind and particle number concentration within the total OPC size range (c) and within the diameter interval 0.5–2.5 μm (d). Variance spectra are normalized with the variance of particle concentration, whereas co-spectra are normalized with the covariance of vertical wind and particle concentration. All spectra are frequency weighted and plotted against dimensionless frequency (f). Data were collected between 12:00 and 12:30 local time on 23 March.

representing time periods of 30 minutes. Variance spectra for horizontal and vertical wind have been presented in Ahlm et al. (2009) for data measured from the same tower during the same measurement campaign. That study also included spectra of the covariance between horizontal and vertical wind.

In this section, some examples of turbulence spectra incorporating measurements using the OPC are presented. Figure 3 shows variance spectra for particle number concentration within the total OPC size range and for the diameter interval 0.5–2.5 μm , and corresponding co-spectra between particle concentrations and vertical wind. The normalized frequency weighted spectra have been averaged over logarithmically changing intervals. The data were collected between 12:00 and 12:30 local time on 23 March. This period was characterized by slightly unstable conditions and an average friction velocity of 0.14 m s^{-1} .

Figure 3 also includes curves representing the expected slopes in the inertial subrange according to Kolmogorov similarity theory, $-2/3$ for the variance spectra and $-4/3$ for the co-spectra (e.g. Kaimal et al., 1972). The spectra in Fig. 3 follow the expected slopes relatively well. The $+1$ slope at higher frequencies in the variance spectra (Fig. 3a–b) is an indication of random noise. The noise level is clearly higher for particles within the diameter interval 0.5–2.5 μm than for particles within the total OPC size range that represent a larger number of particles. However, even though the noise level in some variance spectra in this study were higher than in

Fig. 3b, the co-spectra were not significantly influenced by this noise. This is an indication that the noise is not correlated with the vertical wind speed, and therefore not adding any false flux.

3 Results and discussion

The wet season OPC measurements started at 11 March and lasted until 27 May 2008. The shorter dry season OPC measurements were performed between 1 and 12 August 2008. During the wet season period, there were two clear episodes of higher concentrations of mineral dust, likely transported from the Sahara as been observed in several other studies in the Amazon (e.g. Talbot et al., 1990; Swap et al., 1992; Formenti et al., 2001). The enhanced amounts of atmospheric mineral dust were observed as significantly higher concentrations of the trace elements Al, Fe and Mn from the PIXE analysis. However, the OPC measurements were not running during these two episodes due to technical problems. Influence of mineral dust is therefore expected to be low in the results to be shown in the following sections.

All fluxes, concentrations and meteorological parameters have been averaged over 30 min long intervals in the results that will be presented. Where it is stated that a figure represent median values, medians have been calculated of these 30 min-averaged data.

3.1 Aerosol size distributions

Earlier studies of the aerosol number size distribution over the Amazon rain forest have shown that the two dominating modes occur at diameters 60–90 nm and 130–190 nm (Zhou et al., 2002; Rissler et al., 2006; Pöschl et al., 2010). This means that only a small fraction of the total aerosol number population is analyzed in this study, since the Grimm 1.109 OPC starts at 250 nm. The median aerosol number concentration within the OPC size range in this study was 33 cm^{-3} in the wet season, and 122 cm^{-3} in the dry season. This can be compared with the total aerosol number concentration measured with a condensational particle counter CPC ($D_p > 10 \text{ nm}$) that was 682 cm^{-3} and 1513 cm^{-3} in the wet season and dry season, respectively (Ahlm et al., 2010). Hence, the number of particles measured with the Grimm 1.109 in this study represents roughly 5% and 8% of the total number of particles in the wet and dry season, respectively.

Figure 4 shows median size distributions in the wet and dry season of particle number (a), and volume (b) for each of the 15 OPC channels. Figure 4c shows vertical flux of particles in the first five OPC size bins. The rapid decrease in concentration with increasing particle size is obvious in Fig. 4a. In the volume distribution (Fig. 4b), the concentration increased with increasing diameter for $D_p > 1 \mu\text{m}$, and a coarse mode appears. The ratio of the coarse mode volume to the total volume within the OPC size range is clearly

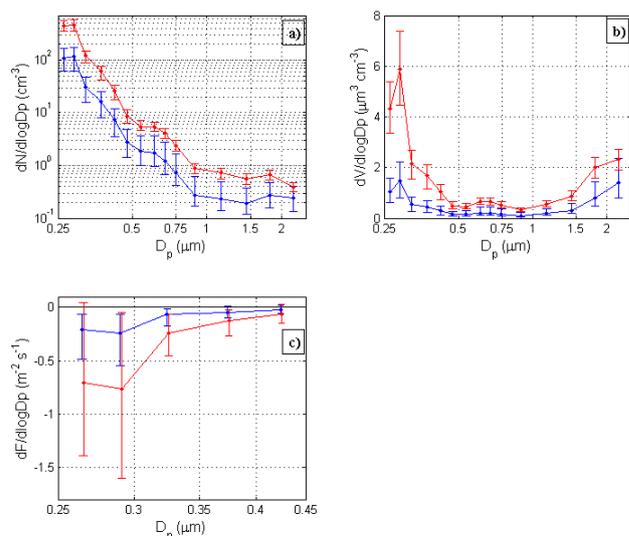


Fig. 4. Median size distribution in the wet season (blue) and in the dry season (red) of (a) numbers of particles, (b) particle volume, and (c) vertical flux of the first five OPC channels from data collected when the friction velocity is above 0.1 ms^{-1} . The vertical bars are 25 and 75 percentiles.

larger in the wet season than in the dry season. This is consistent with observations by for instance Reid and Hobbs (1998) that biomass burning is a large source primarily of accumulation mode particles. The median vertical flux of the first five OPC size bins (Fig. 4c) points downward (negative sign) in both seasons. In these flux medians, half hours with friction velocities lower than 0.1 ms^{-1} have been excluded to reduce the uncertainty in the eddy covariance flux. The flux is clearly larger in magnitude in the dry season than in the wet season. In general, the flux decreases in magnitude with increasing particle diameter and thereby decreasing number of particles.

3.2 Diurnal cycles

3.2.1 Total OPC size range

In this section, diurnal cycles of total concentration and flux of aerosol particles over the whole OPC size range ($0.25 \mu\text{m} < D_p < 2.5 \mu\text{m}$) are investigated. Figure 5 shows median diurnal cycles of aerosol number concentration (a) and vertical flux (b) in the wet and dry season. The aerosol number concentrations within the OPC size range are relatively stable throughout the day, both in the wet and in the dry season (Fig. 5a).

Downward fluxes clearly dominate in the dry season and, as a matter of fact, also in the clean wet season (Fig. 5b). However, the downward flux is considerably larger in magnitude in the dry season. Note that these fluxes include the whole OPC size range, which are likely dominated by

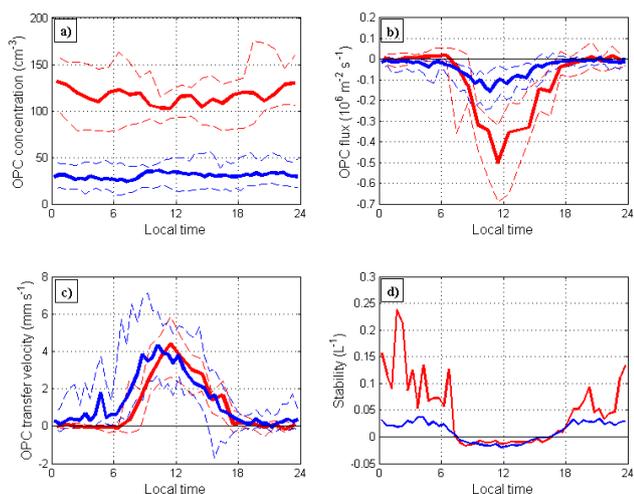


Fig. 5. Median diurnal cycles of aerosol number concentration (a), aerosol number flux (b), transfer velocity (c), and atmospheric stability (d). Blue curves represent wet season data and red curves represent dry season data. (a)–(c) represents the total OPC size range and dashed lines are 25 and 75 percentiles.

the lowest OPC channels, where number concentrations are highest (Fig. 4a). The net downward flux in the wet season implies that sources of primary biogenic aerosol particle are small in comparison with the total dry deposition sink, when considering the whole OPC size range. However, when separately investigating the flux of the larger particles in the OPC size range in Sects. 3.3.3 and 3.4, it will be shown that the role of primary aerosol emission is more important for larger particles.

The particle transfer velocity in the total OPC size range (Fig. 5c) has a very clear diurnal trend with maximum values of 4 mm s^{-1} in both seasons in daytime but considerably lower values at nighttime when turbulence is suppressed by the stable stratification. Particularly stable conditions prevail at nighttime during the dry season (Fig. 5d), whereas unstable conditions dominate in daytime in both seasons.

3.2.2 Size-resolved concentrations

Figure 6 shows median diurnal cycles of the aerosol number concentration for each of the 15 OPC size bins. For each size bin, all median concentrations are normalized with the maximum median concentration during the day within the certain size bin. The median particle concentration of each size bin is shown in Table 1.

In the wet season (Fig. 6a), aerosol number concentrations of the five lowest size bins, representing diameters of $0.25\text{--}0.45 \mu\text{m}$, are high between 09:00 and 12:00 LT. In contrast, particles with diameters of $0.5\text{--}2.0 \mu\text{m}$ have relatively low concentrations in the morning but maximum concentrations around 15:00 LT. Hence, it seems that particles of different size have different diurnal trends. In the dry season (Fig. 6b),

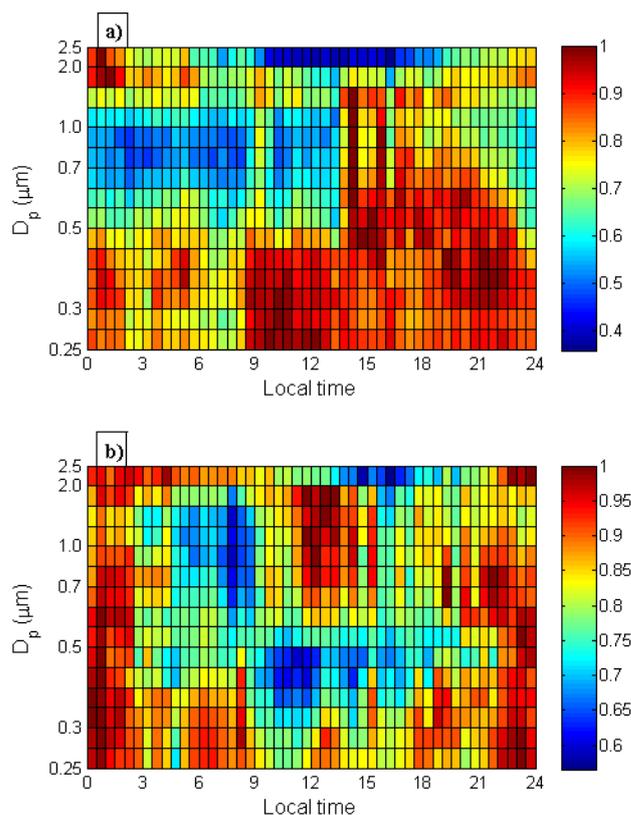


Fig. 6. Median diurnal cycles of the particle concentration within all the 15 OPC size bins in 30 min time resolution. For each size bin, all median concentrations are normalized with the maximum median concentration during the day within the certain size bin. Therefore, the figure gives no information of absolute particle concentrations but instead the diurnal trends in concentration for each size bin. (a) represents wet season and (b) represents dry season. Every half hour represents a median of at least 41 half hours in the wet season and at least 11.5 h in the dry season.

number concentrations of particles with diameters between $0.25 \mu\text{m}$ and $0.65 \mu\text{m}$ are decreasing when the mixed layer grows in the morning. However, in the size bins representing larger particles, concentrations increase at the same time.

For the two largest channels ($D_p = 1.6\text{--}2.5 \mu\text{m}$), wet season concentrations are highest between 23:00 and 02:00 LT at night. It is not likely that this nocturnal concentration maximum is a result of any anthropogenic activity, because the maximum is not apparent in the lower OPC size bins. Since these particles are large, primary biogenic aerosol emission is the most likely explanation. The particle concentrations for these two size bins start to increase already in the evening to reach maximum at night. Higher night time concentrations of coarse mode particles ($2 < D_p < 10 \mu\text{m}$) were observed by Graham et al. (2003a) at the rain forest site at Balbina. Elevated coarse mode aerosol concentrations at night have been observed in the Amazon also by Artaxo et al. (2002) and Guyon et al. (2003b). In the study by Graham et al. (2003a),

Table 1. Median aerosol number concentration for each of the 15 OPC size bins in the wet and dry season.

Channel	Size bin (μm)	Median wet season concentration (cm^{-3})	Median dry season concentration (cm^{-3})
1	0.25–0.28	13.8	52.5
2	0.28–0.30	8.7	33.4
3	0.30–0.35	5.1	19.4
4	0.35–0.40	2.3	8.6
5	0.40–0.45	0.93	3.2
6	0.45–0.50	0.32	0.97
7	0.50–0.58	0.30	0.86
8	0.58–0.65	0.21	0.68
9	0.65–0.70	0.10	0.33
10	0.70–0.80	0.11	0.35
11	0.8–1.0	0.07	0.21
12	1.0–1.3	0.07	0.20
13	1.3–1.6	0.04	0.12
14	1.6–2.0	0.06	0.15
15	2.0–2.5	0.05	0.09

the higher night time concentrations could be directly linked to fungal spores. This is a likely explanation also for the elevated night time concentrations of 1.6–2.5 μm particles in this study. Fungal spores are typically 2–20 μm in diameter (Elbert et al., 2007).

However, it cannot be stated from this study whether primary aerosol emission of 1.6–2.5 μm particles is actually higher at night or if the elevated night time concentration is only a result of the variation in boundary layer depth between day and night. The nocturnal boundary layer in the Amazon Basin is typically 80–180 m (Garstang et al., 1990), while the well developed mixed layer often reaches a depth of ~ 1000 m (Fisch et al., 2004). Therefore, the source of primary aerosol emission would not necessarily have to be stronger at night time just because the concentration maximum occurs at that time. However, as Graham et al. (2003a) pointed out, it is widely believed that the active discharge of several spore types is favored by humid conditions, and the relative humidity frequently reaches 100% at night time over the Amazon rain forest.

3.3 Pristine versus polluted conditions in the wet season

Even though the wet season in the Cuieiras Reserve represents one of the cleanest conditions that can be found on Earth, with aerosol concentrations close to those over remote oceans (Andreae et al., 2009), there is occasionally influence from anthropogenic sources. Figure 7a illustrates how the aerosol number concentration depends on wind direction for each of the 15 size bins in the wet season. The highest particle concentrations are found in the wind sector associated with wind directions between 150 and 180 degrees for all of the 15 size bins. This wind sector represents air transport from Manaus. Maximum BC_e concentrations at K34

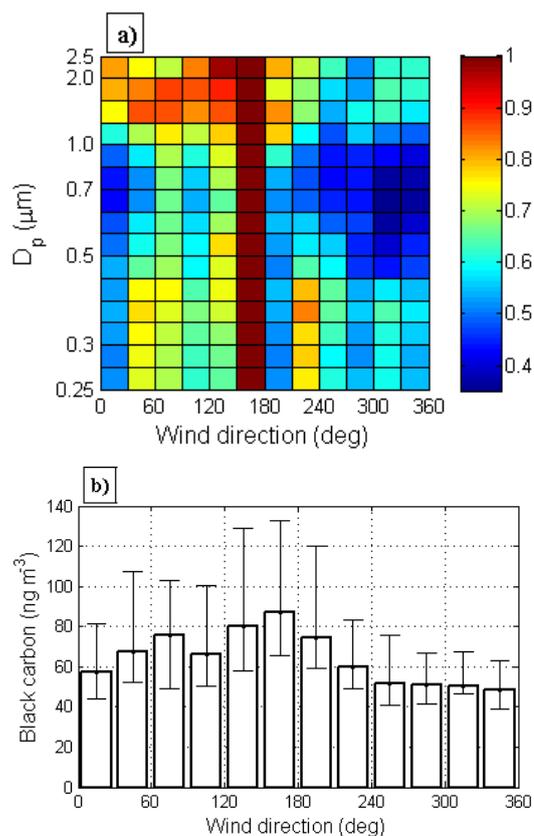


Fig. 7. (a) Wet season median OPC number concentrations as functions of wind direction for each size bin of the OPC. Each median concentration has been normalized with the maximum median concentration of the certain bin. Hence, the colours in the colour bar represent numbers between 0 and 1 where 1 represents the maximum median concentration of the certain size bin. (b) Wet season black carbon concentration medians in constant wind sector intervals. Error bars represent 25 and 75 percentiles.

(Fig. 7b) are associated with the same wind sector. Thus, the city Manaus clearly has an impact on the aerosol number population when the wind is blowing from that direction. The lowest concentrations are found in the wind sectors associated with wind directions between 240 and 360 degrees. These sectors include the direction that is associated with air transport from the research station (approximately 340 degrees). In fact, the minimum concentration is found in the north-westerly wind sector, the specific direction to the diesel generator at the research station. Therefore, it seems that the influence of the research station on aerosol concentration at the top of K34 is negligible for particles within the OPC size range.

In Sects. 3.3.1–3.3.2, particle number fluxes in the diameter range 0.25–0.45 μm are investigated. The corresponding analysis for particles in the diameter range 0.5–2.5 μm is presented in Sect. 3.3.3. Flux dependence on wind direction, differences between pristine and anthropogenic conditions,

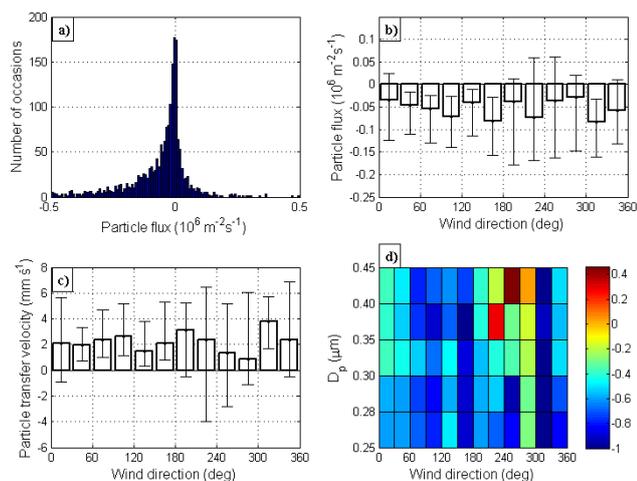


Fig. 8. (a) Histogram of particle ($D_p = 0.25\text{--}0.45\ \mu\text{m}$) flux, (b) Median wet season particle ($D_p = 0.25\text{--}0.45\ \mu\text{m}$) number flux within constant wind sector intervals with error bars representing 25 and 75 percentiles. (c) Median transfer velocities of particles with $D_p = 0.25\text{--}0.45\ \mu\text{m}$. (d) Median wet season size resolved number fluxes within constant wind sector intervals. The color in every square represents a ratio of the median flux within the certain size bin and wind sector to the absolute value of the largest negative median flux within the same size bin but for any wind sector. Hence, colors below zero in the color bar (blue and green) represent downward fluxes, and colors above zero (red) represent upward fluxes. Particle fluxes obtained when the friction velocity is lower than $0.1\ \text{ms}^{-1}$ have been excluded. Every bin represents a median of at least 15.5 h.

and particle transfer velocities will be discussed. All focus will from now on be on the wet season period, since it is hard to find even close to pristine conditions in the dry season, and because the amount of dry season data in this study is too small to reliably represent all wind sectors.

3.3.1 Flux of particles with $D_p = 0.25\text{--}0.45\ \mu\text{m}$

Figure 8a shows a histogram of particle fluxes with $D_p = 0.25\text{--}0.45\ \mu\text{m}$. Obviously, there is a very large dominance of deposition fluxes. Figure 8b illustrates how the flux of particles within the same size range depends on wind direction. Downward fluxes clearly dominate in all wind sectors, indicating that the downward fluxes are not only results of anthropogenic sources. Therefore, it seems that the local source of primary aerosol particles within the particle diameter range $0.25\text{--}0.45\ \mu\text{m}$ is small. The median transfer velocities (Eq. 7) vary between 1.5 and $4\ \text{mm s}^{-1}$ (Fig. 8c). Figure 8d illustrates how the particle flux varies with particle size and wind direction in the same particle diameter range. Downward (negative) particle fluxes dominate in all wind sectors for the three lowest OPC channels. For the fourth and fifth channel, upward fluxes are more frequent in the westerly sector but otherwise downward fluxes dominate

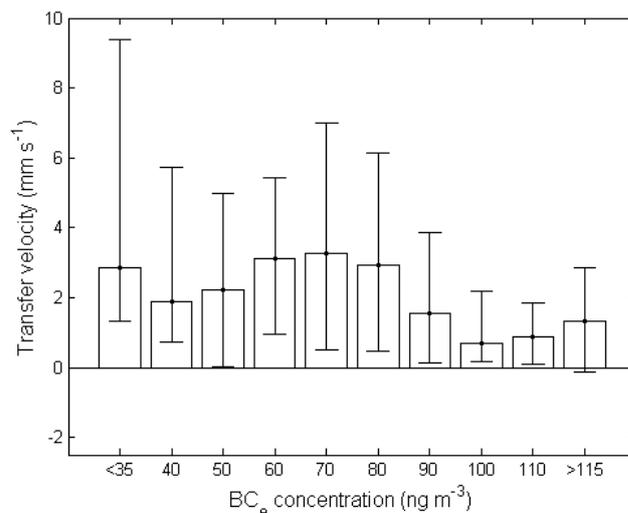


Fig. 9. Median values of transfer velocity for particles with $D_p = 0.25\text{--}0.45\ \mu\text{m}$ within intervals of BC_e concentration. The error bars represent 25 and 75 percentiles. Every median is calculated over at least 38 half hours. Particle fluxes obtained when the friction velocity is lower than $0.1\ \text{ms}^{-1}$ have been excluded.

also for these size bins. The fact that the flux medians for all the five lowest OPC channels are somewhat less negative in the westerly sector, a sector that is associated with low anthropogenic influence (Fig. 7), leads to the question whether deposition of anthropogenic particles may mask a potential primary aerosol source. An efficient way to investigate this is to explore whether the particle transfer velocity changes sign when going from more polluted conditions to extremely clean conditions. Figure 9 shows how the transfer velocity for particles within the diameter size range $0.25\text{--}0.45\ \mu\text{m}$ depends on mass concentration of BC_e . Positive transfer velocities represent downward fluxes according to Eq. 7. Clearly, the transfer velocities are positive even in the cleanest atmospheric conditions. This means that downward fluxes of particles with diameter between 0.25 and $0.45\ \mu\text{m}$ dominate even in the absolute cleanest atmospheric conditions, again indicating that the source of primary biogenic aerosol particles in this diameter range is small.

3.3.2 Size-resolved transfer velocities for particles in the size range $D_p = 0.25\text{--}0.45\ \mu\text{m}$

Figure 10 shows median diurnal cycles of v_t for the five lowest OPC size bins, representing diameters $0.25\text{--}0.45\ \mu\text{m}$. The median transfer velocities are positive throughout the whole cycle, also here indicating a large dominance of deposition fluxes. At nighttime, the stratification is stable and turbulence is thereby suppressed, resulting in small particle fluxes and low transfer velocities at nighttime (as was observed in Fig. 5c–d as well). Throughout the morning, when turbulence gradually increases, transfer velocities also increase.

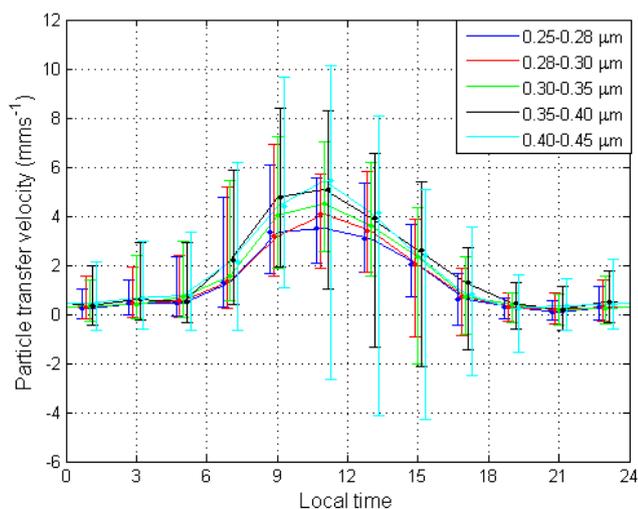


Fig. 10. Median diurnal cycles of particle transfer velocity in each of the five size bins in the diameter range 0.25–0.45 μm . The vertical bars represent 25 and 75 percentiles.

In daytime, v_t for all the five size bins show maximum values between 3.5 and 5.5 mm s^{-1} . These peak values are considerably higher than the median peak in v_t at $\sim 1\text{--}2 \text{ mm s}^{-1}$ observed for the total particle size range ($D_p > 10 \text{ nm}$) in Ahlm et al. (2010). Most studies of aerosol dry deposition show minimum deposition velocities at diameters around 0.1–0.3 μm (Zhang and Vet, 2006), a diameter range where there is no efficient dry deposition mechanism. Brownian diffusion becomes gradually more important with decreasing particle size, and interception and impaction become gradually more important for increasing particle sizes (Slinn, 1982; Zhang et al., 2001). This minimum, in which particles tend to be accumulated, is one of the key processes behind the formation of the accumulation mode. The fact that the transfer velocities obtained in this study (for particles with $D_p > 0.25 \mu\text{m}$) are considerably larger than the transfer velocities for the total particle size range ($D_p > 0.01 \mu\text{m}$) in Ahlm et al. (2010) indicates that for conditions observed over the rain forest, the minimum deposition velocity (as a function of particle size) is located at lower diameters than 0.25 μm . The higher transfer velocities in this study can thereby be interpreted as higher efficiency of interception and impaction of particles with $D_p > 0.25 \mu\text{m}$ compared with the corresponding average efficiencies of these processes in the total particle size range, which is dominated by particles smaller than 0.25 μm in diameter (Zhou et al., 2002; Rissler et al., 2006). In general, v_t increases with increasing particle size for the five size bins in Fig. 10. This means that the increased efficiency of interception and impaction with increasing particle size can also be observed within the relatively small particle diameter range of 0.25–0.45 μm .

Figure 11a shows how the particle transfer velocity, within each of the five size bins, depends on friction velocity. The

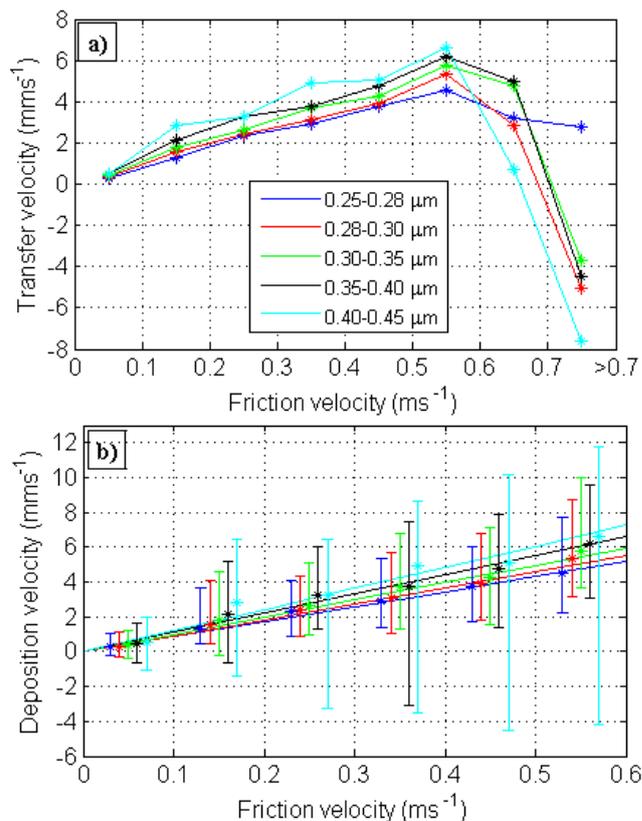


Fig. 11. (a) Median particle transfer velocities in different u_{*a} -intervals. Every median is based on at least 11.5 h. (b) Stars represent median particle transfer velocities in different u_{*a} -intervals, and the vertical bars are the 25 and 75 percentiles. The linear curves represent deposition velocity as a function of D_p and u_{*a} according to Eq. (8).

curves are based on median values of v_t in different friction velocity intervals. The transfer velocity of each particle size interval increases close to linearly with increasing friction velocity until the friction velocity reaches 0.6 ms^{-1} . Above this “threshold”, the transfer velocity of all five size bins instead decreases with further increasing friction velocity, and even becomes negative for all but the bin representing the smallest particles. Negative transfer velocities mean that upward fluxes dominate, possibly indicating that particle emission from the rain forest dominates over the dry deposition sink in very turbulent conditions. However, for friction velocities lower than 0.6 ms^{-1} , 67–76% of the half hour net fluxes were negative (pointing downward) for the five size bins. These percentages of net deposition fluxes are high even compared to aerosol deposition studies at boreal forest sites. Rannik et al. (2003) observed 68% net deposition fluxes at a boreal forest site in Finland, and Pryor et al. (2008) observed less than 60% net deposition fluxes over a forest in Denmark. Hence, due to the high percentage of net particle deposition fluxes in this study, together with the results

in Sect. 3.3.1, it seems reasonable to assume that local primary aerosol emission made a very minor contribution to the observed net particle flux for particles in the diameter range 0.25–0.45 μm when the friction velocity was below 0.6 ms^{-1} . This means that the transfer velocity may be expected to be close to equal to the deposition velocity below this threshold friction velocity. It also seems reasonable to assume that dry deposition velocities continue to increase linearly with increasing friction velocity also at higher friction velocities than 0.6 ms^{-1} . Therefore, it has been assumed that linear fits to the median points in Fig. 11a describes the deposition velocity as a function of friction velocity for each size bin. By combining these linear fits, it is possible to make a parameterization of deposition velocity as a function of particle size and friction velocity for the particle diameter range 0.25–0.45 μm . The dry deposition velocity can then be described as a function of friction velocity and particle size in the form of $v_d = (aD_p + b)u_*$, where v_d is the deposition velocity, u_* is the friction velocity, and a and b are constants. By using the least-square fitting method, the following relation was found for particles in the size range $D_p = 0.25$ –0.45 μm :

$$v_d = (22.12D_p + 2.75)u_* \quad (8)$$

where v_d is given in [mm s^{-1}], D_p is in [μm], and u_* is in [ms^{-1}].

Deposition velocity as a function of particle size and friction velocity, according to Eq. (8) is shown in Fig. 11b, together with the median points from Fig. 11a. It is important to remember that both half hours representing net downward fluxes and half hours of net upward fluxes have been included in the parameterization of the deposition velocity in Eq. (8). The reason for this is that the majority of the net upward fluxes (obtained when the friction velocity is lower than 0.6 ms^{-1}) have been assumed to be caused by random errors, rather than emission. Furthermore, the relative flux error increases with increasing particle size (and thereby decreasing numbers of particles). If only half hours with a net downward flux are included in the calculation, there is a risk that the calculated deposition velocities would increase with increasing particle size, only as a result of enhanced uncertainty in flux. These higher deposition velocities could then be incorrectly interpreted as a result of increased impaction and interception with increasing particle size. Therefore we think it is better to include also half hours with a net upward flux in the parameterization of the deposition velocity, even though there is a risk that a few of the upward fluxes actually could represent real emission also at friction velocities below 0.6 ms^{-1} .

3.3.3 Dependence on wind direction and BC_e for particles with $D_p = 0.5$ –2.5 μm

Figure 6a revealed that the OPC size bins representing particles with diameters between 0.5 and 2.5 μm have rather sim-

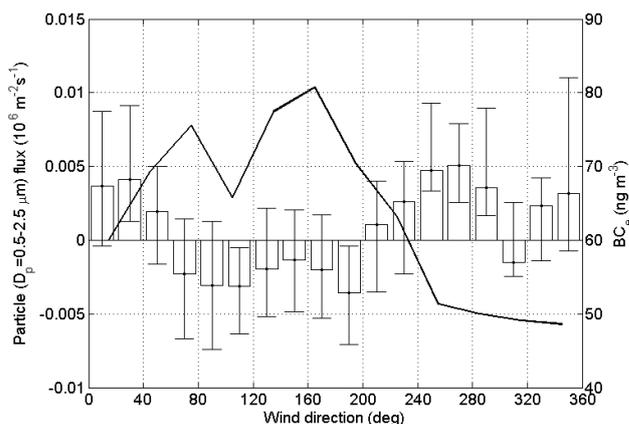


Fig. 12. Median values of vertical flux of particles with $D_p = 0.5$ –2.5 μm within constant wind sector intervals. The error bars represent 25 and 75 percentiles. The black curve represents median BC_e concentrations within the same wind sectors. Every bar is a median over at least 16.5 h. Particle fluxes obtained when the friction velocity is lower than 0.1 ms^{-1} have been excluded.

ilar diurnal cycles of particle concentrations in the wet season, with a well defined maximum in concentration during the afternoon (at least for particles with $D_p < 1.5 \mu\text{m}$). This afternoon peak in concentration is not apparent in the diurnal cycles of the smallest particles within the OPC size range.

Since the diurnal cycle of the total OPC size range flux (Fig. 5b) is likely dominated by fluxes of small and thereby numerous particles, it is necessary to investigate the fluxes of the larger particles separately. However, as was stated earlier, it is not possible to calculate size-resolved fluxes of each size bin of the larger particles separately, since the numbers of particles in these bins are too low compared to the resolution of the OPC. Therefore, we have calculated fluxes over one single large diameter interval, $D_p = 0.5$ –2.5 μm , to increase the number of particles to at least an order of magnitude higher than the resolution of the OPC. However, the concentration decreases with increasing particle size also within this size interval. Approximately 50% of particles in the diameter range 0.5–2.5 μm are located between 0.5 and 0.65 μm .

Figure 12 shows how the net aerosol number flux of particles within this large size interval varies with wind direction. A black curve showing the BC_e dependence on wind direction (from Fig. 7b) has been added to the figure. Interestingly, downward fluxes dominate in the southerly to north-easterly wind sector where concentrations of BC_e are higher, while upward fluxes dominate in most other wind sectors, where concentrations of BC_e are lower. Hence, it seems that in wind sectors where the influence of anthropogenic sources is low, local primary aerosol emission dominates over the dry deposition sink. On the other hand, in wind sectors where the influence of anthropogenic sources is larger, it seems as deposition of anthropogenic particles masks the local source of primary aerosol particles.

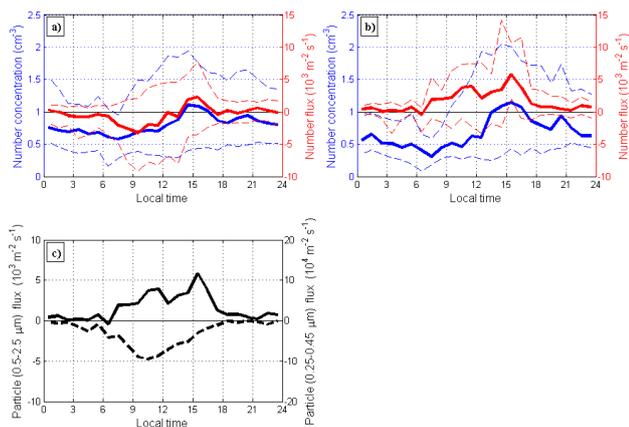


Fig. 13. (a) Wet season median diurnal cycles of (a) particle concentration (blue) and flux (red) of 0.5–2.5 μm particles with all wind sectors included, (b) particle concentration (blue) and flux (red) of 0.5–2.5 μm particles with the wind sector 60–200 degrees excluded, and (c) flux of 0.5–2.5 μm particles (solid line) and 0.25–0.45 μm particles (dashed line). Every median is taken over a one hour time interval and represent at least 62 values in (a) and at least 14 values in (b and c).

The fact that local primary aerosol emission seems to be important in the 0.5–2.5 μm interval, in terms of number concentrations, is not a contradiction to the observations by Chen et al. (2009) who found that the aerosol mass of particles ($D_p < 1 \mu\text{m}$) at the same site during the wet season was dominated by secondary aerosol. The particle diameter interval 0.5–1.0 μm is in a range where the aerosol volume has very low values in the volume distribution in Fig. 4b. Therefore, the aerosol mass in this interval is a very small fraction of the submicron aerosol mass measured by Chen et al. (2009).

3.4 Diurnal cycles of fluxes of particles with diameters of 0.5–2.5 μm

Figure 13a shows the wet season median diurnal cycles of flux and concentration of particles with $D_p = 0.5\text{--}2.5 \mu\text{m}$ when all wind sectors are included. The diurnal cycle of the net flux is somewhat dominated by deposition fluxes even though there is a peak in upward flux around 15:00 LT. However, when the wind sector 60–200 degrees (associated with higher anthropogenic influence in Fig. 12) is excluded, the diurnal cycle is clearly dominated by upward particle fluxes (Fig. 13b). These appear after sunrise at 06:00 LT and increase in magnitude throughout most of the day to reach a maximum also here around 15:00 LT. The fact that the largest net upward fluxes appear in the afternoon means that the sign of the flux cannot be a result of entrainment. At 15:00 LT, the mixed layer is well developed (Fisch et al., 2004). The net upward fluxes are instead most likely a result of primary biogenic aerosol emission, and the fact that they peak during afternoon also suggests that the source here cannot likely be

related to the potential emission observed as upward particle fluxes during early morning in Ahlm et al. (2010). However, that study focused on the total aerosol number population ($D_p > 10 \text{ nm}$) whereas the fluxes discussed here only include particles in the diameter range 0.5–2.5 μm . The contribution in number concentration from this interval makes a negligible contribution to the total number concentration studied in Ahlm et al. (2010). Therefore, there is no contradiction between the observed net emission fluxes here, peaking in the afternoon, and the early morning net upward fluxes observed in Ahlm et al. (2010). An implication, however, is that the observed early morning upward fluxes in Ahlm et al. (2010) must be associated with smaller particles than 0.25 μm , since no peak in upward flux is observed in the morning for particles within the OPC size range in this study.

At the same time as the maximum net upward flux appears, the particle concentration is at its maximum (Fig. 13b). Throughout the later part of the afternoon and early evening, the net emission flux decreases in parallel with decreasing particle concentrations. Thus, an interesting question is whether emission can explain the increase in aerosol number concentration during the day and the peak in concentration at 15:00 LT. By using values of the median net emission flux in Fig. 13b, it is possible to make a rough estimation of the expected increase in mixed layer particle concentration resulting from the flux. A net source of $\sim 4000 \text{ particles m}^{-2} \text{ s}^{-1}$ would need roughly seven hours to increase the particle number concentration by $\sim 0.1 \text{ cm}^{-3}$ in an assumed 1000 m thick mixed layer. Here it has been assumed that the emitted particles are evenly distributed in this 1000 m deep layer. In reality, a particle source at the surface will build up a particle gradient that is largest close to the source but decrease with height within the surface layer. This means that an emission from the surface will have a somewhat higher influence on the particle concentration at the top of K34 than higher up in the mixed layer. However, even when considering this, the increase in particle number concentration with $\sim 0.5 \text{ cm}^{-3}$ between 08:00 and 15:00 LT seems too high to be explained only by the observed net emission flux. Hence, even though emission contribute to the median number concentration increase during the day, it seems that some additional source mechanisms are necessary to explain the diurnal cycle of the particle number concentration within the 0.5–2.5 μm diameter range. One of these additional mechanisms could be in-cloud processing by Aitken or smaller accumulation mode particles (Zhou et al., 2002).

Figure 13c shows a comparison between the median diurnal cycle of the net vertical flux of particles with $D_p = 0.5\text{--}2.5 \mu\text{m}$ (from Fig. 13b) and the corresponding diurnal cycle for particles with $D_p = 0.25\text{--}0.45 \mu\text{m}$. In both these median diurnal cycles, the anthropogenic sector has been excluded. To exclude the anthropogenic sector does not change the sign of the flux of 0.25–0.45 μm particles. At night time, fluxes of particles within both diameter intervals are low. However, when turbulence increases in the morning, net upward

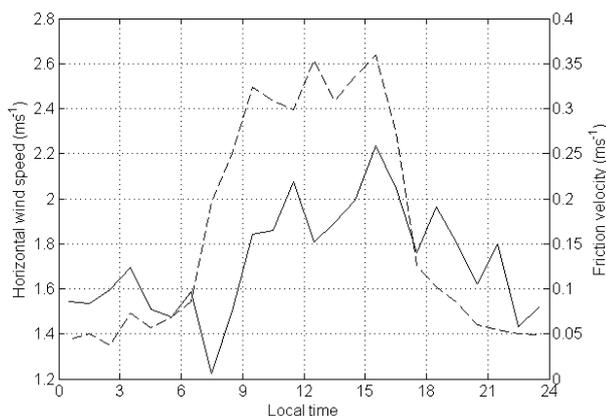


Fig. 14. Median diurnal cycles of horizontal wind speed (solid line) and friction velocity (dashed line). Data collected when wind directions between 60 and 200 degrees prevail have been excluded.

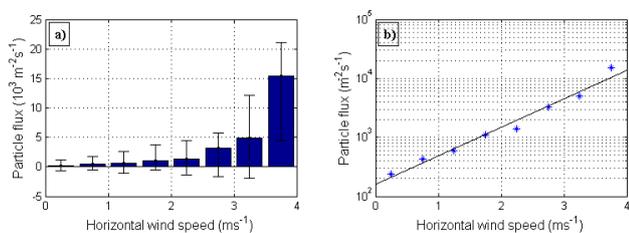


Fig. 15. (a) Wet season median vertical flux of 0.5–2.5 μm particles within constant wind speed intervals. The wind sector 60–200 degrees have been excluded. (b) Stars represent the median values from (a) and the curve is a log-linear fit to the median values. Every bar represents a median over at least 16.5 h.

fluxes of the 0.5–2.5 μm population and net downward fluxes of the 0.25–0.45 μm population appear. We will now look a bit closer at the net emission flux of the 0.5–2.5 μm particles in next section.

3.5 Source mechanism of primary biogenic particles with diameter 0.5–2.5 μm

The median diurnal cycles of horizontal wind speed and friction velocity in the wet season (when the anthropogenic wind sector is excluded) can be seen in Fig. 14. When comparing the diurnal cycles of these two parameters with the corresponding diurnal cycle of the net flux of the 0.5–2.5 μm particles (Fig. 13b–c), it appears as the net emission flux might be more related to wind speed than friction velocity. Actually, both peaks in flux (at 11:00 and 15:00 LT) appear at exactly the same time as the peaks in wind speed, indicating that the net emission flux is correlated with horizontal wind speed. Linking the net emission flux to wind speed instead of friction velocity is also preferable since wind speed is easier to measure and is a parameter that is more often available than friction velocity.

The net emission flux of 0.5–2.5 μm particles increases with increasing wind speed in Fig. 15a. Linear regression of the logarithm of the net flux versus horizontal wind speed yielded the following equation:

$$\log_{10} F = 0.48 \cdot U + 2.21 \quad (9)$$

where F is the net emission flux of 0.5–2.5 μm particles [$\text{m}^{-2}\text{s}^{-1}$] and U is the horizontal wind speed [ms^{-1}] at the top of K34. R^2 between Eq. (9) and the binned data (Fig. 12) is 0.98.

An important point is that the net emission flux described by Eq. (9) also includes contribution from dry deposition. Therefore, Eq. (9) is likely to somewhat underestimate the actual emission. Dry deposition velocities are often estimated by using a dry deposition model as e.g. Slinn (1982) or Zhang et al. (2001). The reason why no attempt has been made to correct Eq. (9) for deposition by using a model is that dry deposition models are very sensitive to parameters as e.g. roughness length, displacement height, drag coefficients and collector properties. All these parameters are associated with very large uncertainties for this site, why an attempt to correct for deposition would likely result in further increased uncertainty. Another important point associated with Eq. (9) is that it cannot be stated whether the wind is only responsible for transporting particles out of the canopy layer, or if the wind (and the turbulence it creates) has any direct impact on the actual emission from the specific source.

Rainfall is another potential triggering mechanism of aerosol emission. Many fungal spore types have been observed to increase in concentration during and after rainfall (Elbert et al., 2007). Even though particle fluxes measured during rainfall have been excluded in the data presented in this study, rainfall could still have an impact on the data because of differences in wetness of the surface one half hour before and after the rainfall. One argument for rainfall being important for the observed net emission fluxes in this study is that wet season rainfall on average was at maximum at 15:00 LT (Fig. 16a), just like the net emission flux. However, the period between 06:00 and 12:00 was the period with least occurrence of rainfall. This period includes the first maximum in net emission flux at 11:00 LT, and in this period the wind speed increases rather simultaneously with the increase in net emission flux. This is an argument against rainfall being the main triggering mechanism of the observed daytime emission of 0.5–2.5 μm particles. Actually, when including particle fluxes measured during rainfall and comparing the particle flux before and after rainfall, it turns out that the net emission flux on average is slightly larger before rainfall than after. This is a second argument against rainfall being the driving mechanism of primary aerosol emission in this study. Nor does solar radiation seem to generate the emission directly, since the photosynthetic radiation is at maximum at noon, or even an hour before, as a result of increasing cloudiness in the afternoon (Fig. 16b).

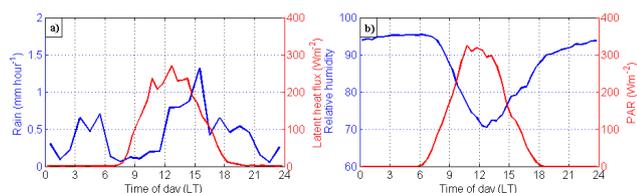


Fig. 16. (a) Mean diurnal cycle of rain amounts (blue) and median diurnal cycle of latent heat flux (red) during the wet season. (b) Median diurnal cycles of relative humidity (blue) and PAR (red) in the wet season.

Another potential mechanism for generating particle emission is transpiration from plants. Several biogenic related elements (e.g. K, P, S, Zn) in plants are present in the fluids circulating in plants and can be released from the plant during transpiration (Nemeruyk, 1970). The latent heat flux is at maximum at 12:00 LT and is much lower at 15:00 LT (Fig. 16a) when the net emission flux is at maximum, indicating that the net emission flux is not well correlated with the latent heat flux. However, the latent heat flux includes both evaporation and transpiration. Therefore, the possibility that transpiration contributes to the emission cannot be ruled out.

Gabey et al. (2009) observed that the aerosol number concentration of particles in the size range $0.8 < D_p < 20 \mu\text{m}$ below the canopy at a rain forest site in Borneo, Malaysia, on average peaked at 15:00 LT. The concentration above the canopy, however, did not show the same behavior. The peak in concentration below the canopy in that study was thought to be a result of fungal spore release triggered by raised relative humidity during the afternoon. It cannot be excluded that a similar source mechanism contributes also to the observed net emission of $0.5\text{--}2.5 \mu\text{m}$ particles in this study in the Amazon rain forest. The relative humidity during the wet season was observed to be at minimum around noon followed by increasing relative humidity in the afternoon (Fig. 16b) when the net emission flux is at maximum. However, since net upward fluxes of $0.5\text{--}2.5 \mu\text{m}$ particles are observed also before noon (Fig. 13b), when the relative humidity decreases (Ahlm et al., 2010), increasing relative humidity is at least not the only source mechanism present in this study.

Thus, the observed fluxes of particles in the diameter range $0.5\text{--}2.5 \mu\text{m}$ indicate a source of biogenic particles in the rain forest, but a more specific source mechanism cannot be defined from this study. However, the fact that the net emission flux is correlated with horizontal wind speed makes it possible to describe the emission in models. Even though number concentrations in this diameter range are low, typically $1 \text{ particle cm}^{-3}$, these particles could potentially play an important role as giant nuclei in warm clouds. Primary biogenic aerosol emission has also been observed to be an important source of ice nuclei in the Amazon Basin (Prenni et al., 2009).

To further investigate the actual emission process, particle number fluxes should be measured simultaneously in and above the canopy, similar to the study by Grönholm et al. (2009). This type of measurements could reveal whether the wind actually triggers the emission, or if the correlation between net upward particle fluxes and wind speed results from storage and ventilation processes. Additional aerosol chemistry measurements inside the canopy would increase the possibility to determine what types of particles that are emitted and what sources are responsible for the emission.

4 Summary and conclusions

Size resolved aerosol number fluxes within the particle size range $0.25\text{--}2.5 \mu\text{m}$ diameter were measured with the eddy covariance method from the top of a 53 m high tower over the Amazon rain forest in the Cuieiras Reserve, Brazil. The measurements included a longer period in the relatively clean wet season and a shorter period in the more polluted dry season.

The median aerosol number concentration within the OPC size range ($0.25\text{--}2.5 \mu\text{m}$ diameter) in this study was 33 cm^{-3} in the wet season and 122 cm^{-3} in the dry season, which represents roughly 5% and 8% of the total number of particles in the wet and dry season, respectively. Aerosol number concentration within the two largest size bins, representing particle diameters between 1.6 and $2.5 \mu\text{m}$, were at maximum at night in the wet season. This night time maximum was likely as a result of primary biogenic aerosol emission and may be related to fungal spores as has been observed in previous studies in the Amazon Basin.

The vertical aerosol number flux of the total OPC size range on average pointed downward in both the wet season and dry season. The net deposition flux, however, was considerably larger in the dry season, probably to a large extent due to the much higher number of particles in the dry season.

The investigation of the size resolved fluxes within the OPC size range showed that the sign of the vertical particle flux may differ for different particle sizes. For the five lowest OPC channels, representing particle diameters of $0.25\text{--}0.45 \mu\text{m}$, downward fluxes dominated in more or less all wind sectors. This is an indication that the source of primary biogenic aerosol particles is low in this particle size interval. Transfer velocities in this particle size range were observed to increase with increasing friction velocity and particle diameter.

To be able to investigate the vertical flux also of the larger particles within the OPC size range, nine size bins were summed up to obtain a higher number of counts. In the resulting particle diameter interval, $0.5\text{--}2.5 \mu\text{m}$, the net vertical number flux depended highly on wind direction. In wind sectors with higher anthropogenic influence, deposition fluxes dominated. In the cleaner wind sectors, however, emission fluxes dominated. The net emission flux within the clean sectors is likely explained by primary biogenic aerosol

emission. The net deposition flux in wind sectors associated with higher anthropogenic influence is probably due to deposition of anthropogenic particles, masking the biogenic aerosol emission.

The net emission number flux of the 0.5–2.5 μm particles peaked at 15:00 LT in the afternoon and was well correlated with horizontal wind speed through the equation

$$\log_{10} F = 0.48 \cdot U + 2.21$$

where F is the net emission number flux of 0.5–2.5 μm particles [$\text{m}^{-2} \text{s}^{-1}$] and U is the horizontal wind speed [ms^{-1}] at the top of K34.

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